

Efficient Algorithms for Semiclassical Quantum Dynamics

David Sattlegger

joint work with Caroline Lasser and Manfred Liebmann

Banff, January 27, 2016

Time-dependent semiclassical Schrödinger equation

Goal: Find solution to

$$i \varepsilon \frac{d}{dt} \psi(t) = -\frac{1}{2} \varepsilon^2 \Delta \psi(t) + V \psi(t)$$

with $\psi(0) = \psi_0 \in L^2(\mathbb{R}^d)$, $\|\psi_0\|_{L^2} = 1$ and $\varepsilon := \sqrt{\frac{m_e}{M_n}} \ll 1$.

Self-adjoint $H := -\frac{1}{2} \varepsilon^2 \Delta + V$ generates U_t ,

$$\psi(t) = U_t \psi_0 = e^{-\frac{i}{\varepsilon} t H} \psi_0$$

for all $t \in \mathbb{R}$.

Time-dependent semiclassical Schrödinger equation

Goal: Find solution to

$$i \varepsilon \frac{d}{dt} \psi(t) = -\frac{1}{2} \varepsilon^2 \Delta \psi(t) + V \psi(t)$$

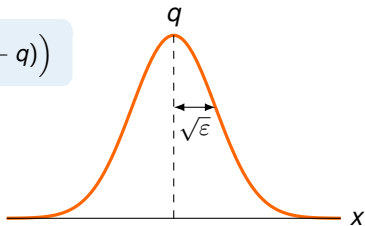
with $\psi(0) = \psi_0 \in L^2(\mathbb{R}^d)$, $\|\psi_0\|_{L^2} = 1$ and $\varepsilon := \sqrt{\frac{m_e}{M_n}} \ll 1$.

- ▶ high dimensional system
- ▶ small ε produces highly oscillatory solutions
- ▶ ...

Gaussian wave packets

$$g_z^\varepsilon(x) := (\pi\varepsilon)^{-\frac{d}{4}} \exp\left(-\frac{1}{2\varepsilon} |x - q|^2 + \frac{i}{\varepsilon} p \cdot (x - q)\right)$$

$$z = \begin{pmatrix} q \\ p \end{pmatrix} \in \mathbb{R}^{2d}$$



$$\psi = (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{2d}} g_z^\varepsilon \langle g_z^\varepsilon, \psi \rangle dz \quad \text{for all } \psi \in L^2(\mathbb{R}^d)$$

$$U_t \psi = (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{2d}} (U_t g_z^\varepsilon) \langle g_z^\varepsilon, \psi \rangle dz$$

Herman–Kluk Propagator

Definition (Herman & Kluk, 1984)

$$\mathcal{I}_t \psi := (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{2d}} \left(u(t, z) e^{\frac{i}{\varepsilon} S(t, z)} g_{\Phi^t(z)}^\varepsilon \right) \langle g_z^\varepsilon, \psi \rangle dz$$

depends on

- ▶ classical flow $\Phi^t = (X^t, \Xi^t) : \mathbb{R}^{2d} \rightarrow \mathbb{R}^{2d}$,
- ▶ classical action integral $S(t, z) := \int_0^t \left(\frac{d}{d\tau} X^\tau(z) \cdot \Xi^\tau(z) - h(\Phi^\tau(z)) \right) d\tau$,
- ▶ Herman–Kluk factor $u(t, z)$

$$u(t, z) := \sqrt{2^{-d} \det(\partial_q X^t(z) + \partial_p \Xi^t(z) + i (\partial_q \Xi^t(z) - \partial_p X^t(z)))}.$$

Herman–Kluk Propagator

$$\begin{aligned}U_t \psi &= (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{2d}} (U_t g_z^\varepsilon) \langle g_z^\varepsilon, \psi \rangle dz \\ \mathcal{I}_t \psi &= (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{2d}} \left(u(t, z) e^{\frac{i}{\varepsilon} S(t, z)} g_{\Phi_t(z)}^\varepsilon \right) \langle g_z^\varepsilon, \psi \rangle dz.\end{aligned}$$

