

BIRS Discovery Workshop - Banff, July 2010

## Summary talk: a cosmologist's perspective

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Picture credit: Peter Liversidge

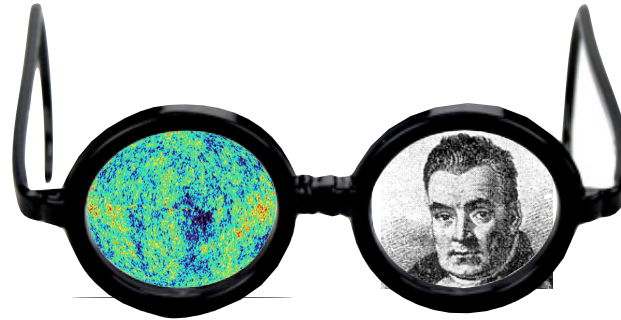
Thanks to Louis, Jim and Richard for organizing  
this workshop!

Lots of interesting talks, stimulating discussions...  
and some clashes!



# A cosmologist's perspective

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- Physicists/statisticians and astro/cosmo vs particle physics
- Look elsewhere effect
- Model selection
- Relevance of 5 sigma for us folks

# Particle physicists vs cosmologists/astro

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- **Methodological**

- Repeatable experiments (counting) vs observations (there is only 1 Universe)
- Frequentist vs mostly Bayesian
- Profiling vs marginalizing
- Priors: “What priors?” vs Often highly relevant prior information
- Selection effects are usually important in cosmology
- Combination of probes necessary in cosmology to break degeneracies (problem: what about systematics?)

# FIRST-YEAR SLOAN DIGITAL SKY SURVEY-II (SDSS-II) SUPERNOVA RESULTS: HUBBLE DIAGRAM AND COSMOLOGICAL PARAMETERS

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*Accepted for publication in ApJS*

Contributions  
by particle  
physicists

## ABSTRACT

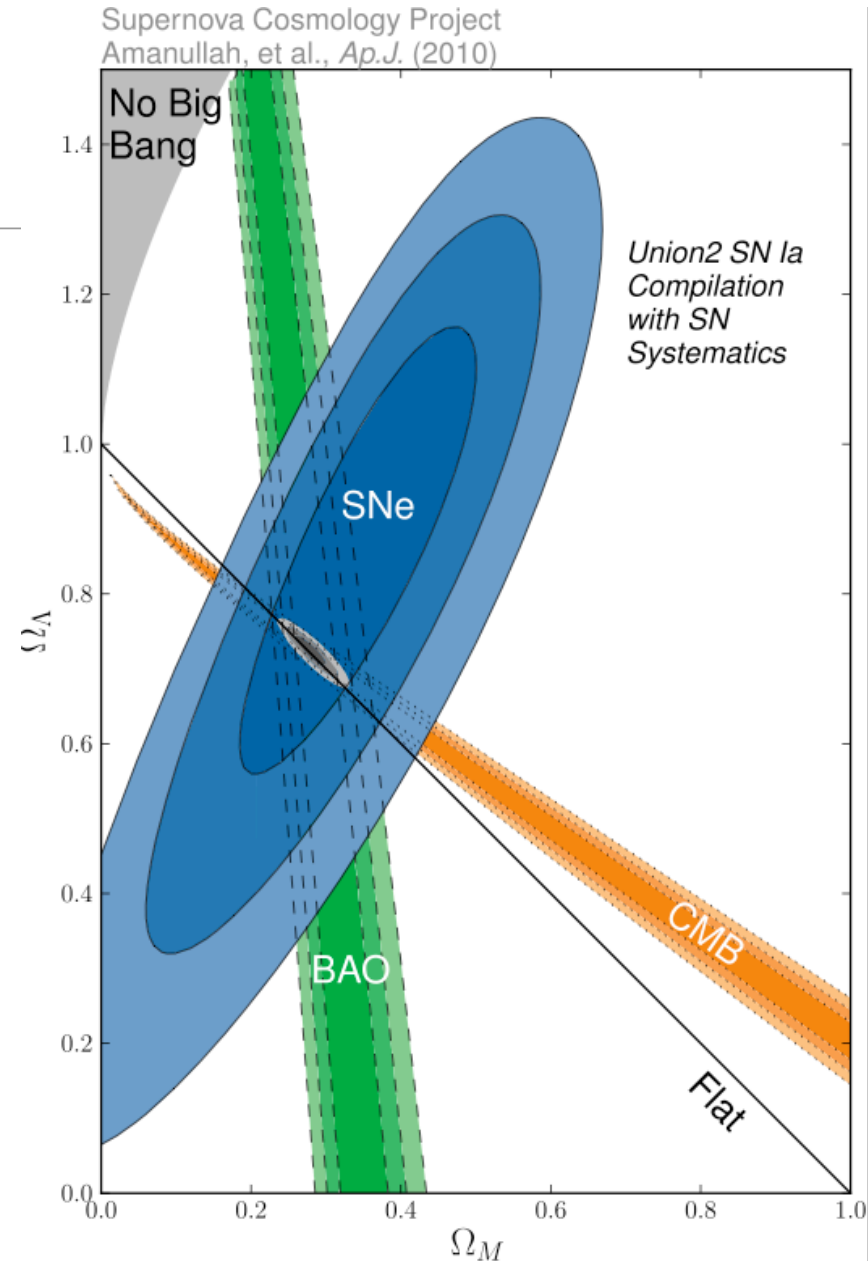
We present measurements of the Hubble diagram for 103 Type Ia supernovae (SNe) with redshifts  $0.04 < z < 0.42$ , discovered during the first season (Fall 2005) of the Sloan Digital Sky Survey-II (SDSS-II) Supernova Survey. These data fill in the redshift “desert” between low- and high-redshift SN Ia surveys. Within the framework of the MLCS2K2 light-curve fitting method, we use the SDSS-II SN sample to infer the mean reddening parameter for host galaxies,  $R_V = 2.18 \pm 0.14_{\text{stat}} \pm 0.48_{\text{syst}}$ , and find that the intrinsic distribution of host-galaxy extinction is well fit by an exponential function,  $P(A_V) = \exp(-A_V/\tau_V)$ , with  $\tau_V = 0.334 \pm 0.088$  mag. We combine the SDSS-II measurements with new distance estimates for published SN data from the ESSENCE survey, the Supernova Legacy Survey (SNLS), the Hubble Space Telescope (HST), and a compilation of nearby SN Ia measurements. A new feature in our analysis is the use of detailed Monte Carlo simulations of all surveys to account for selection biases, including those from spectroscopic targeting. Combining the SN Hubble diagram with measurements of baryon acoustic oscillations from the SDSS Luminous Red Galaxy sample and with cosmic microwave background temperature anisotropy measurements from WMAP, we estimate the cosmological parameters  $w$  and  $\Omega_M$ , assuming a spatially flat cosmological model ( $Fw$ CDM) with constant dark energy equation of state parameter,  $w$ . We also consider constraints upon  $\Omega_M$  and  $\Omega_\Lambda$  for a cosmological constant model ( $\Lambda$ CDM) with  $w = -1$  and non-zero spatial curvature. For the  $Fw$ CDM model and the combined sample of 288 SNe Ia, we find  $w = -0.76 \pm 0.07(\text{stat}) \pm 0.11(\text{syst})$ ,  $\Omega_M = 0.307 \pm 0.019(\text{stat}) \pm 0.023(\text{syst})$  using MLCS2K2 and  $w = -0.96 \pm 0.06(\text{stat}) \pm 0.12(\text{syst})$ ,  $\Omega_M = 0.265 \pm 0.016(\text{stat}) \pm 0.025(\text{syst})$  using the SALT-II fitter. We trace the discrepancy between these results to a difference in the rest-frame UV model combined with a different luminosity correction from color variations; these differences mostly affect the distance estimates for the SNLS and HST supernovae. We present detailed discussions of systematic errors for both light-curve methods and find that they both show data-model discrepancies in rest-frame  $U$ -band. For the SALT-II approach, we also see strong evidence for redshift-dependence of the color-luminosity parameter ( $\beta$ ). Restricting the analysis to the 136 SNe Ia in the Nearby+SDSS-II samples, we find much better agreement between the two analysis methods but with larger uncertainties:  $w = -0.92 \pm 0.13(\text{stat}) \pm 0.10_{-0.33}(\text{syst})$  for MLCS2K2 and  $w = -0.92 \pm 0.11(\text{stat}) \pm 0.07_{-0.15}(\text{syst})$  for SALT-II.

# Precision cosmology

$$\Omega_{\text{cdm}}h^2 = 0.1099 \pm 0.0062$$

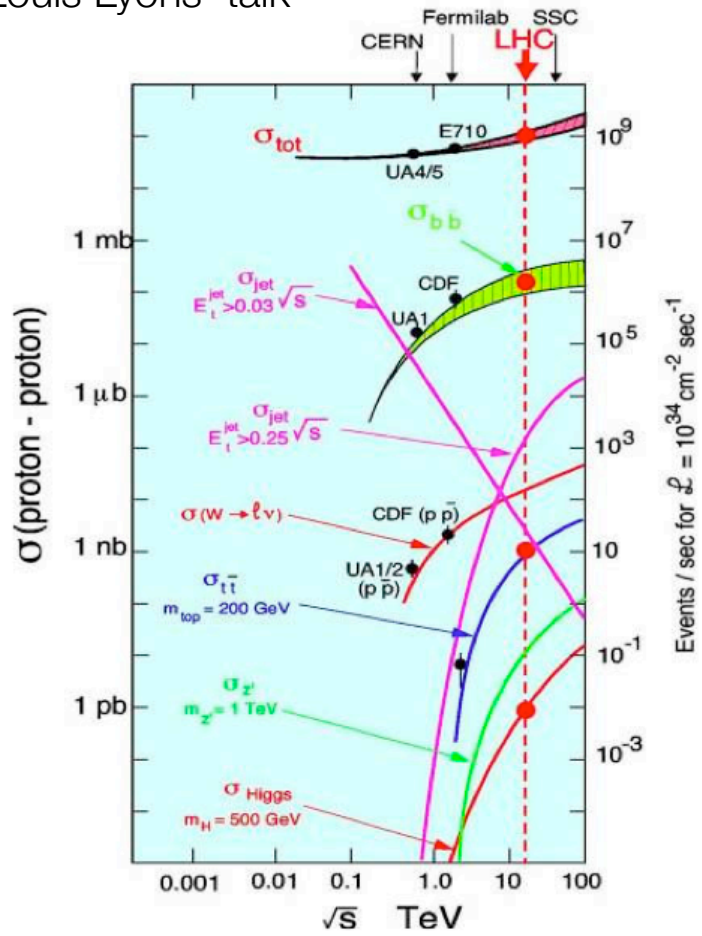
$$\Omega_m = 0.258 \pm 0.030$$

Those are NOT  
confidence regions!

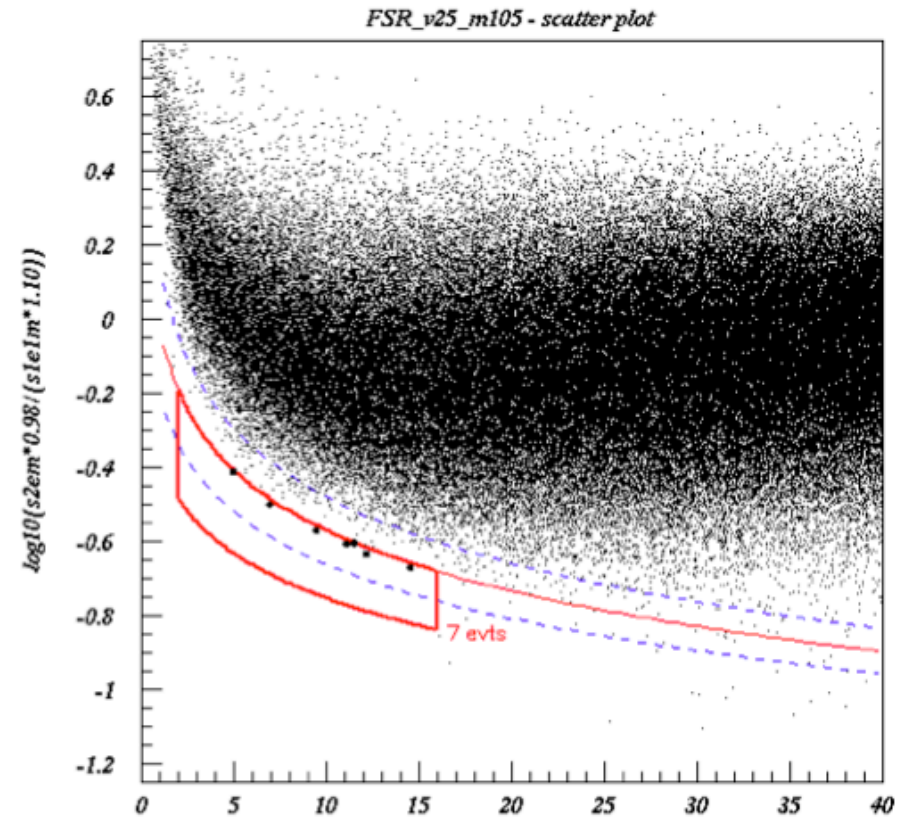


# Needle in the haystack

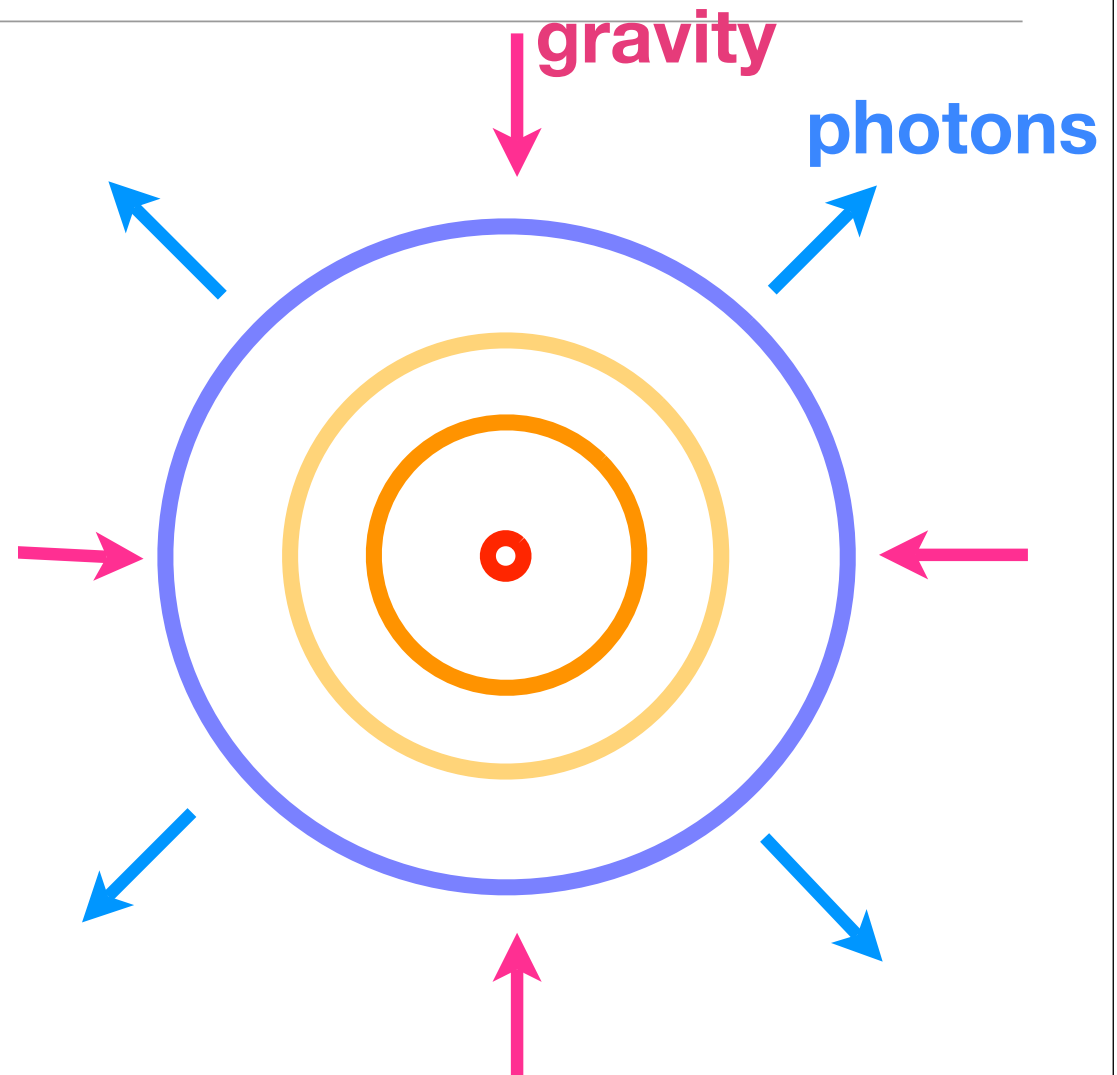
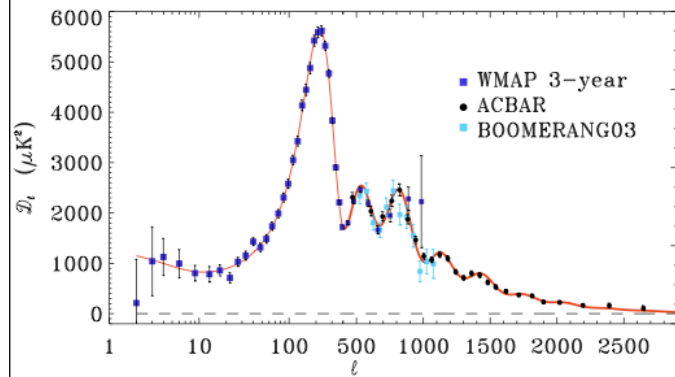
Louis Lyons' talk



