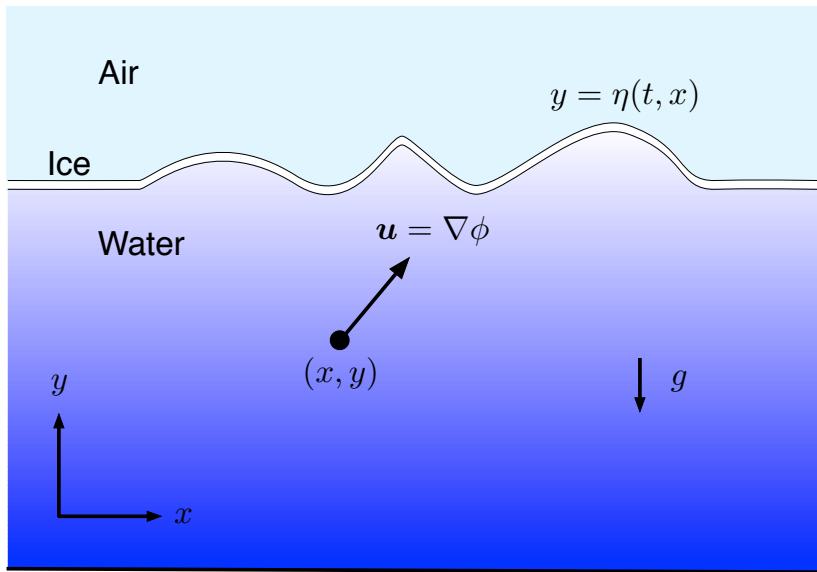


Variational existence and stability theory for hydroelastic solitary waves

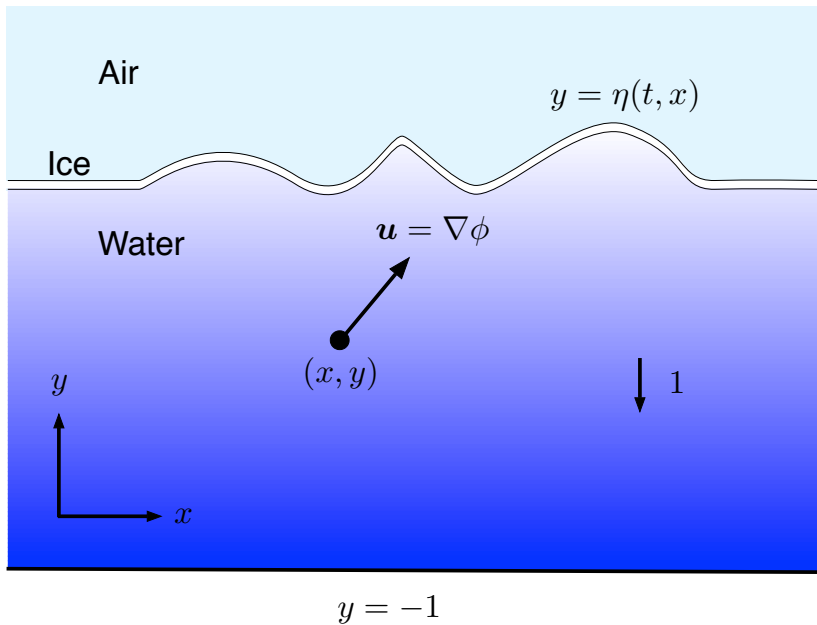
Erik Wahlén, Lund University

Joint with Mark Groves and Benedikt Hoyer,
Saarland University

Banff, November 3, 2016



$$y = -d$$



Equations of motion

(Plotnikov & Toland '11, Blyth, Parau & Vanden-Broeck '11)

$$\Delta\phi = 0 \quad -1 < y < \eta$$

$$\phi_y = 0 \quad y = -1$$

$$\eta_t + \phi_x \eta_x - \phi_y = 0 \quad y = \eta$$

$$\phi_t + \frac{1}{2}|\nabla\phi|^2 + \eta + \gamma H(\eta) = 0 \quad y = \eta$$

$$H(\eta) = \kappa_{ss} + \frac{1}{2}\kappa^3, \quad \gamma = \frac{\mathcal{D}}{\rho g d^4}$$