Theorem (Swart & Rouse, 2009)

\mathcal{I}_t is a bounded operator on $L^2(\mathbb{R}^d)$. For every $T > 0$, there exists $C > 0$ such that for all $\varepsilon > 0$

$$\sup_{t \in [0, T]} \|\mathcal{I}_t - U_t\| \leq C\varepsilon.$$

General strategy

1. High-dimensional quadrature

$$\mathcal{I}_t \psi = \int_{\mathbb{R}^{2d}} f(z) \, dz \approx \frac{1}{M} \sum_{m=1}^M f(z_m).$$

2. Calculate $\Phi^t(z_m)$, $S(t, z_m)$, and $u(t, z_m)$ by solving a system of ODEs with a symplectic integrator.

Quadrature

Factorization

$$(2\pi\varepsilon)^{-d} \langle g_z^\varepsilon, \psi_0 \rangle =: r_0(z) \cdot \mu_0(z), \quad z \in \mathbb{R}^{2d}$$

where $\mu_0 : \mathbb{R}^{2d} \rightarrow \mathbb{R}$ is a probability distribution.

Rewrite \mathcal{I}_t as weighted integral

$$\mathcal{I}_t \psi_0 = \int_{\mathbb{R}^{2d}} r_0(z) u(t, z) e^{\frac{i}{\varepsilon} S(t, z)} g_{\Phi_t(z)}^\varepsilon d\mu_0(z).$$

Quadrature

Define $\psi_M(t) \in L^2(\mathbb{R}^d)$ by

$$\psi_M(t) := \frac{1}{M} \sum_{m=1}^M r_0(z_m) u(t, z_m) e^{\frac{i}{\epsilon} S(t, z_m)} \mathbf{g}_{\Phi^t(z_m)}^\epsilon$$

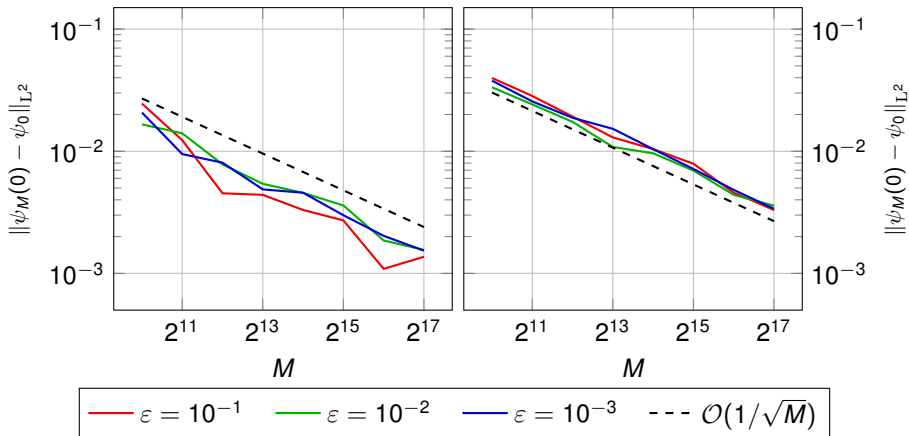
with $z_1, \dots, z_M \in \mathbb{R}^{2d}$ chosen as

- ▶ independent samples of μ_0 (Monte Carlo quadrature),
or
- ▶ low μ_0 -discrepancy points (quasi-Monte Carlo quadrature).