Henrique Araujo's talk

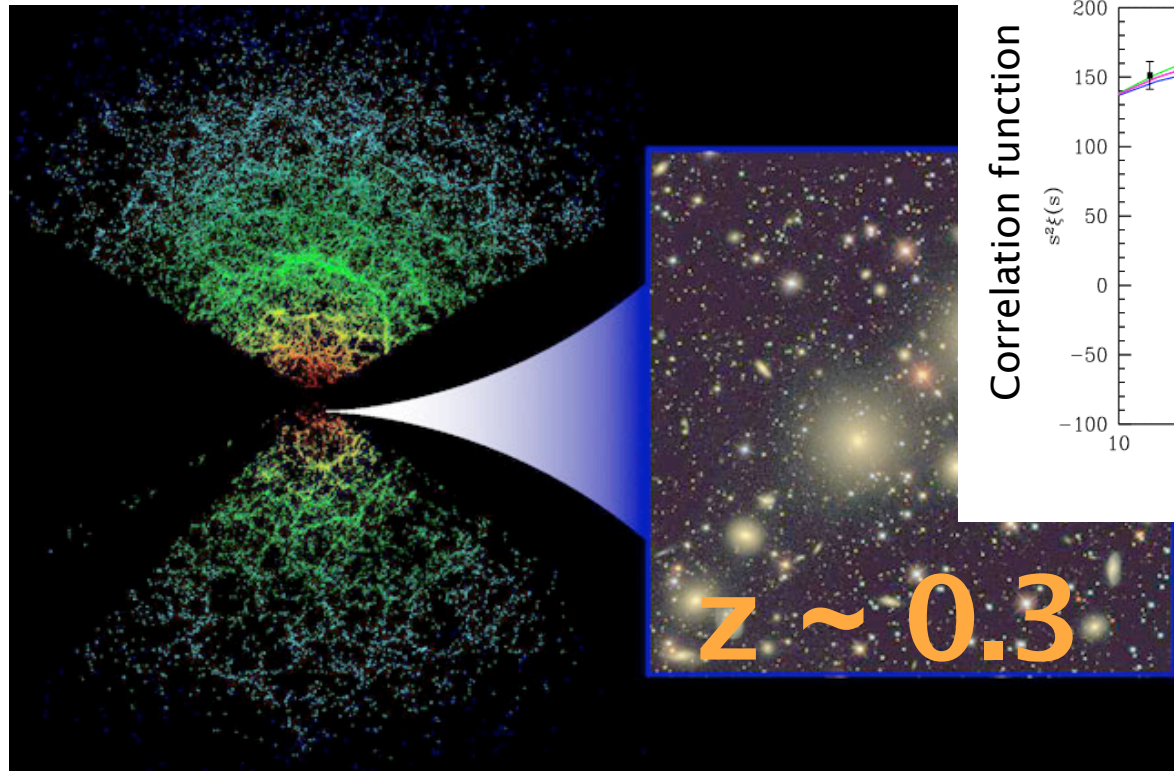


# Cosmic sound

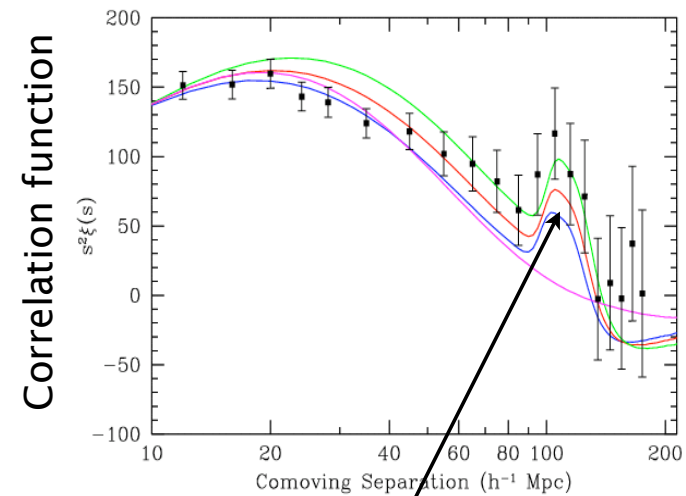




# Needle in the haystack - cosmology



based on  $\sim 50'000$  LRG



Baryonic acoustic oscillations signature  
 $\sim 2.5$  sigma significance

We are looking for extra correlations between galaxies on scales  $\sim 150$  Mpc: this corresponds to (on average) 1 extra galaxy at this preferential separation

# Particle physicists vs cosmologists/astro

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- **Epistemological**

- Particle physicists believe in the existence of THE TRUE MODEL.
- Cosmology is often more pragmatic: “The cosmological concordance model” is more of a phenomenological description of the data (dark matter/dark energy), not necessarily fundamentally motivated in the same way as particle physics models are.
- Frequentist error probabilities vs uncertainties representing degree of lack of knowledge/belief.

# Particle physicists vs cosmologists/astro

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- **Community**

- Data often published in summary form (although this is changing -> Kyle Cranmer's talk) vs full data made public (WMAP, SDSS)
- Large, international collaborations vs smaller, more compact teams (although this is changing: WMAP team ~ 15 people; Sloan ~ 50 people; Planck ~ 500 people, Auger ~ 400 people)
- Large codes often private vs codes usually made public, community input.
- Both communities have now meetings where discussions with statisticians are encouraged (PHYSTAT, Banff, cosmstats, ...)

# Language barriers: a dictionary

Physicist

Statistician

|                          |   |                            |
|--------------------------|---|----------------------------|
| cuts                     | ↔ | filter, selection          |
| large samples statistics | ↔ | ?                          |
| measurements             | ↔ | estimates                  |
| events                   | ↔ | events (different meaning) |
| chi-squared              | ↔ | weighted sum of squares    |
| look elsewhere           | ↔ | multiple comparisons       |
| doh!                     | ↔ | measurement error problem  |
| upper limit              | ↔ | ?                          |
| not on the boundary      | ↔ | the null in the interior   |

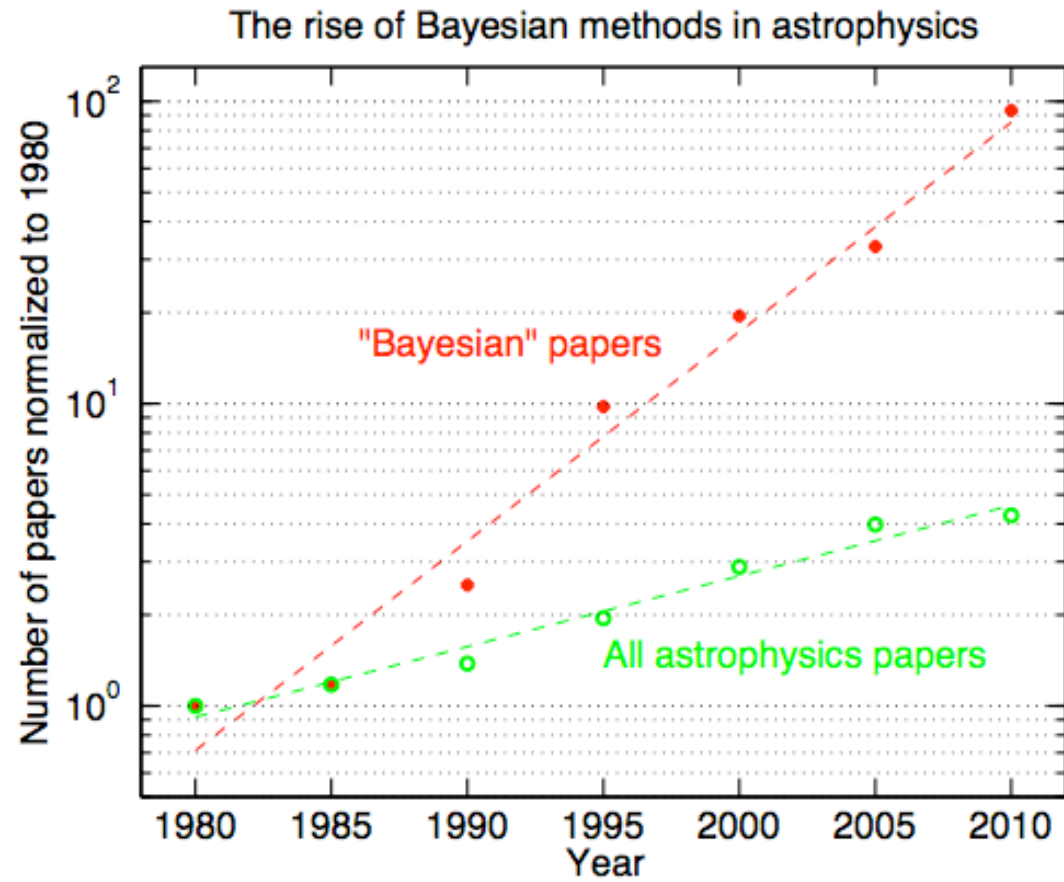
*All of those came up during this workshop*

# Why cosmology is different

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- An observational science (seems obvious, but it has profound consequences)
- Strong selection effects
- Often poorly understood nuisance parameters
- Cosmic variance limited in some cases
- Not clear what the ensemble would be in a frequentist sense!
- Often, somebody's noise is somebody else's signal: This means that often we are interested in  $P(\text{signal, noise} \mid \text{data})$ , so there is no obvious way to classify parameters as “nuisance parameters”.
- We have the sexier pictures!
- As we never properly learnt statistics, we are mostly Bayesians.