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$$H(\eta) = \frac{1}{(1 + \eta_x^2)^{\frac{1}{2}}} \left[\frac{1}{(1 + \eta_x^2)^{\frac{1}{2}}} \left(\frac{\eta_{xx}}{(1 + \eta_x^2)^{\frac{3}{2}}} \right) \right]_x \Big|_x + \frac{1}{2} \left(\frac{\eta_{xx}}{(1 + \eta_x^2)^{\frac{3}{2}}} \right)^3$$

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Travelling (solitary) wave: $\eta = \eta(x - ct)$, $\phi = \phi(x - ct, y)$

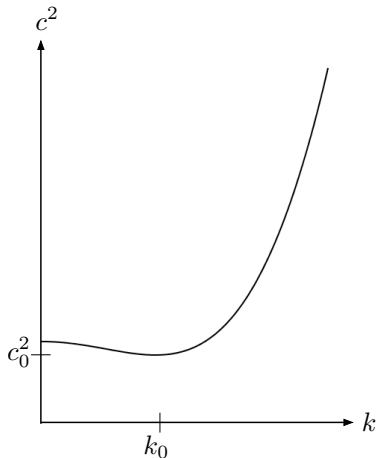
Other models

- ▶ Euler-Bernoulli: $H(\eta) = \eta_{xxxx}$
- ▶ Kirchhoff-Love: $H(\eta) = \partial_x^2 \kappa$

Heuristics

$$\eta(x - ct) \sim e^{ik(x-ct)}, \quad c^2 = (1 + \gamma k^4) \frac{\tanh k}{k}$$

Dispersion relation:



For $c = c_0 - \varepsilon^2$, one can derive the NLS equation

$$v_{xx} - v \pm |v|^2 v = 0$$

where

$$\eta(x) = \varepsilon v(\varepsilon x) e^{ik_0 x} + \text{c.c.} + O(\varepsilon^2)$$

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Focusing (+) if $\gamma > \gamma_0 \approx 3.37 \times 10^{-10}$

Typical values:

$\gamma \approx 10^{-5}$ (McMurdo sound)

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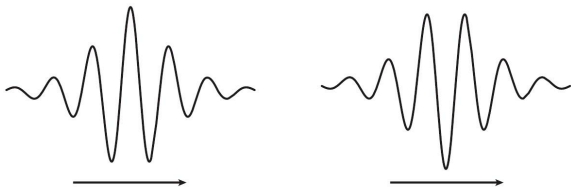
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In the focusing case NLS has solitary waves



Stability theory

Hamiltonian formulation

$$\mathcal{H}(\eta, \xi) = \frac{1}{2} \iint_{\Omega} |\nabla \phi|^2 dx dy + \frac{1}{2} \int_{\mathbb{R}} \eta^2 dx + \frac{\gamma}{2} \int_{\mathbb{R}} \kappa^2 ds$$

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$$\mathcal{H}(\eta, \xi) = \frac{1}{2} \int_{\mathbb{R}} \left\{ \xi G(\eta) \xi + \eta^2 + \gamma \frac{\eta_{xx}^2}{(1 + \eta_x^2)^{5/2}} \right\} dx$$

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Here $\xi = \phi|_{y=\eta}$ and $G(\eta)$ is the *Dirichlet-Neumann operator*

$$G(\eta)\xi = \sqrt{1 + \eta_x^2} \partial_n \phi,$$

$$\Delta \phi = 0, \quad -1 < y < \eta$$

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Travelling waves are critical points of $\mathcal{H} - c\mathcal{I}$

Main Theorem

Theorem

Let $B_R(0) = \{\eta \in H^2(\mathbb{R}) : \|\eta\|_{H^2} < R\}$, $R > 0$ given.

Assume that $\gamma > \gamma_0$ and $0 < \mu \ll 1$.