Monte Carlo quadrature

$d = 1$

$d = 2$



Time discretization

$\Phi^t(z_m)$ is solution to

$$\dot{z}(t) = J \nabla h(z(t)), \quad z(0) = z_m. \quad (\mathbb{R}^{2d})$$

$S(t, z_m)$ is solution to

$$\dot{S}(t, z_m) = \frac{1}{2} |\Xi^t(z_m)|^2 - V(X^t(z_m)), \quad S(0, z_m) = 0. \quad (\mathbb{R})$$

$$u(t, z_m) = \sqrt{2^{-d} \det(\partial_q X^t(z_m) + \partial_p \Xi^t(z_m) + i(\partial_q \Xi^t(z_m) - \partial_p X^t(z_m)))}$$

is calculated from

$$\frac{d}{dt} D\Phi^t(z_m) = J \cdot \nabla^2 h(\Phi^t(z_m)) \cdot D\Phi^t(z_m). \quad (\mathbb{R}^{2d \times 2d})$$

Time discretization

Theorem

Calculate $\tilde{\Phi}^t$, \tilde{S} and \tilde{u} with symplectic method of order $\gamma \in \mathbb{N}$. Set

$$\tilde{\mathcal{I}}_t \psi := (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{2d}} \tilde{u}(t, z) e^{\frac{i}{\varepsilon} \tilde{S}(t, z)} g_{\tilde{\Phi}^t(z)}^\varepsilon \langle g_z^\varepsilon, \psi \rangle dz.$$

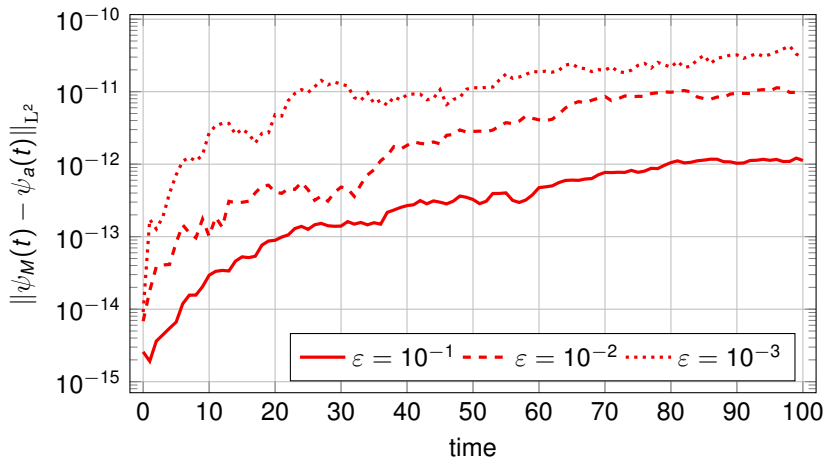
For all $t > 0$ there exist $c_0, c_1, c_2 > 0$ such that for all $\varepsilon > 0$

$$\|\tilde{\mathcal{I}}_t - U_t\| \leq \varepsilon c_0 + \frac{\tau^\gamma}{\varepsilon} c_1 + \tau^\gamma c_2 + \mathcal{O}(\tau^{\gamma+2})$$

where $\tau > 0$ is the time step size.

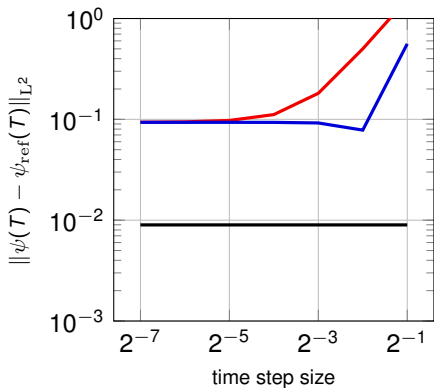
Time evolution of error

for harmonic potential

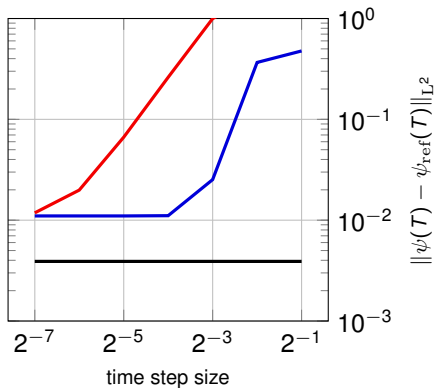


Dependence on time step size

$$\varepsilon = 10^{-1}, M = 2^{13}$$



$$\varepsilon = 10^{-2}, M = 2^{15}$$



— order $\gamma = 2$ at $T = 20$ — order $\gamma = 4$ at $T = 20$ — initial sampling error