# Bayes in the sky



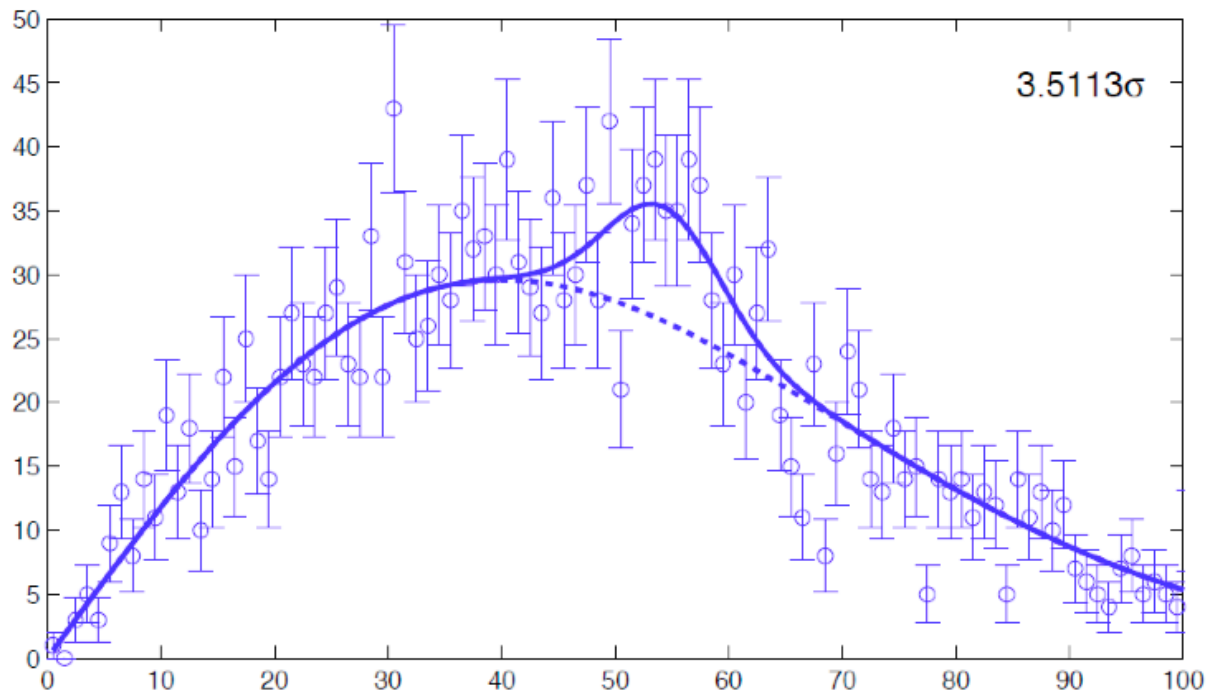
# Look elsewhere effect



Now try to find a few that look like a bear or a dog or something

# Example from Eilam Gross talk: a 3.5 sigma signal?

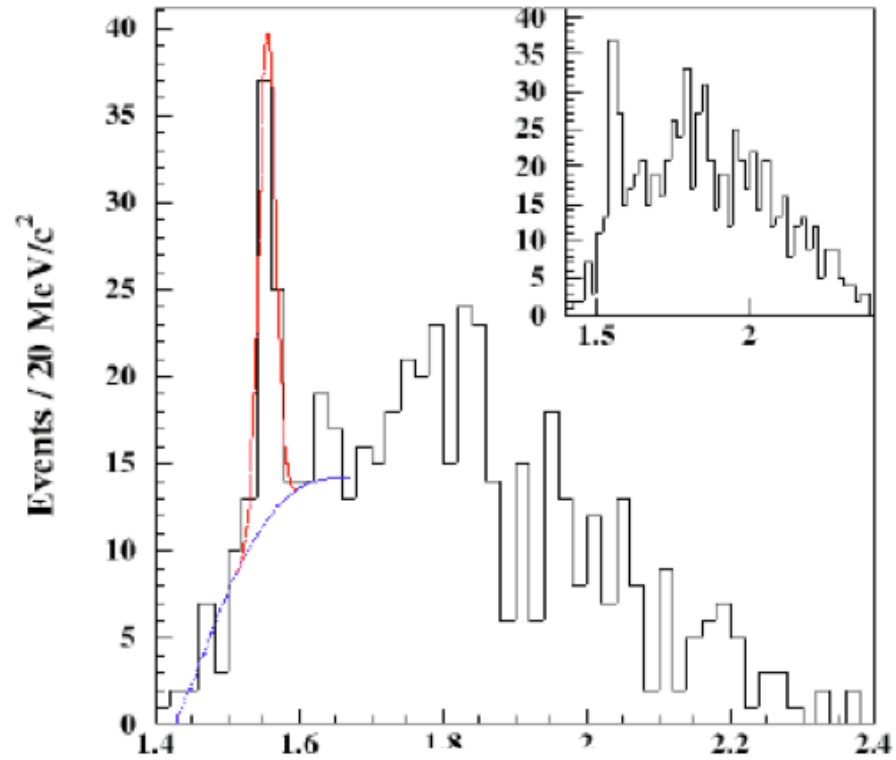
Eilam Gross' talk



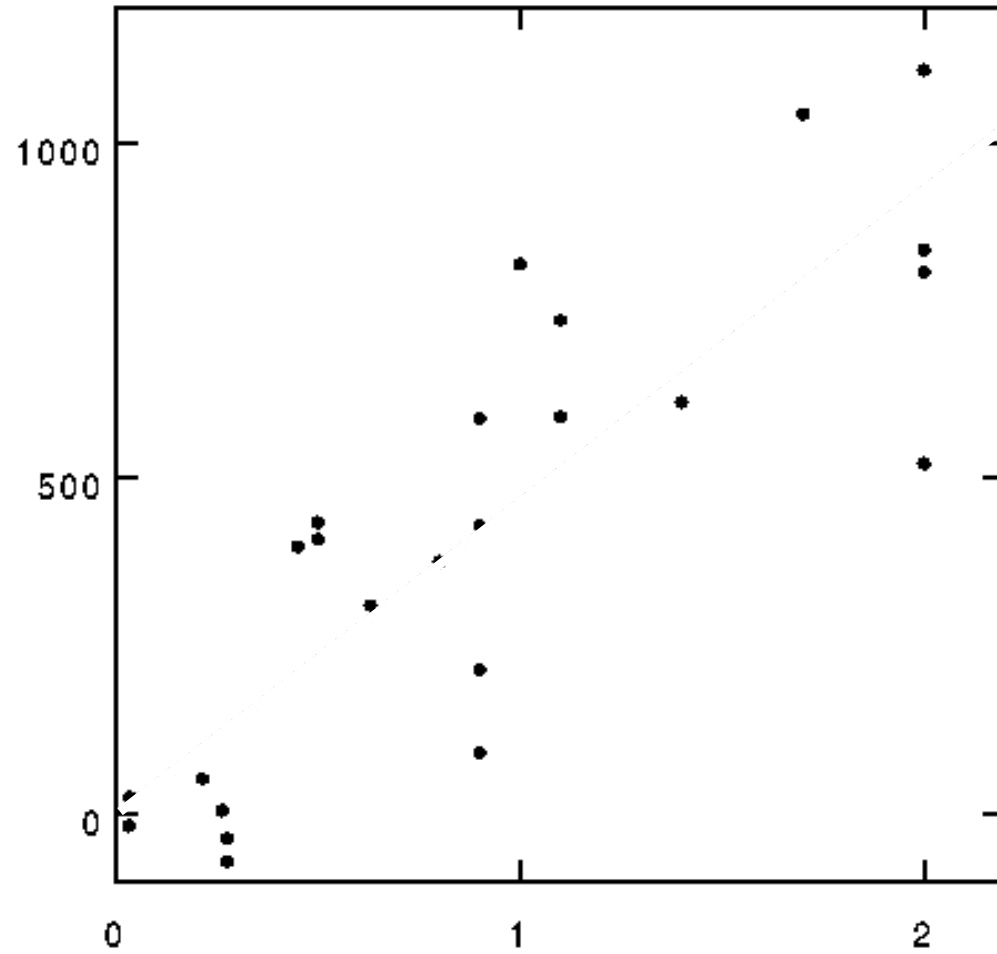


# Pentaquark “discovery”, ooops, maybe not...

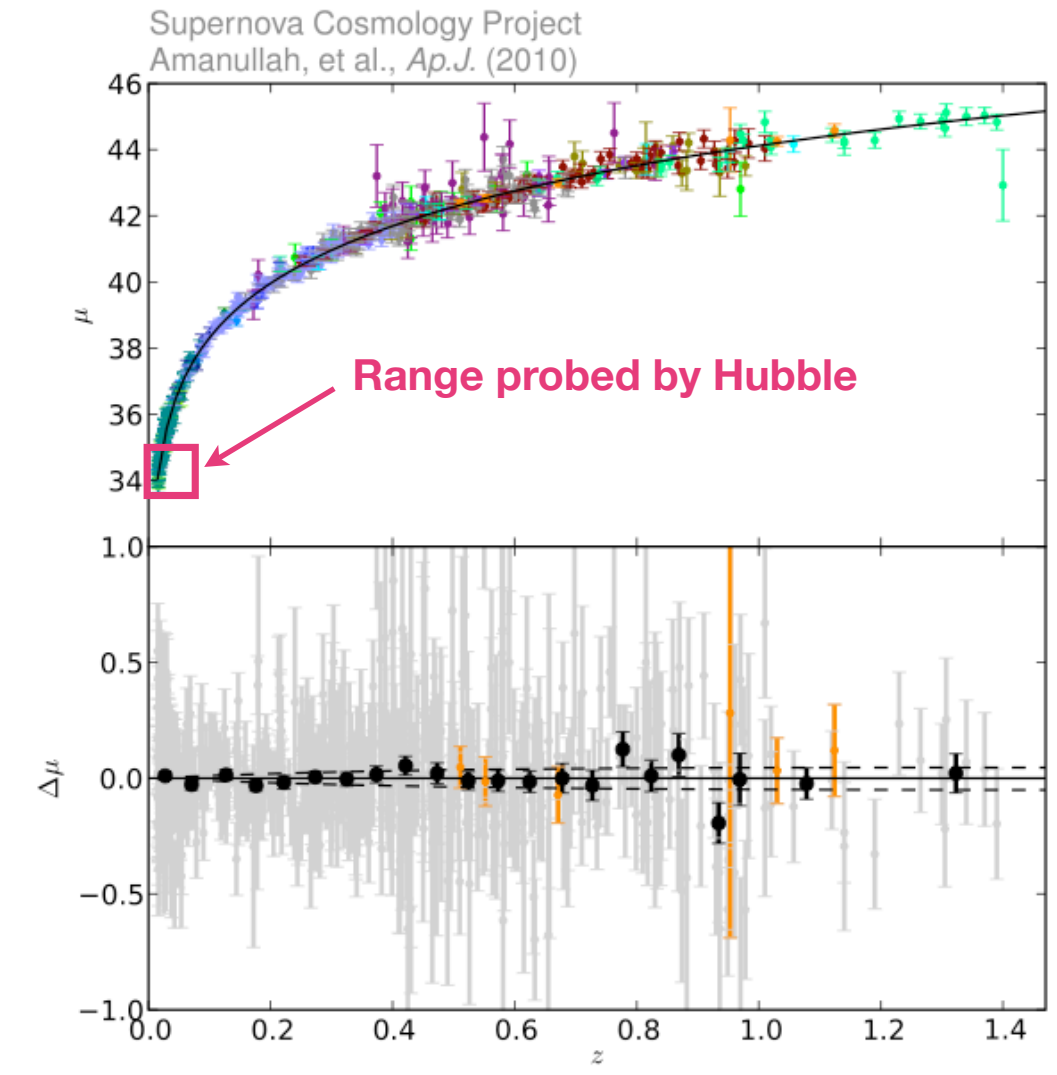
Bob Cousins,  
Louis Lyons



$$v = H_0 * R \text{ (Hubble, 1929)}$$



# Hubble diagram today



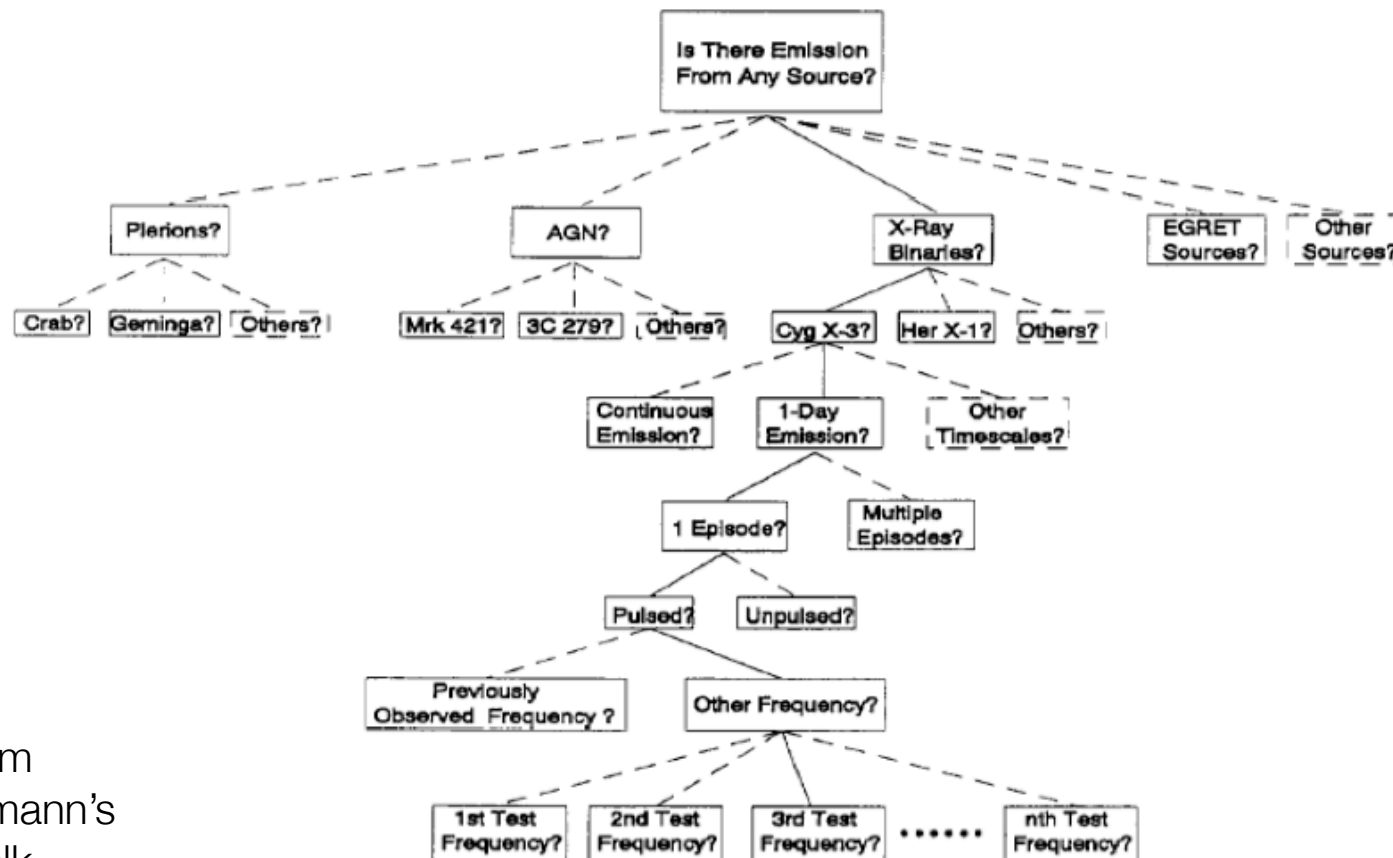
# Look elsewhere effect

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- Where is elsewhere?
- Need to define what “elsewhere” means!
- Do future/possible searches matter? It seems to me that only things you have actually looked at should matter.
- This however brings in the Stopping rule problem. You have to make sure you follow your protocol through!

# Importance ordering: write out a **protocol**

*S.D. Biller / Astroparticle Physics 4 (1996) 285-291*



Jim  
Linnemann's  
talk

# Look elsewhere effect

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- The whole discussion to me gives strong motivation for being a Bayesian about hypothesis testing/model selection. There seems to be no unique/well defined way of defining what “elsewhere” means.
- Jim Berger argued very strongly that the Bayesian answer corresponds to a specific (unique) choice of conditioning statistics for the frequentist testing.
- Things are going to get worse with more complex data sets for which a “single shot” discovery protocol is simply unfeasible.
- Cosmo/astro is typically much more exploratory, with pretty much everybody getting a free hand with the data (see later).

Dangers of looking elsewhere...



Picture credit: Elifam Gross

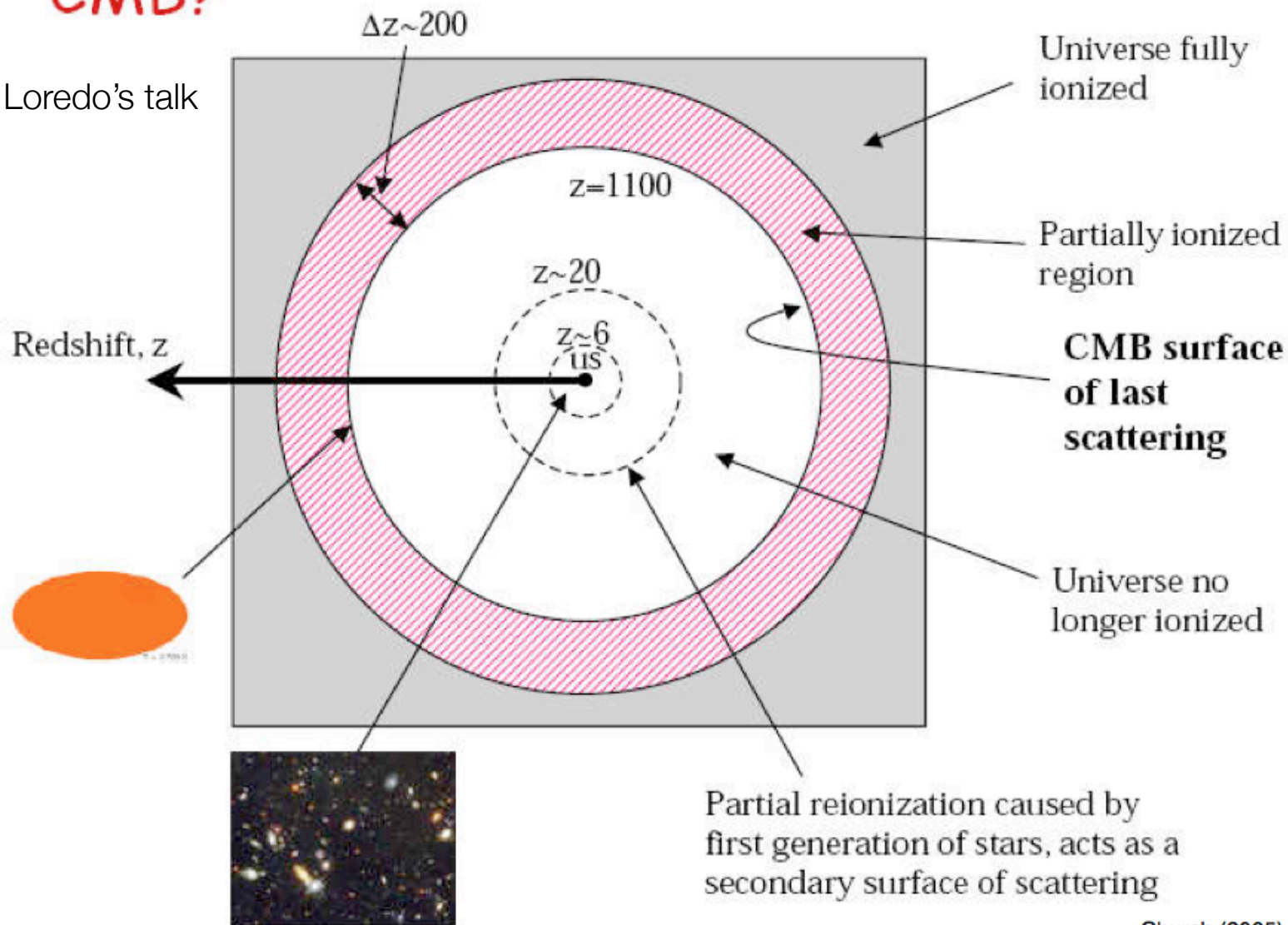
Look elsewhere effect in cosmology:

Large-scale anomalies in the cosmic microwave background



# What are we seeing when we look at the CMB?

Tom Loredó's talk



Church (2005)

# Multi-frequency observations

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WMAP: 5 bands from 23 GHz to 90 GHz