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Function space for ξ :

$$H_\star^{1/2}(\mathbb{R}) = \{u \in H_{\text{loc}}^s(\mathbb{R}) : u' \in H^{-1/2}(\mathbb{R})\} / \mathbb{R}$$
$$\|u\|_{H_\star^{1/2}} = \|u'\|_{H^{-1/2}}$$

Corollary

The set D_μ of minimisers is conditionally energetically stable.

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- ▶ Better existence and stability results than for capillary-gravity waves
- ▶ So far no global existence results under these conditions
- ▶ Local well-posedness: Ambrose & Siegel (to appear in PRSE)

Proof of the theorem

Step 1. Reduction from phase space to configuration space

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- ▶ Fix η and minimise $\mathcal{H}(\eta, \xi)$ subject to $\mathcal{I}(\eta, \xi) = 2\mu$.
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$$\xi_\eta = c_\eta G(\eta)^{-1} \eta_x$$

- ▶ Minimise

$$\mathcal{J}(\eta) = \mathcal{H}(\eta, \xi_\eta) = \frac{\mu^2}{\mathcal{L}(\eta)} + \mathcal{K}(\eta)$$

where

$$\mathcal{L}(\eta) = \frac{1}{2} \int_{\mathbb{R}} \eta_x G(\eta)^{-1} \eta_x dx$$
$$\mathcal{K}(\eta) = \int_{\mathbb{R}} \left\{ \frac{1}{2} \eta^2 + \gamma \frac{\eta_{xx}^2}{(1 + \eta_x^2)^{5/2}} \right\} dx$$

Step 1.5. Periodic problem with large period

$(H^s \subset\subset H^r, s > r)$

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- ▶ For $0 < \mu \ll 1$, \exists test function

$$\begin{aligned}\eta_\star^\mu(x) &= \mu \operatorname{sech}(\mu x) \cos(k_0 x) \\ &\quad - A\mu^2 \operatorname{sech}^2(\mu x) \cos(2k_0 x) - B\mu^2 \operatorname{sech}^2(\mu x),\end{aligned}$$

such that

$$\mathcal{J}(\eta_\star^\mu) < 2\mu$$

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Problem: The compact embedding fails on \mathbb{R} !

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To rule out dichotomy, we show that

$$I(\mu) = \inf\{\mathcal{J}(\eta) : \eta \in B_R(0)\}, \quad 0 < \mu \ll 1,$$

is strictly subhomogeneous:

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Implies strict subadditivity:

$$I(\mu_1 + \mu_2) < I(\mu_1) + I(\mu_2)$$

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Solutions:

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- ▶ The other problem is the interesting part

Idea:

For the test function we find that

$$\mathcal{J}(\eta_{\star}^{\mu}) = 2c_0\mu + I_{\text{NLS}}\mu^3 + O(\mu^4),$$

where $I_{\text{NLS}} < 0$ is the ground state energy for NLS

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Problem: Are near minimisers similar to η_{\star}^{μ} ?

Properties of minimising sequences

Any minimising sequence satisfies

$$\|\mathcal{J}'(\eta_n)\|_{H^{-2}} \rightarrow 0 \text{ and } \|\eta_n\|_{H^2}^2 \leq C\mu$$

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We use these properties to show that η_n has a form similar to η_\star^μ

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$\eta_n = \eta_{n,1} + \eta_{n,2} + \eta_{n,3}$ *where:*

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- ▶ $\hat{\eta}_{n,2}$ has support near $0, \pm 2k_0$,
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▶

$$\|\eta_{n,3}\|_{H^2}^2 \leq C\mu^5$$

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so that $I(\mu)$ is strictly subhomogeneous

Also yields convergence of v_n to solution of NLS using compactness of minimising sequences for the NLS variational problem

Reference

M. D. Groves, B. Hewer & E. Wahlén, Variational existence theory for hydroelastic solitary waves, C. R. Math. Acad. Sci. Paris **354** (2016)

Thank you for staying awake!