Efficient implementation

Idea: Solve

$$\dot{\mathcal{Z}}(t, z_m) = F(\mathcal{Z}(t, z_m)) \in \mathbb{R}^{4d^2+2d+2}$$

independently for all z_m

- ▶ (MPI and OpenMP) parallelisation + (SSE/AVX) vectorisation
- ▶ Automated C++ code generation by computer algebra system

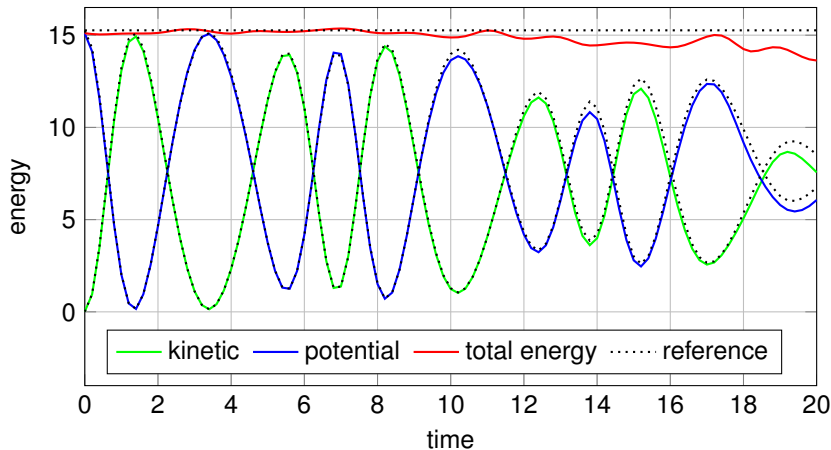
Example: Henon–Heiles potential for $d = 6$

$$V(x) = \sum_{k=1}^6 \frac{1}{2} x_k^2 + \sigma \sum_{k=1}^5 \left(x_k x_{k+1}^2 - \frac{1}{3} x_k^3 \right) + \frac{\sigma^2}{16} \sum_{k=1}^5 (x_k^2 + x_{k+1}^2)^2$$

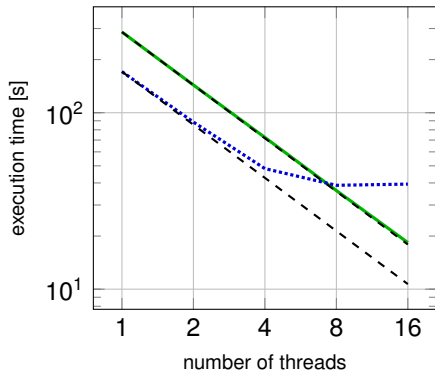
- ▶ $M = 2^{22}$ sampling points resulting in 331,350,016 scalar ODEs
- ▶ 200 time steps

Example

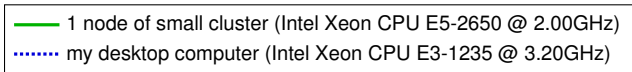
Henon–Heiles potential for $d = 6$, $\varepsilon = 10^{-2}$, $\tau = 0.1$, $\sigma = \frac{1}{\sqrt{80}}$, $\psi_0 = g_{(2,\dots,2,0,\dots,0)}^\varepsilon$



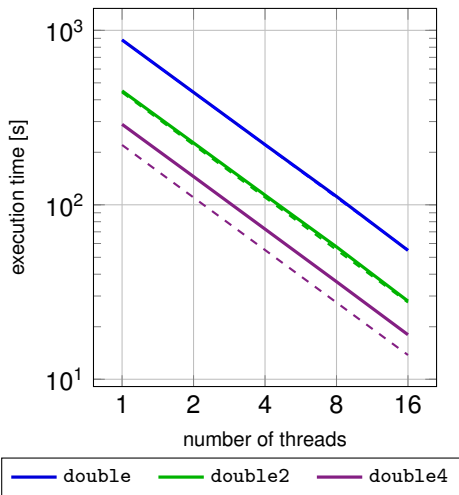
Parallelization



Intel Xeon CPU E5-2650 @ 2.00GHz		
threads	runtime	speedup
1	286.697 s	
2	143.569 s	1.997
4	72.406 s	3.960
8	36.217 s	7.916
16	18.362 s	15.614