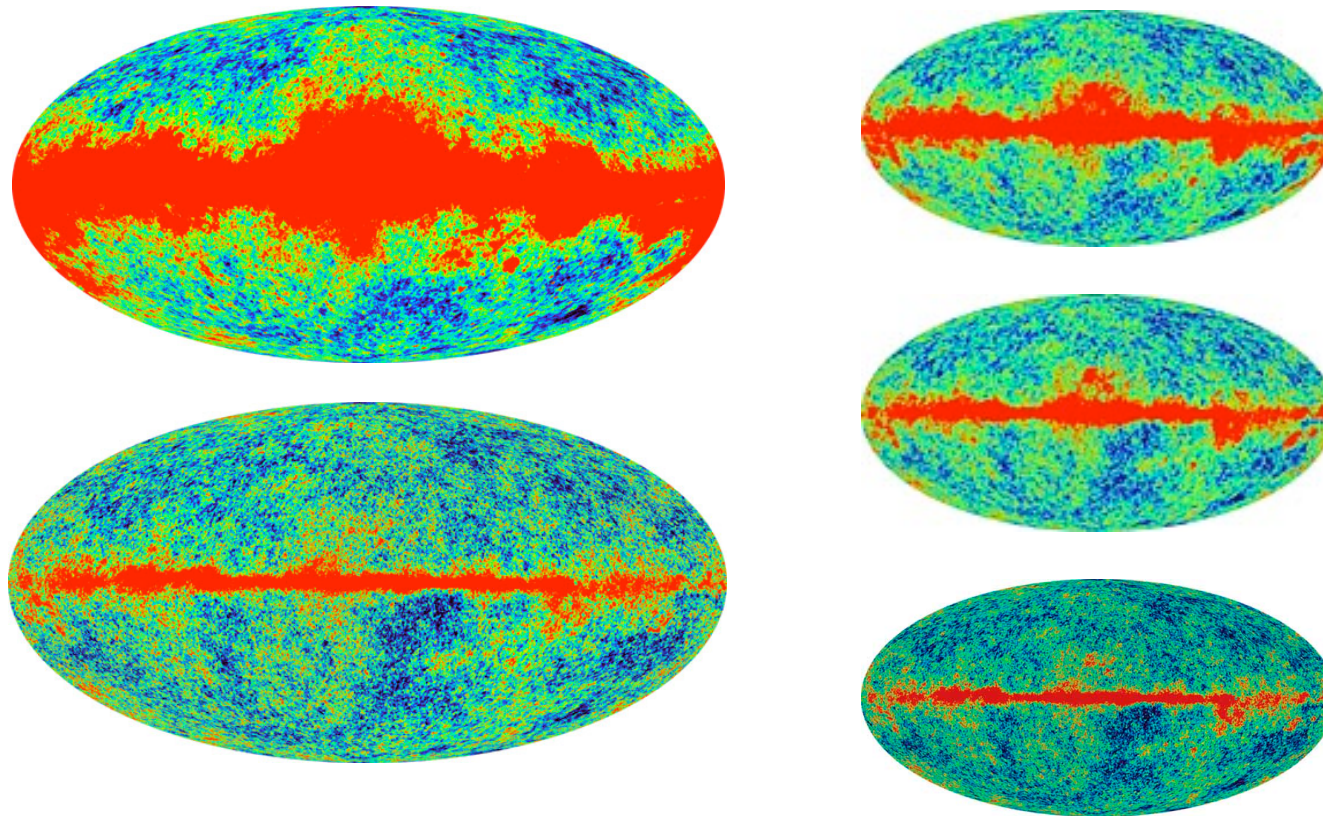
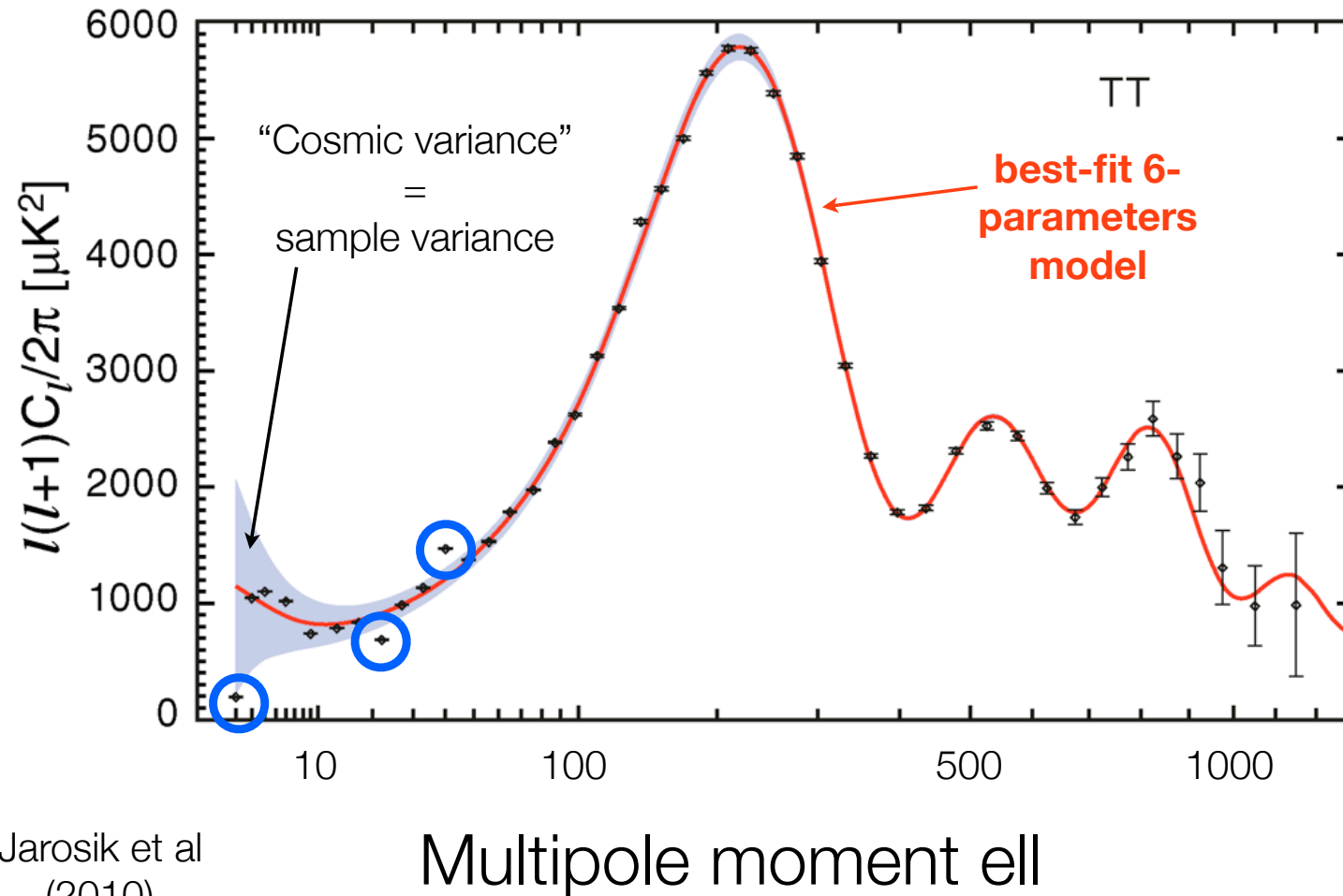


Image: WMAP team

# WMAP 7-years temperature power spectrum



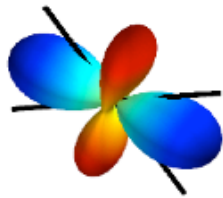
Jarosik et al  
(2010)

## Multipole vectors

$$\frac{\Delta T(\mathbf{n})}{\bar{T}} \equiv \sum_{m=-\ell}^{\ell} a_{\ell m} Y_{\ell m} \equiv A^{(\ell)} \left[ \prod_{i=1}^{\ell} (\hat{v}^{(\ell,i)} \cdot \hat{e}) - \mathcal{T}_{\ell} \right]$$

Characterize spherical harmonics by “multiple vectors”

Quadrupole  
L=2



Octupole  
L=3



L=4



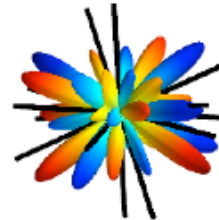
L=5



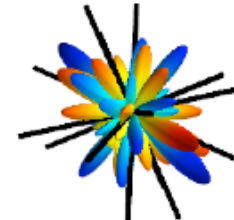
see Copi et al  
(2010) for a  
review and  
references



L=6



L=7



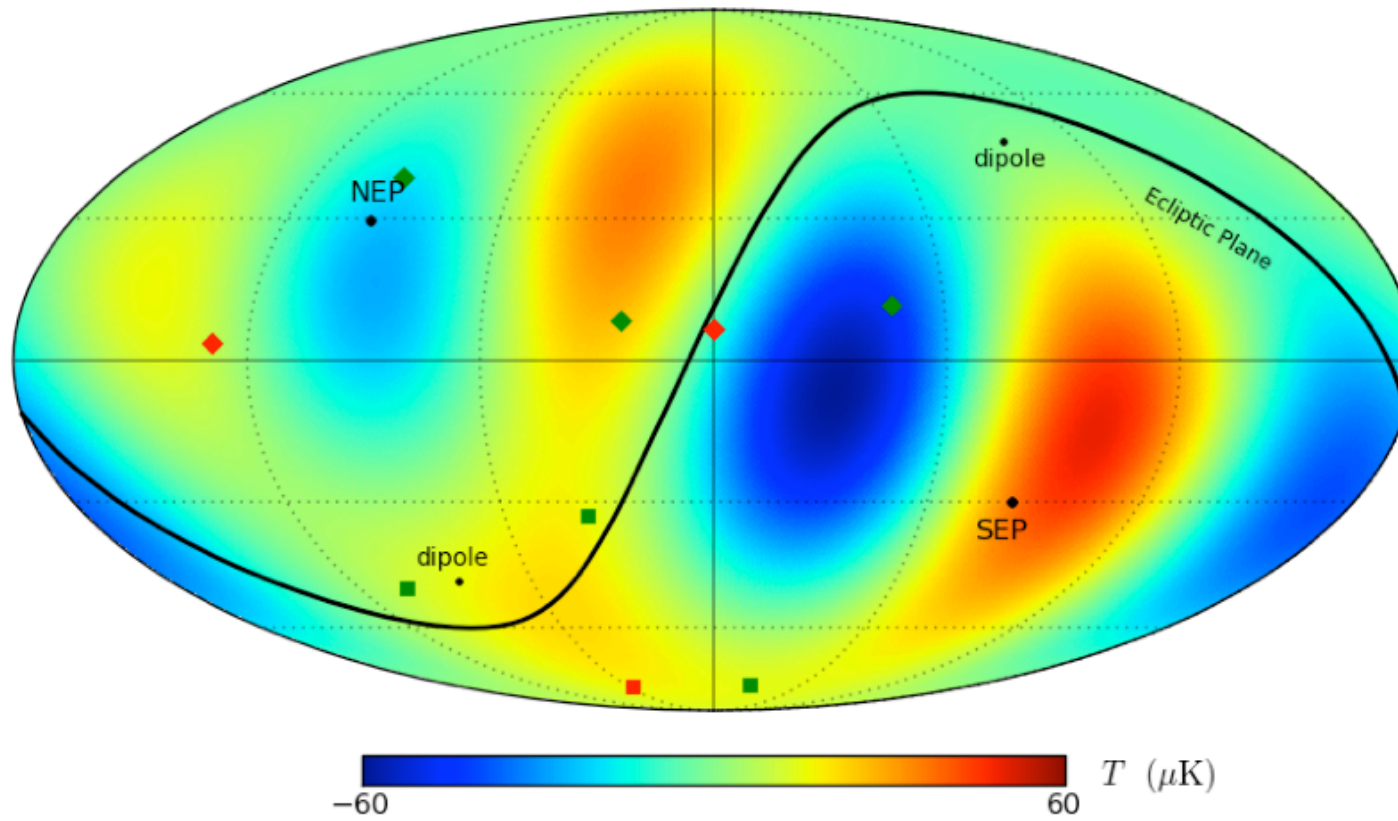
L=8

# Large scale anomalies?

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- Quadrupole ( $l=2$ ) is low
- Four area vectors of Quadrupole and Octupole are mutually close (p-value = 0.004)
- Quadrupole and Octupole aligned with ecliptic (p-value = 0.041)
- Normals to area vector planes aligned with dipole (p-value = 0.003)
- Hot/Cold spots divided by the ecliptic (p-value  $\sim 0.05$ )
- Two-point correlation function vanishes above  $60^\circ$  (p-value  $\sim 0.0002$ )
- All of the above appear to disprove the cosmological principle (isotropy and homogeneity)

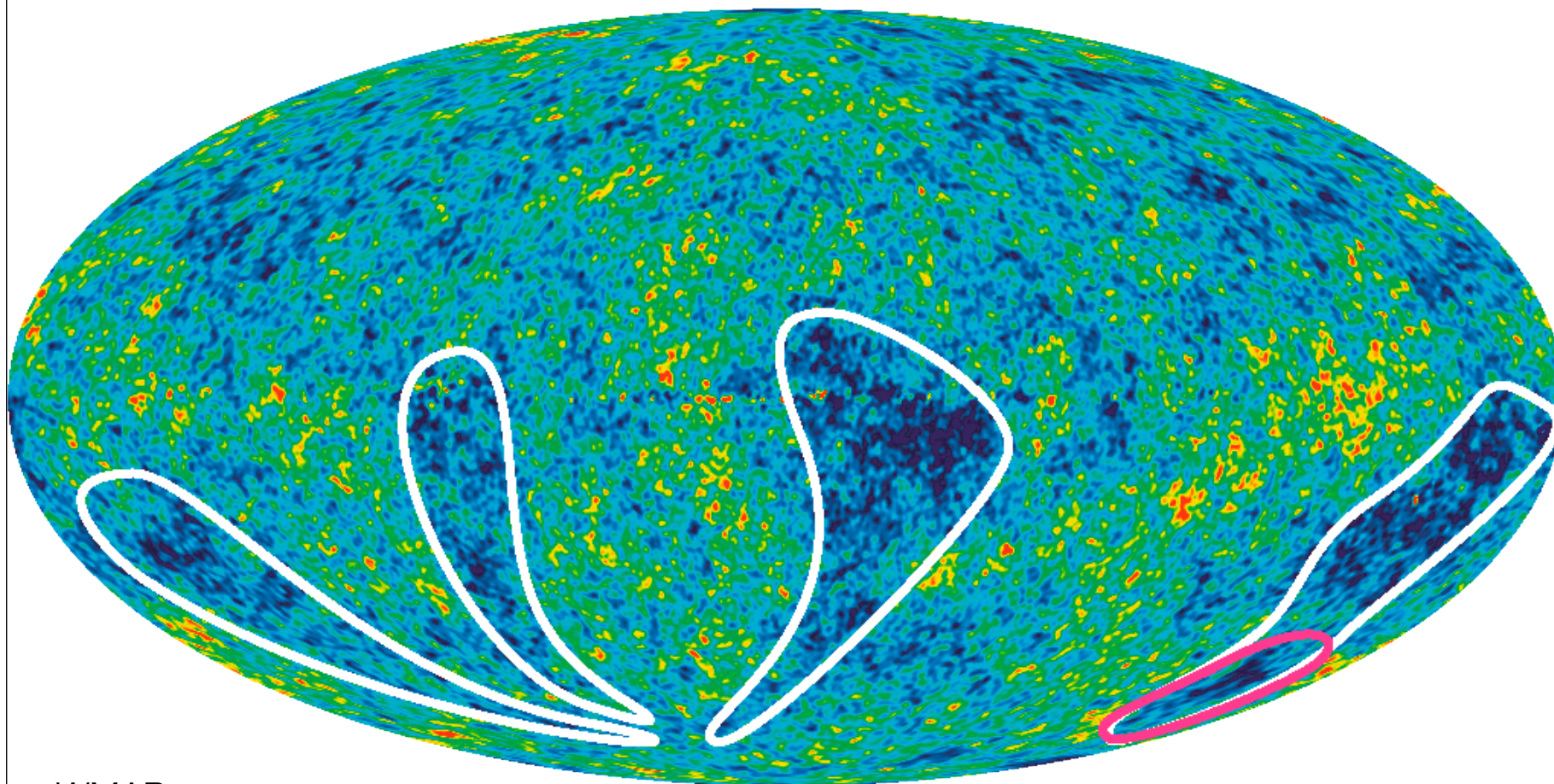
# Large scale CMB anomalies



Copi et al (2010)

# Anomalous cold regions in the southern hemisphere

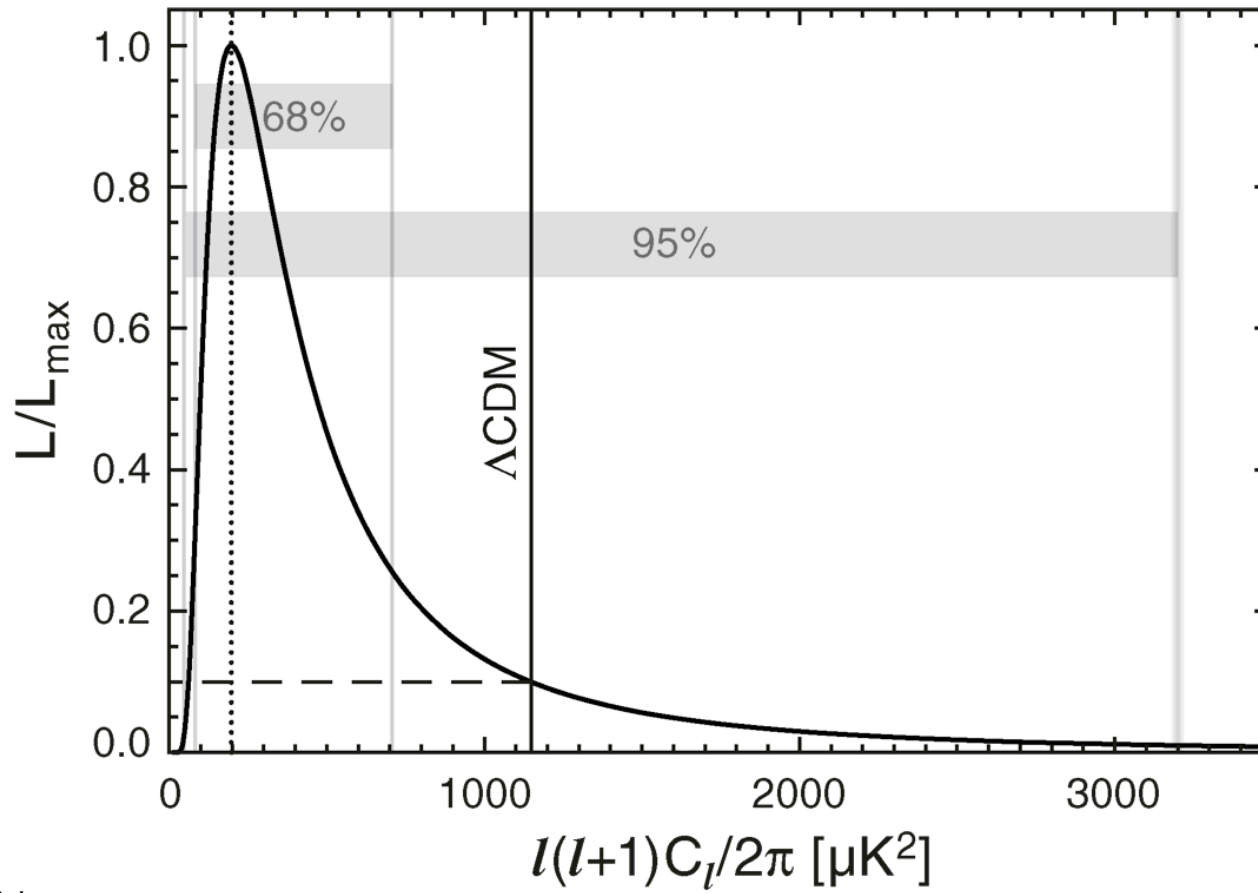
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WMAP team

# Is the quadrupole anomalously low?

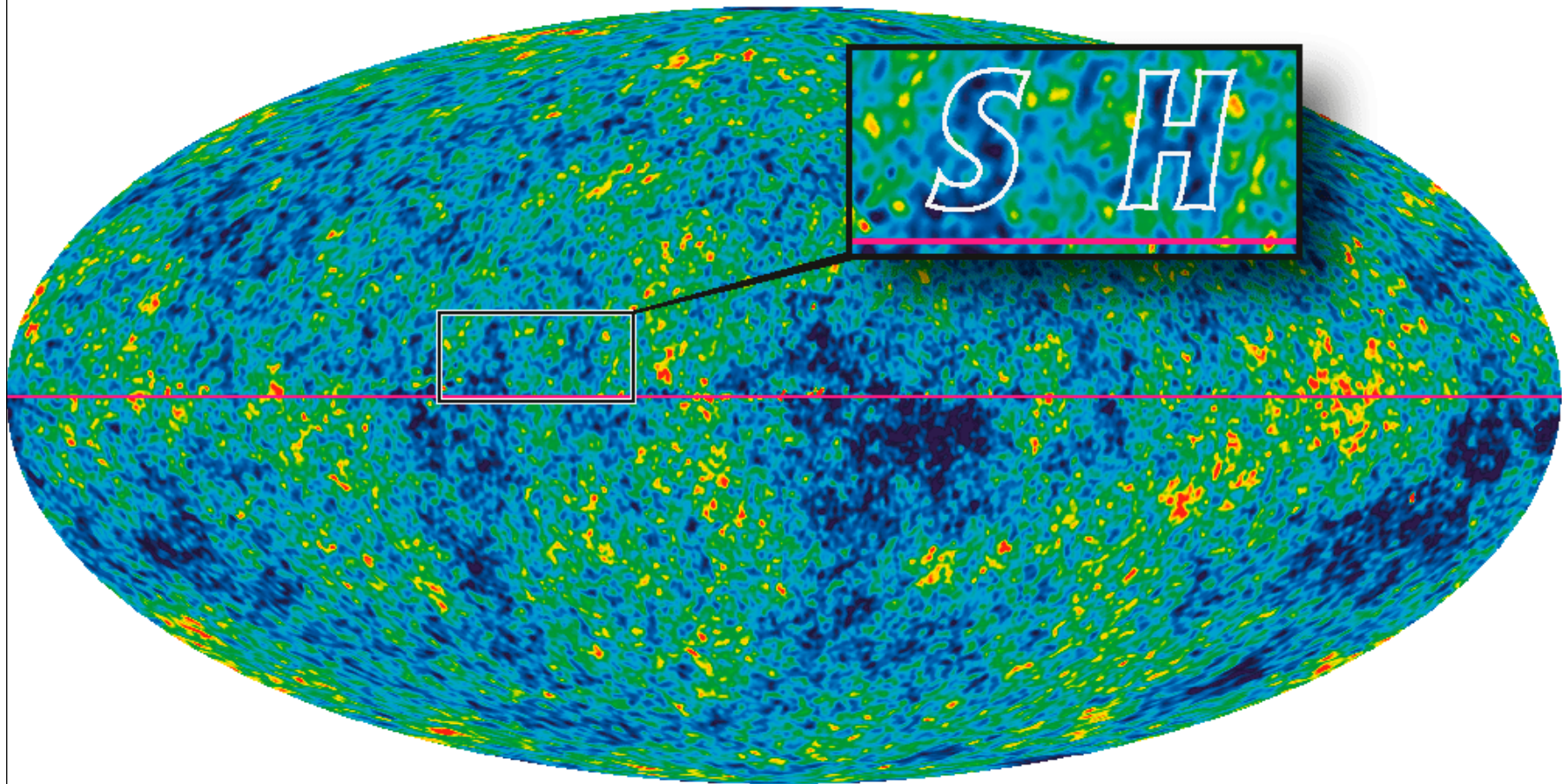
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WMAP team



The “SH” initials of Stephen Hawking are shown in the ILC sky map.



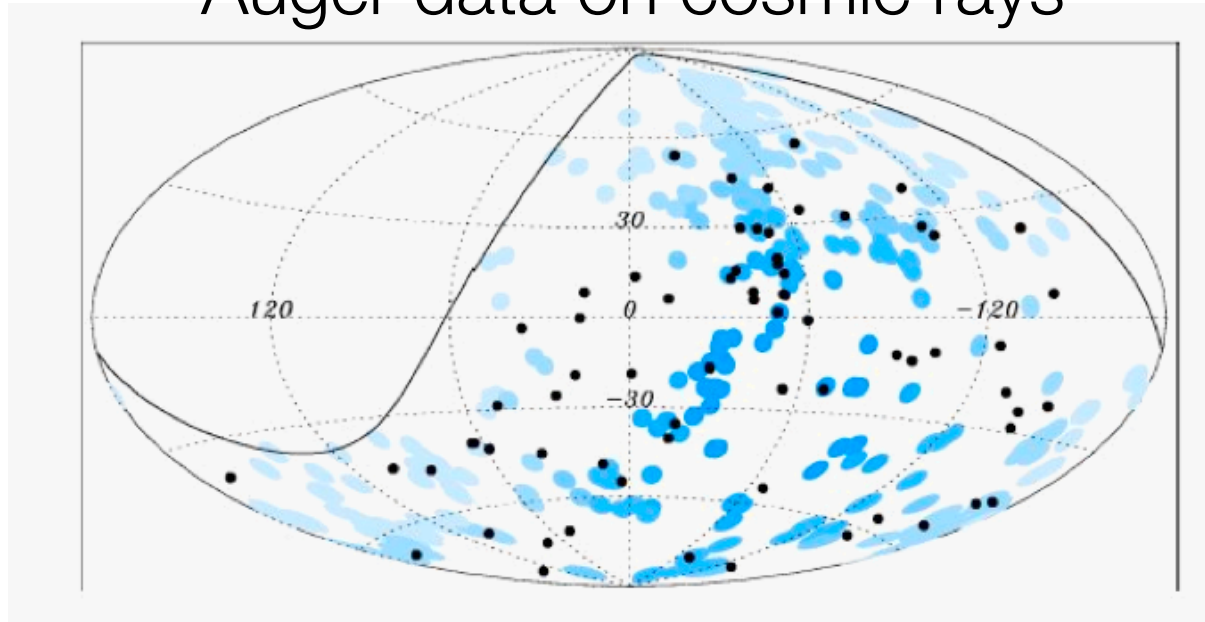
WMAP team

Model selection

# Model selection/hypothesis testing

- For well-defined problems with “easily” identifiable prior information, Bayesian model selection is the tool of choice.

## Auger data on cosmic rays



Tom Loredo's talk:

Odds favouring association of two sources:

| Angular error                    | Odds $O$                 |                         |
|----------------------------------|--------------------------|-------------------------|
|                                  | $\theta_{12} = 26^\circ$ | $\theta_{12} = 0^\circ$ |
| $\sigma_1 = \sigma_2 = 10^\circ$ | $\approx 1.5$            | $\approx 75$            |
| $\sigma_1 = \sigma_2 = 25^\circ$ | $\approx 7$              | $\approx 12$            |

### **Advantages:**

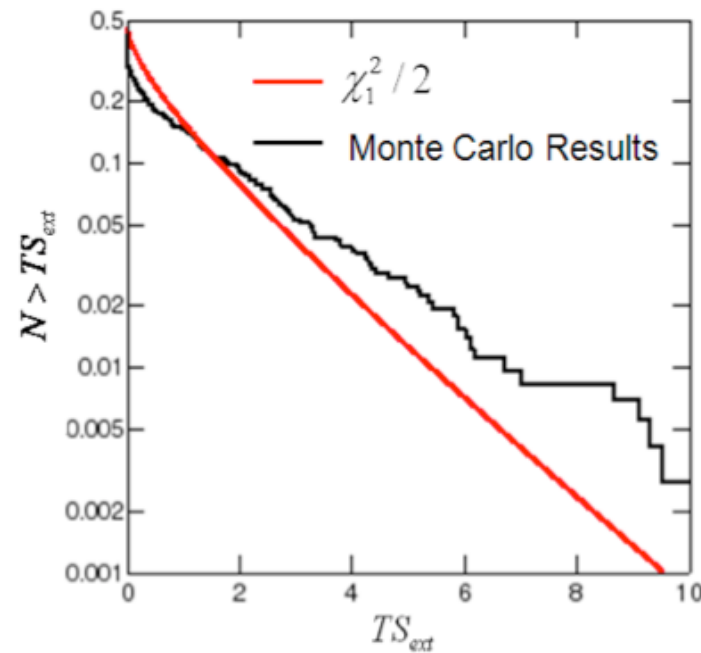
Tuning replaced by averaging

No cuts - the full information is exploited

No fuss with a posteriori statistics

# Object detection

- Detection of extended sources in Fermi data falls foul of Wilk's theorem even for toy cases (Elliott Bloom's talk)
- In real life, all other parameters in the fit must also be learnt from the data (background, spectral shape, source location, source extension, amplitude)

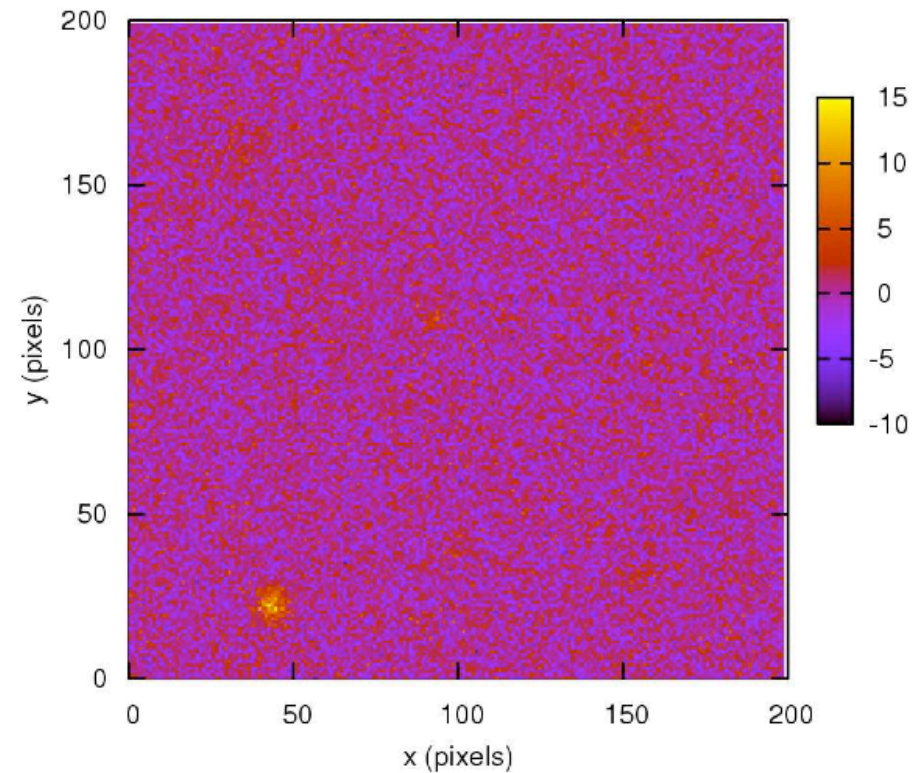


Elliott Bloom's talk

# A “simple” example: how many sources?

Feroz and Hobson  
(2007)