Parallelization + Vectorization



threads	double (64 bit)	speedup
1	881.14 s	
2	441.31 s	2.00
4	221.23 s	3.98
8	111.65 s	7.89
16	54.78 s	16.09

threads	double2 (128 bit)	speedup
1	449.78 s	1.96
2	226.34 s	3.89
4	113.61 s	7.76
8	57.38 s	15.36
16	28.05 s	31.41

threads	double4 (256 bit)	speedup
1	289.34 s	3.05
2	145.16 s	6.07
4	72.58 s	12.14
8	36.25 s	24.31
16	18.00 s	48.95

Expectation values

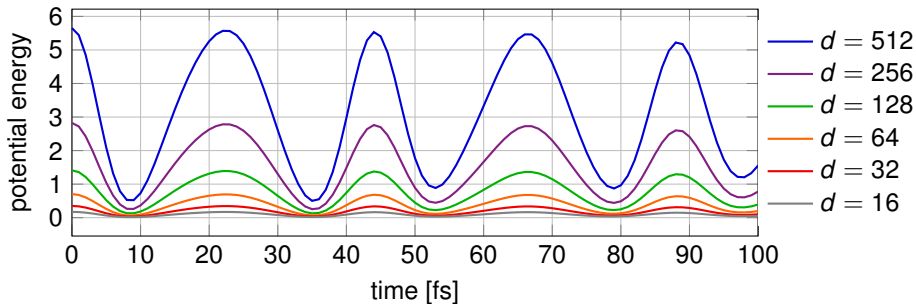
Egorov's Theorem / LSC - IVR

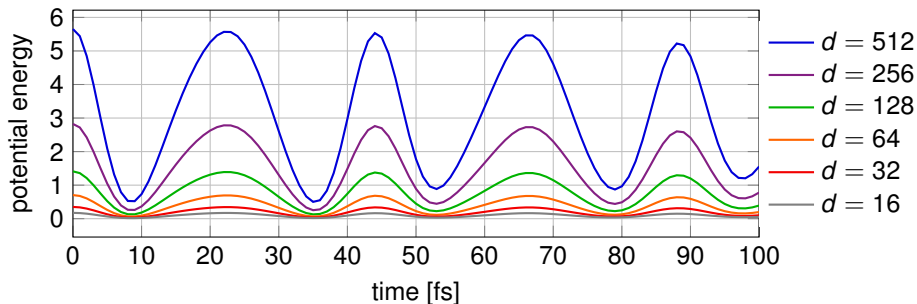
$$\langle \psi(t) | A | \psi(t) \rangle - \int_{\mathbb{R}^{2d}} (a \circ \Phi^t)(z) W(\psi_0)(z) dz = \mathcal{O}(\varepsilon^2)$$

Example: Henon–Heiles potential

$$V(x) = \frac{1}{2} \sum_{k=1}^d x_k^2 + 1.8436 \sum_{k=1}^{d-1} \left(x_k x_{k+1}^2 - \frac{1}{3} x_k^3 \right) + 0.4 \sum_{k=1}^{d-1} (x_k^2 + x_{k+1}^2)^2$$

with $\varepsilon = 0.0037$.





dim	2 × Intel Xeon E5-2650	1 × Nvidia Tesla K20	speedup
16	4.260 s	0.641 s	6.646
32	10.386 s	1.180 s	8.802
64	24.793 s	4.412 s	5.619
128	109.393 s	29.266 s	3.738
256	219.183 s	71.980 s	3.045
512	454.614 s	155.915 s	2.916

Thank you for your attention!