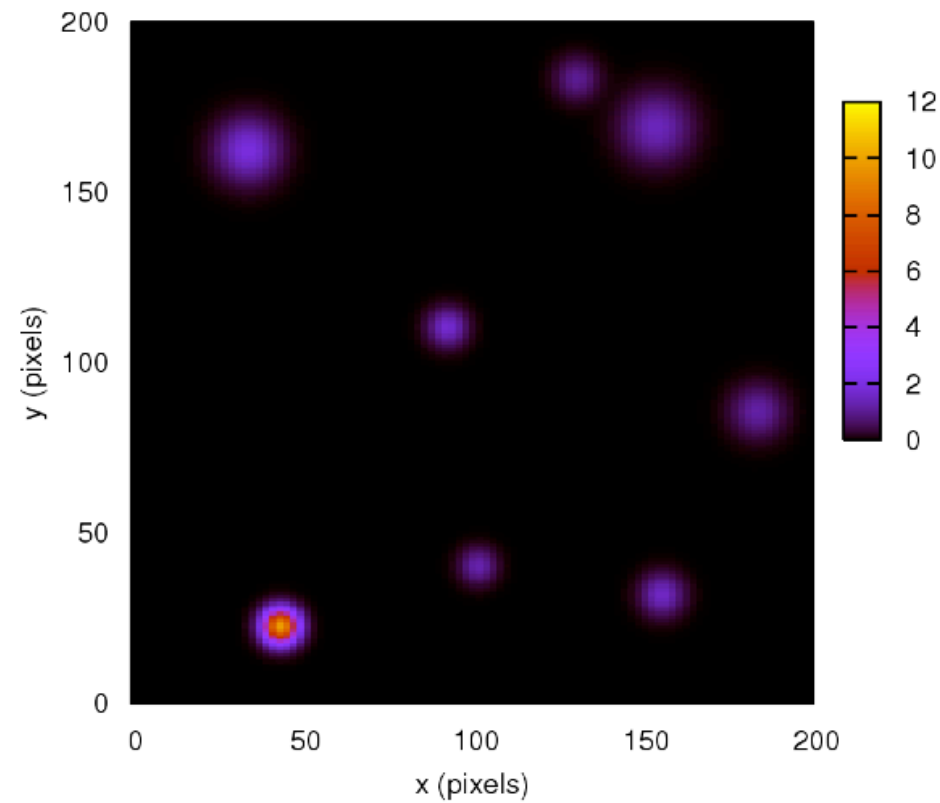
## Signal + Noise



# A “simple” example: how many sources?

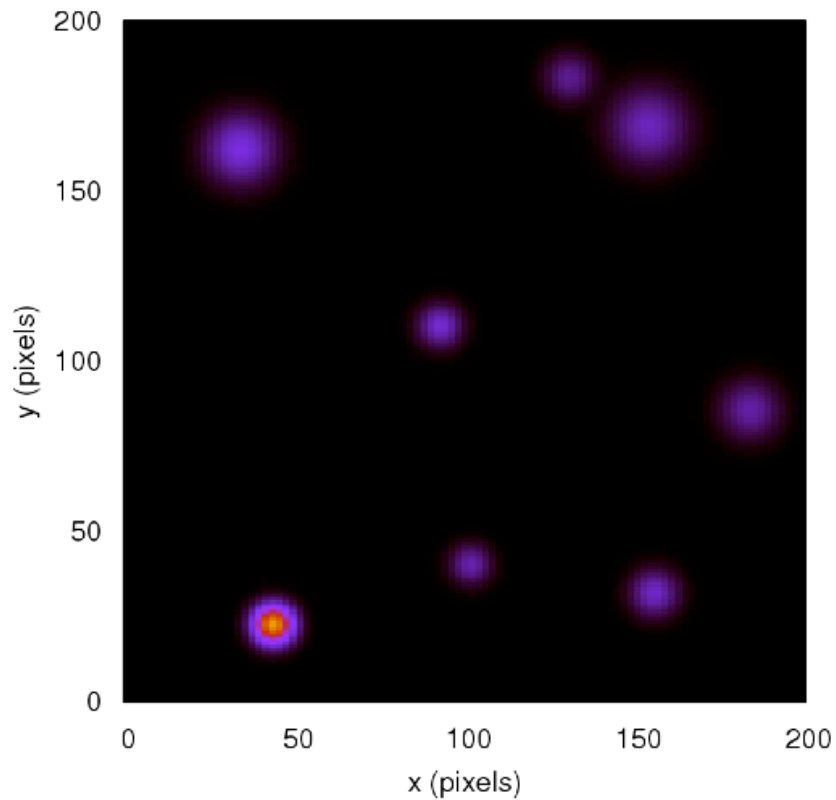
Feroz and Hobson  
(2007)

Signal: 8 sources



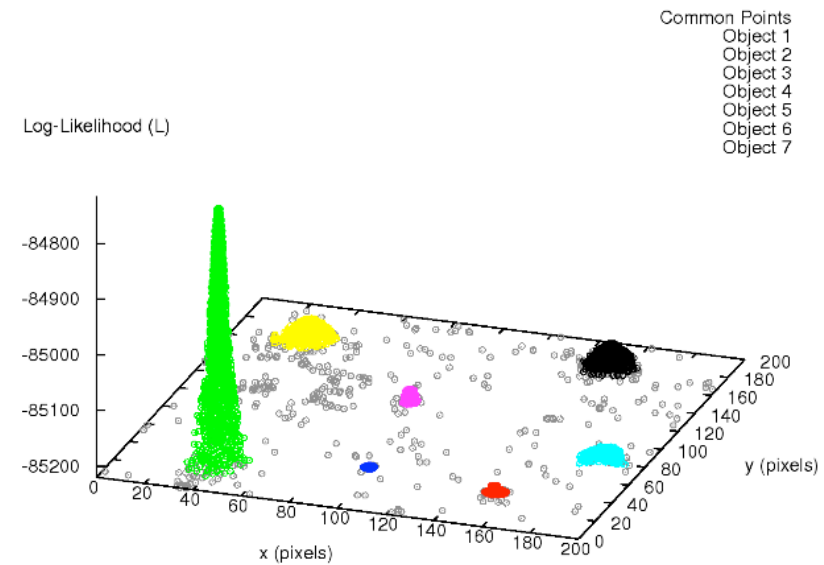
# A “simple” example: how many sources?

Feroz and Hobson  
(2007)



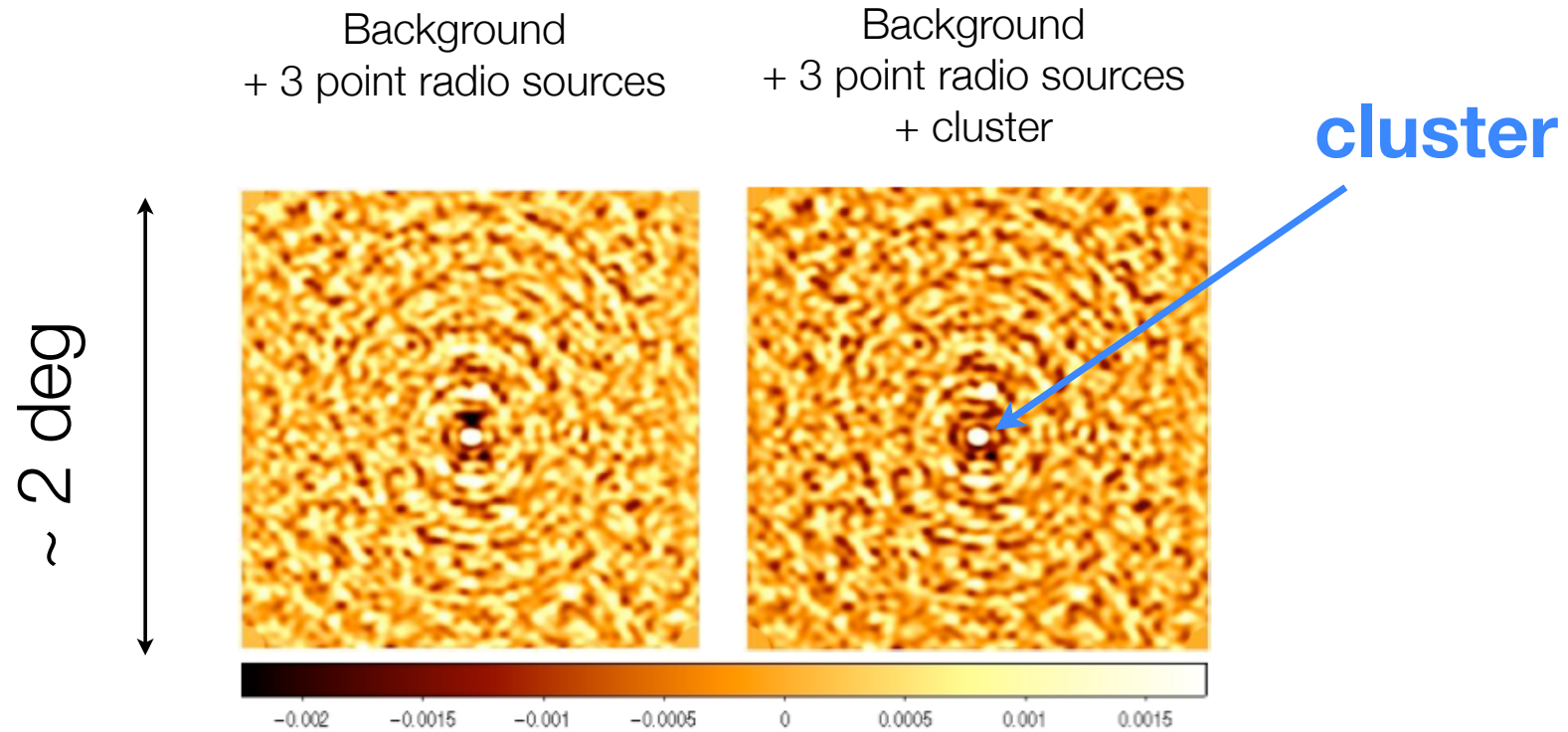
## Bayesian reconstruction

7 out of 8 objects correctly identified.  
Mistake happens because 2 objects very close.





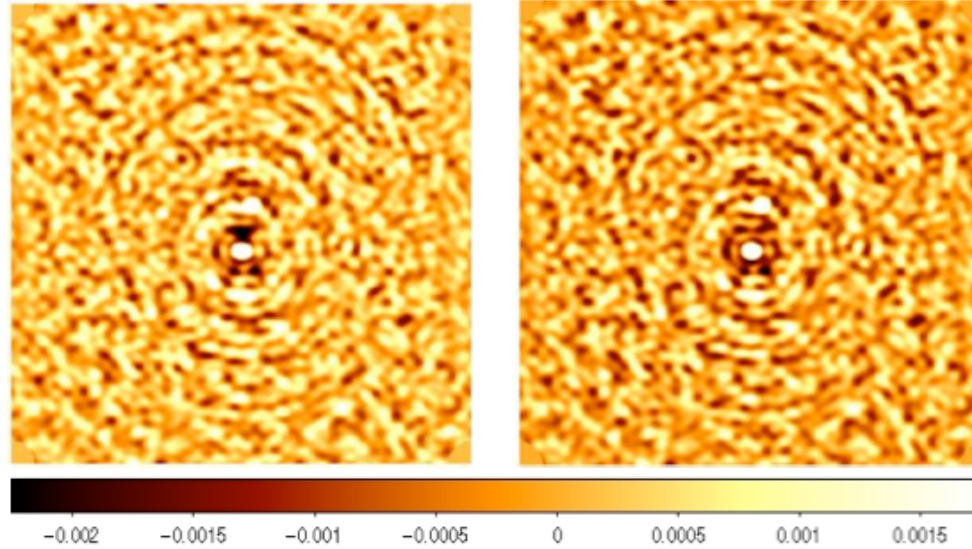
# Cluster detection from Sunyaev-Zeldovich effect in cosmic microwave background maps



Feroz et al 2009

Background  
+ 3 point radio sources

Background  
+ 3 point radio sources  
+ cluster



Bayesian model comparison:

$$R = \frac{P(\text{cluster} \mid \text{data})}{P(\text{no cluster} \mid \text{data})}$$

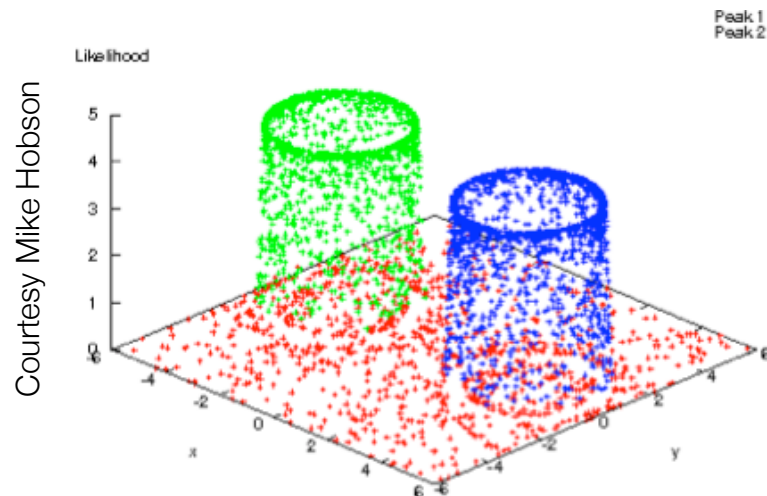
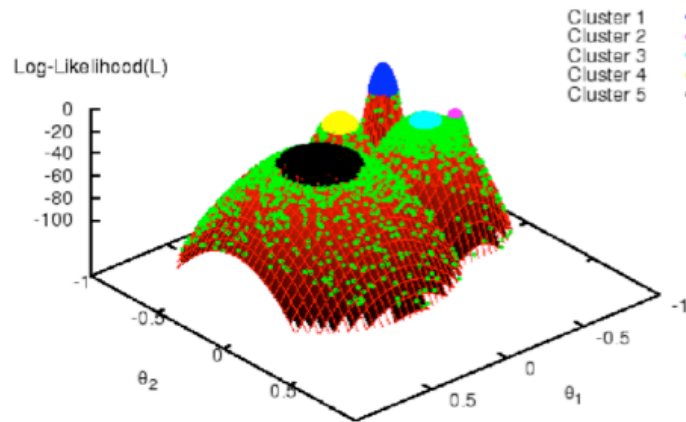
$$R = 0.35 \pm 0.05$$

$$R \sim 10^{33}$$

Cluster parameters also recovered (position, temperature, profile, etc)

# Computation of the evidence with Multinest Imperial College London

Feroz and Hobson  
(2007)



Gaussian mixture model:

True evidence:  $\log(E) = -5.27$

**Multinest:**

Reconstruction:  $\log(E) = -5.33 \pm 0.11$

Likelihood evaluations  $\sim 10^4$

**Thermodynamic integration:**

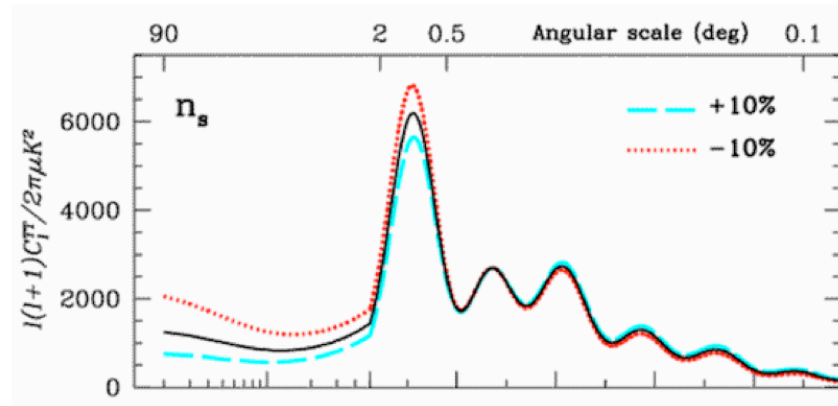
Reconstruction:  $\log(E) = -5.24 \pm 0.12$

Likelihood evaluations  $\sim 10^6$

| D  | $N_{\text{like}}$ | efficiency | likes per dimension |
|----|-------------------|------------|---------------------|
| 2  | 7000              | 70%        | 83                  |
| 5  | 18000             | 51%        | 7                   |
| 10 | 53000             | 34%        | 3                   |
| 20 | 255000            | 15%        | 1.8                 |
| 30 | 753000            | 8%         | 1.6                 |

# Cosmological model selection

- Is the spectrum of primordial fluctuations scale-invariant ( $n = 1$ )?
- Model comparison:  
 **$n = 1$  vs  $n \neq 1$**  (with inflation-motivated prior)
- Results:  
 **$n \neq 1$  favoured with odds of 17:1** (Trotta 2007)  
 **$n \neq 1$  favoured with odds of 15:1** (Kunz, Trotta & Parkinson 2007)  
 **$n \neq 1$  favoured with odds of 7:1** (Parkinson 2007 et al 2006)



# Where Bayesian model selection can go wrong

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- In cosmology/HEP we have many situations with nested models with extra unknown parameters for the fundamental theory.
- Little or nothing is known about the metric to be imposed on such a parameter space
- “The concept of total ignorance about  $\theta$  does not have any precise meaning” (Bob Cousins)
- “ $\theta$  is  $\theta!$ ” (Bob Cousins)

# Where Bayesian model selection can go wrong

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- Occam's razor factor may be arbitrary. HOWEVER: if the range of your prior is arbitrary (by many orders of magnitude) then arguably the physics behind it is not strongly predictive...
- In some cases, the upper bound formalism might be useful (Jim Berger and collaborators)
- In the cosmology community, people often use (blindly) Information Criteria (often with silly answers).

# Information criteria

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- Several information criteria exist for approximate model comparison  
k = number of fitted parameters, N = number of data points,  
-2 ln(L<sub>max</sub>) = best-fit chi-squared

- **Akaike Information Criterium (AIC):**

$$\text{AIC} \equiv -2 \ln \mathcal{L}_{\max} + 2k$$

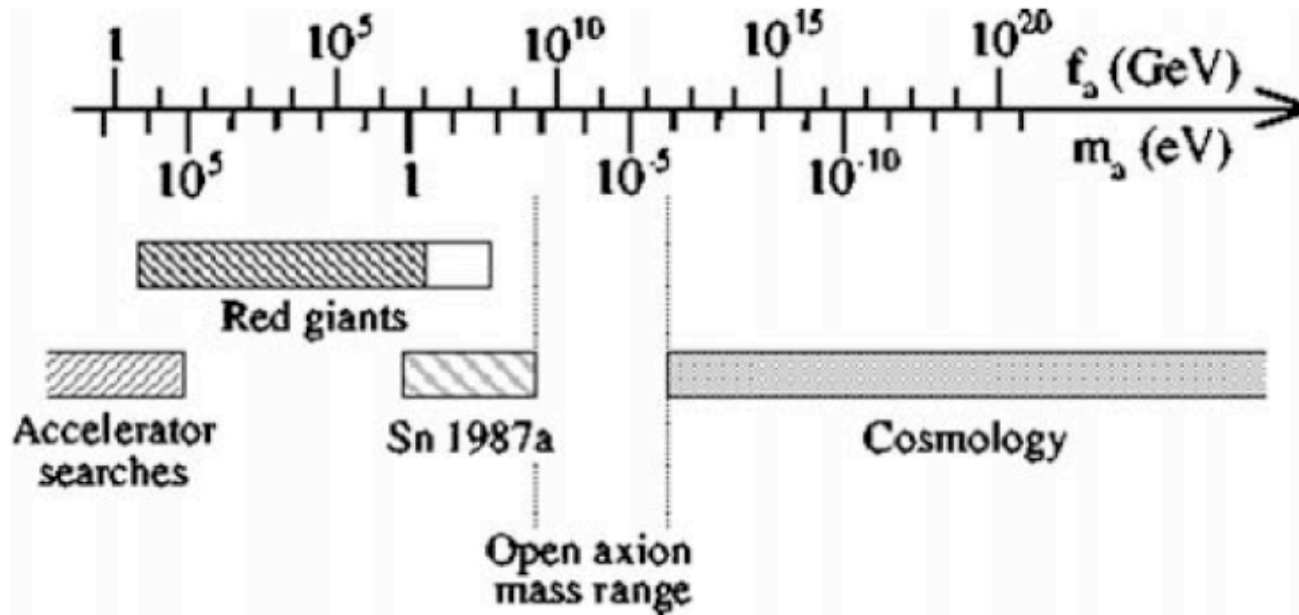
- **Bayesian Information Criterium (BIC):**

$$\text{BIC} \equiv -2 \ln \mathcal{L}_{\max} + k \ln N$$

- **Deviance Information Criterium (DIC):**

$$\text{DIC} \equiv -2 \widehat{D}_{\text{KL}} + 2\mathcal{C}_b.$$

# Axion discovery space



**For most exploratory experiments I can think of, these metrics just don't exist in a relevant way.**

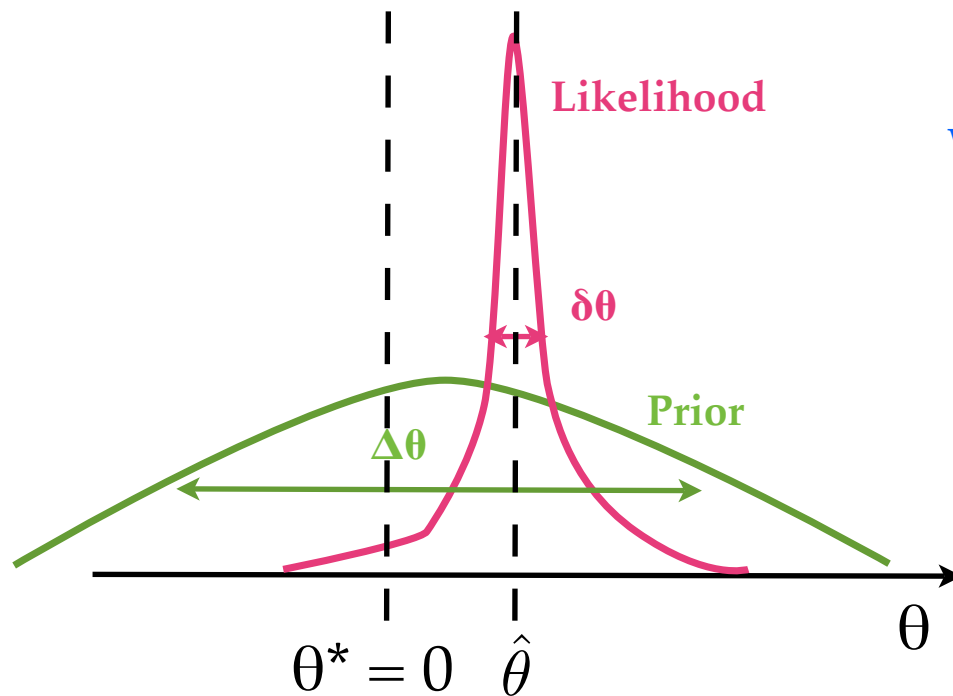
(Bob Cousin's talk)



# Nested models

$M_0: \theta = 0$

$M_1: \theta \neq 0$  with prior  $p(\theta)$



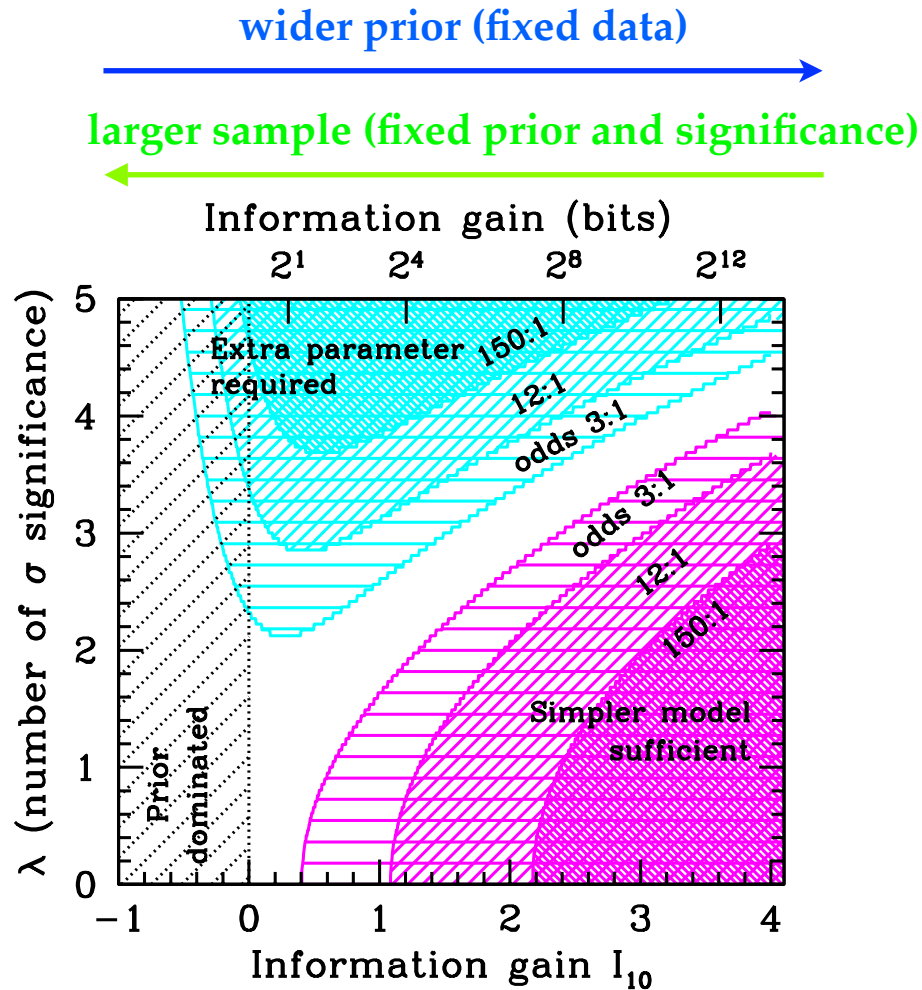
$$\lambda \equiv \frac{\hat{\theta} - \theta^*}{\delta\theta}$$

$$\ln B_{01} \approx \ln \frac{\Delta\theta}{\delta\theta} - \frac{\lambda^2}{2}$$

wasted parameter space  
(favours simpler model)

mismatch of prediction with observed data  
(favours more complex model)

# Model selection for nested models



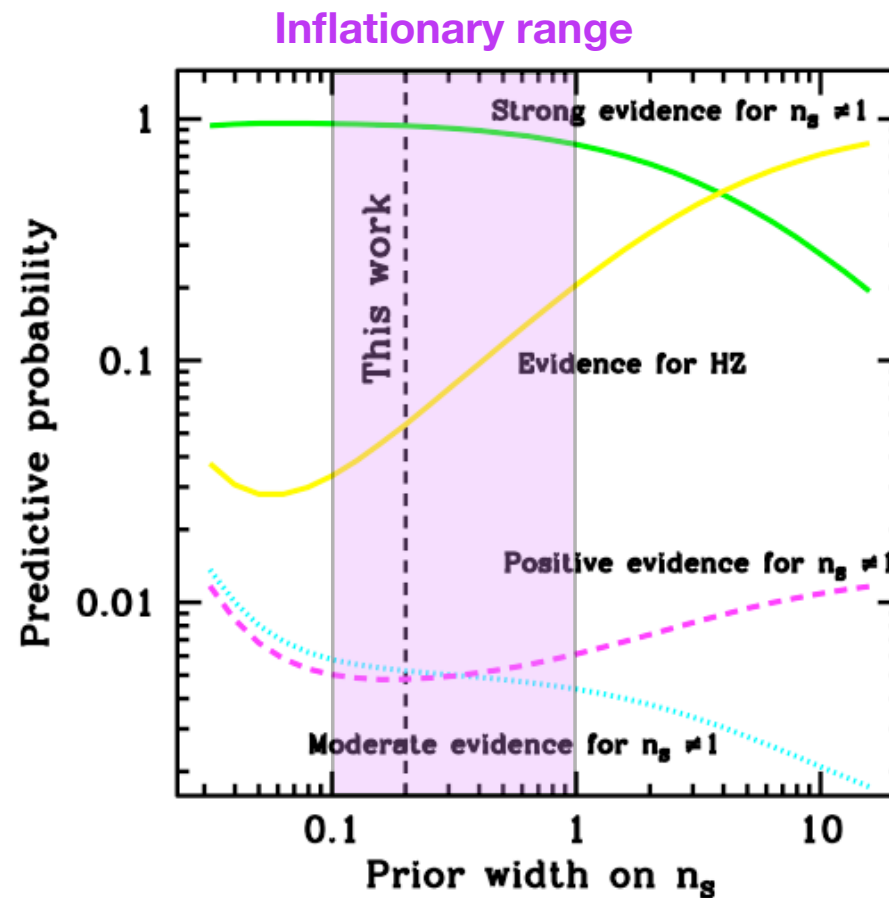
Another look at Lindley's paradox (Bob Cousins' talk)

Jim Berger argued that one should look at the scale of the prior and hope that the result is robust for reasonable choices

$$I_{10} \equiv \log_{10} \frac{\Delta\theta}{\delta\theta}$$

# Example of reasonable sensitivity analysis

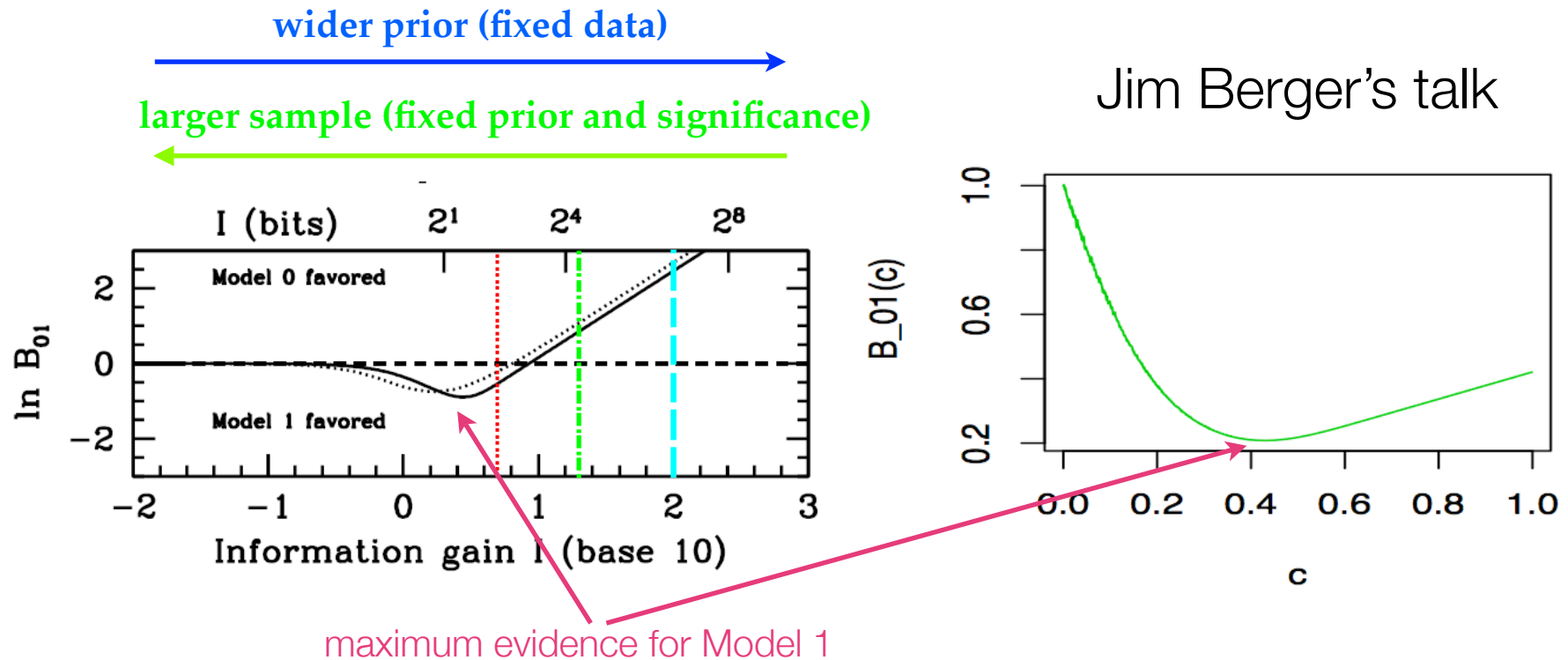
- The answer does not change for physically reasonable changes in the prior width



Trotta (2007)

# “Prior-free” evidence bounds

- What if we do not know how to set the prior? For nested models, we can still choose a prior that will maximise the support for the more complex model:



Jim Berger's talk

# Maximum evidence for a detection

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- **The absolute upper bound:** put all prior mass for the alternative onto the observed maximum likelihood value. Then

$$B < \exp(-\chi^2/2)$$

- **More reasonable class of priors:** symmetric and unimodal around  $\Psi=0$ , then ( $\alpha$  = significance level)

$$B < \frac{-1}{\exp(1)\alpha \ln \alpha}$$

***If the upper bound is small, no other choice of prior will make the extra parameter significant.***

Sellke, Bayarri & Berger, *The American Statistician*, 55, 1 (2001)

# How to interpret the “number of sigma’s”

| <b>p</b> | <b>sigma</b> | <b>Absolute<br/>bound on<br/>lnB (B)</b> | <b>“Reasonable”<br/>bound on lnB<br/>(B)</b> |
|----------|--------------|--|--|
| 0.05     | 2.0          | 2.0<br>(7:1)<br><i>weak</i>              | 0.9<br>(3:1)<br><i>undecided</i>             |
| 0.003    | 3.0          | 4.5<br>(90:1)<br><i>moderate</i>         | 3.0<br>(21:1)<br><i>moderate</i>             |
| 0.0003   | 3.6          | 6.48<br>(650:1)<br><i>strong</i>         | 5.0<br>(150:1)<br><i>strong</i>              |

# A conversion table

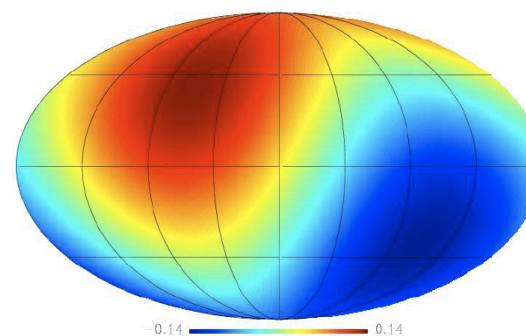
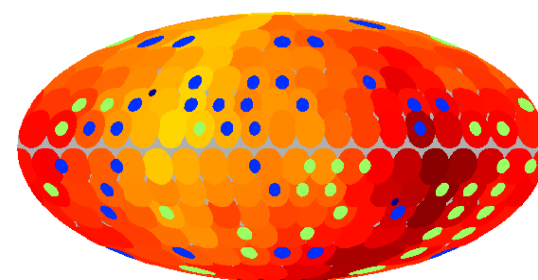
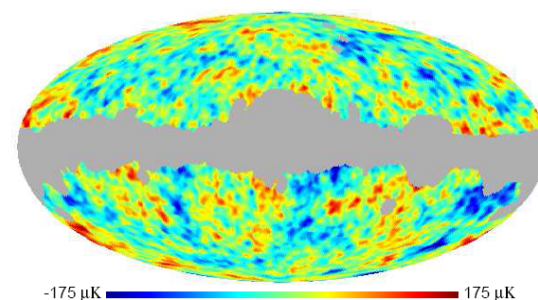
| p-value            | $\bar{B}$ | $\ln \bar{B}$ | sigma | category           |
|--------------------|-----------|---------------|-------|--------------------|
| 0.05               | 2.5       | 0.9           | 2.0   |                    |
| 0.04               | 2.9       | 1.0           | 2.1   | 'weak' at best     |
| 0.01               | 8.0       | 2.1           | 2.6   |                    |
| 0.006              | 12        | 2.5           | 2.7   | 'moderate' at best |
| 0.003              | 21        | 3.0           | 3.0   |                    |
| 0.001              | 53        | 4.0           | 3.3   |                    |
| 0.0003             | 150       | 5.0           | 3.6   | 'strong' at best   |
| $6 \times 10^{-7}$ | 43000     | 11            | 5.0   |                    |

## Rule of thumb:

*a n-sigma result should be interpreted as  
a n-1 sigma result*

# Application: dipole modulation

- Eriksen et al (2004) found hints for a dipolar modulation in WMAP1 ILC map
- Adding a phenomenological dipole pattern **improves the chi-square by 9 units (for 3 extra parameters)**
- Is this significant evidence?
- Not really: **upper bound on B is odds of 9:1**  
The absolute upper bound is about the same (Gordon and Trota 2007)



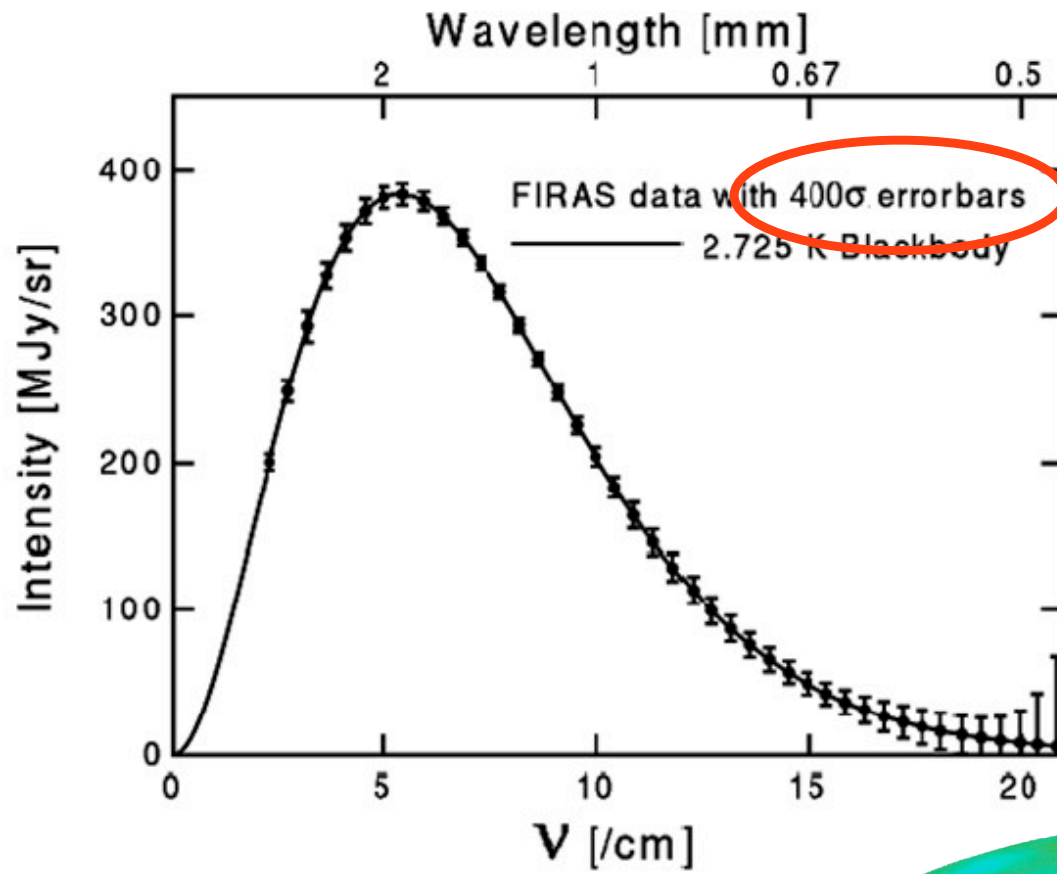


Why 5 sigma?

# Relevance of 5 sigma for cosmology

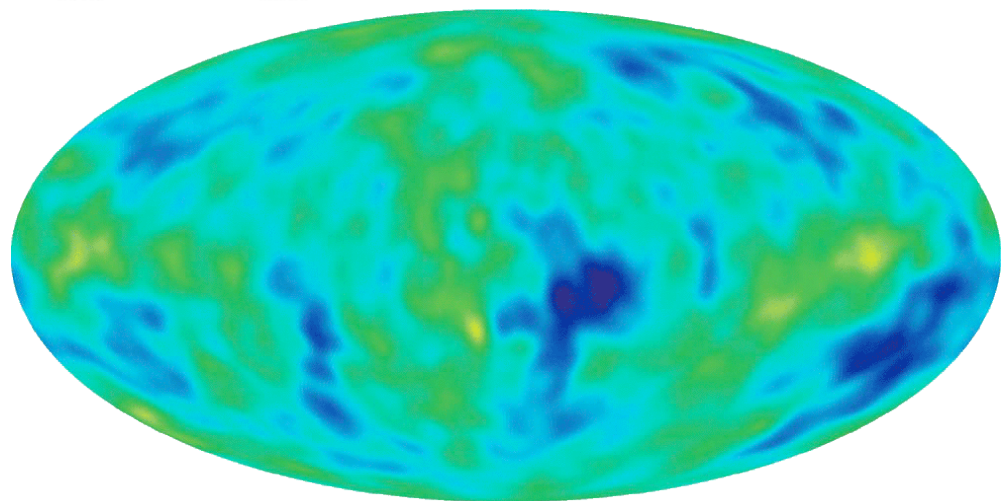
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- I cannot think of many examples where this is relevant in cosmology
- **Discovery of the CMB:** Penzias and Wilson (1965)  $T = 3.5 \pm 1.0$  K (3.5 sigma)  
NOBEL PRIZE 1978
- **Blackbody nature of the CMB:** this was a slam-dunk discovery  
NOBEL PRIZE 2006 (Mather)
- **COBE measurement of anisotropies in the CMB (1994)**  
Quadrupole measurement =  $15.3 +3.8-2.8$   $\mu$ K (~ 5.4 sigma)  
NOBEL PRIZE 2006 (Smoot)



COBE (1994)

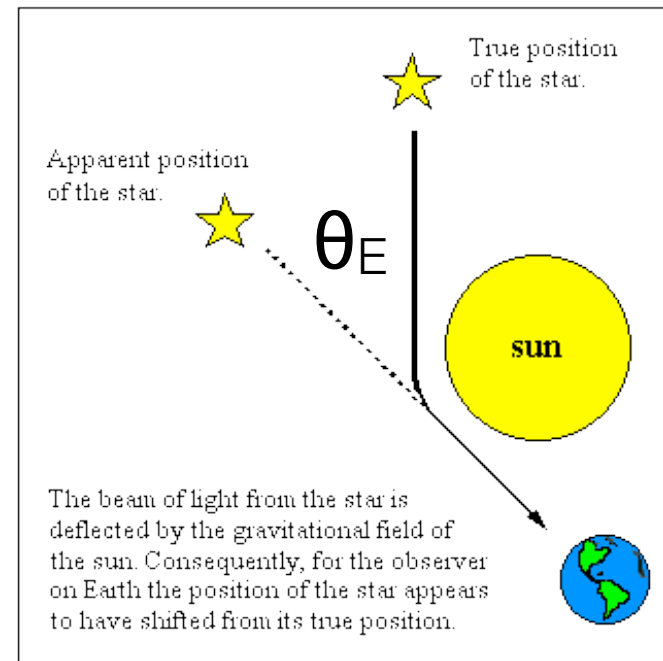
$\Delta T/T \sim 10^{-5}$   
Angular resolution  
 $\sim 10$  deg



# Evidence for Einstein gravity (1919)

- Einstein's theory of General Relativity made a crucial prediction: the deflection angle around the Sun should be twice what predicted by Newton

$$\theta_E = 2 \theta_N = 1.75''$$



# Evidence for Einstein gravity (1919)

- Measurements were performed during the solar eclipse of May 29th 1919:

**Eddington:**  $\theta = 1.61 \pm 0.40$  arcsec (based on 5 stars)

**Crommelin:**  $\theta = 1.98 \pm 0.16$  arcsec (based on 7 stars)

|                                       | <b>Einstein</b>          | <b>Newton</b>           |
|---------------------------------------|--------------------------|-------------------------|
| Hypothesis                            | $\theta = 1.75$ arcsec   | $\theta = 0.875$ arcsec |
| p-value from Eddington's data         | 0.72                     | 0.06                    |
| Posterior odds for Einstein vs Newton | ~ 5 to 1 (weak evidence) |                         |

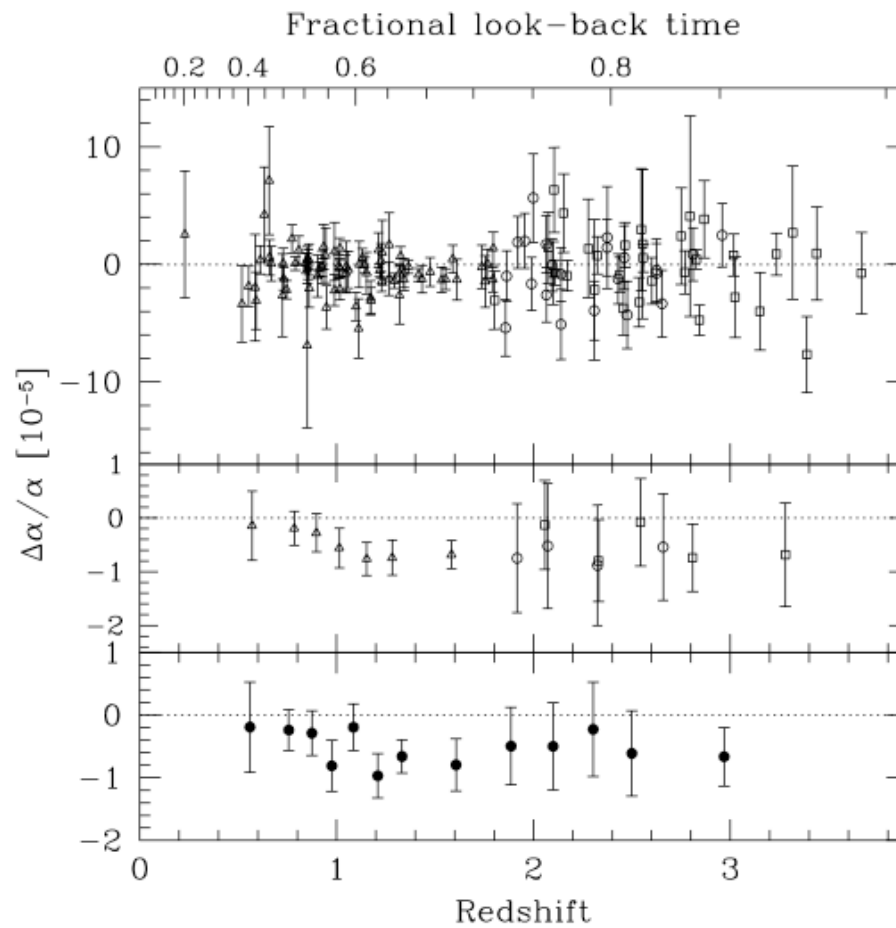
# Evidence for varying $\alpha$ ?

- For several years, Webb and collaborators have claimed a  $\sim 5$  sigma evidence for a time variation in the fine structure constant from analysis of QSO absorption spectra.

random error in 22 high- $z$  systems characterized by transitions with a large dynamic range in apparent optical depth. Increasing the statistical errors on  $\Delta\alpha/\alpha$  for these systems gives our fiducial result, a weighted mean  $\Delta\alpha/\alpha = (-0.543 \pm 0.116) \times 10^{-5}$ , representing  $4.7 \sigma$  evidence for a varying  $\alpha$ . Assuming that  $\Delta\alpha/\alpha = 0$  at  $z_{\text{abs}} = 0$ , the data marginally prefer a linear increase in  $\alpha$  with time rather than a constant offset from the laboratory value:  $\dot{\alpha}/\alpha = (6.40 \pm 1.35) \times 10^{-16} \text{ yr}^{-1}$ . The two-point correlation function for  $\alpha$  is consistent with zero over  $0.2 - 13 \text{ Gpc}$  comoving scales and the angular distribution of  $\Delta\alpha/\alpha$  shows no

Murphy et al (2003)

# Evidence for varying $\alpha$ ?



Murphy et al (2003)

# My 2 pennies

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- Standards of evidence in physics are not absolute: we have a Bayesian prior in the back of our minds when assessing strength of evidence (systematics, plausibility, scientific experience, appeal of the model, theoretical framework, simplicity, how the model fits within the bigger picture, elegance, etc).
- How those factors could be summarized in  $P(M)$  is difficult to imagine.
- Jim Berger argued that priors should be “defendible”, no matter how you got there.
- “Inside every Frequentist there is a Bayesian struggling to get out” (Lindley).
- Bayesian model selection works best in cases where relevant prior information can be objectively specified (e.g., object detection example).



THANK YOU!