# Some $q$-exponential formulas involving the double lowering operator $\psi$ for a tridiagonal pair 

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## Definition of a tridiagonal pair

Let $V$ denote a finite-dimensional vector space over a field $\mathbb{K}$.

## Definition

By a tridiagonal pair (or TD pair) on $V$ we mean an ordered pair of linear transformations $A: V \rightarrow V$ and $A^{*}: V \rightarrow V$ satisfying:

1. Each of $A, A^{*}$ is diagonalizable.
2. There exists an ordering $\left\{V_{i}\right\}_{i=0}^{d}$ of the eigenspaces of $A$ such that

$$
A^{*} V_{i} \subseteq V_{i-1}+V_{i}+V_{i+1} \quad(0 \leq i \leq d)
$$

where $V_{-1}=0$ and $V_{d+1}=0$.
3. There exists an ordering $\left\{V_{i}^{*}\right\}_{i=0}^{\delta}$ of the eigenspaces of $A^{*}$ such that

$$
A V_{i}^{*} \subseteq V_{i-1}^{*}+V_{i}^{*}+V_{i+1}^{*} \quad(0 \leq i \leq \delta)
$$

where $V_{-1}^{*}=0$ and $V_{\delta+1}^{*}=0$.
4. There does not exist a subspace $W$ of $V$ such that $A W \subseteq W$,
$A^{*} W \subseteq W, W \neq 0, W \neq V$.

## Example: Q-polynomial distance-regular graph

- Let $\Gamma=\Gamma(X, E)$ denote a $Q$-polynomial distance-regular graph.
- Let $A$ denote the adjacency matrix of $\Gamma$.
- Fix $x \in X$. Let $A^{*}=A^{*}(x)$ denote the dual adjacency matrix of $\Gamma$ with respect to $x$.
- Let $W$ denote an irreducible $\left(A, A^{*}\right)$-submodule of $\mathbb{C}^{|X|}$.
- Then $A, A^{*}$ form a TD pair on $W$.


## Tridiagonal system

By a tridiagonal system (or TD system) on $V$, we mean a sequence

$$
\Phi=\left(A ;\left\{V_{i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d}\right)
$$

that satisfies (1)-(3) below.

1. $A, A^{*}$ is a tridiagonal pair on $V$.
2. $\left\{V_{i}\right\}_{i=0}^{d}$ is an ordering of the eigenspaces of $A$ such that

$$
A^{*} V_{i} \subseteq V_{i-1}+V_{i}+V_{i+1} \quad(0 \leq i \leq d)
$$

3. $\left\{V_{i}^{*}\right\}_{i=0}^{d}$ is an ordering of the eigenspaces of $A^{*}$ such that

$$
A V_{i}^{*} \subseteq V_{i-1}^{*}+V_{i}^{*}+V_{i+1}^{*} \quad(0 \leq i \leq d)
$$

## Relatives of a TD system

A given TD system can be modified in a number of ways to get a new TD system.

$$
\begin{array}{ll}
\left(A ;\left\{V_{i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d}\right) & \left(A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d} ; A ;\left\{V_{i}\right\}_{i=0}^{d}\right) \\
\left(A ;\left\{V_{d-i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d}\right) & \left(A^{*} ;\left\{V_{d-i}^{*}\right\}_{i=0}^{d} ; A ;\left\{V_{i}\right\}_{i=0}^{d}\right) \\
\left(A ;\left\{V_{i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{d-i}^{*}\right\}_{i=0}^{d}\right) & \left(A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d} ; A ;\left\{V_{d-i}\right\}_{i=0}^{d}\right) \\
\left(A ;\left\{V_{d-i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{d-i}^{*}\right\}_{i=0}^{d}\right) & \left(A^{*} ;\left\{V_{d-i}^{*}\right\}_{i=0}^{d} ; A ;\left\{V_{d-i}\right\}_{i=0}^{d}\right)
\end{array}
$$

These eight TD systems are said to be relatives of one another.

## Relatives of a TD system

A given TD system can be modified in a number of ways to get a new TD system.

$$
\begin{array}{ll}
\left(A ;\left\{V_{i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d}\right) & \left(A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d} ; A ;\left\{V_{i}\right\}_{i=0}^{d}\right) \\
\left(A ;\left\{V_{d-i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d}\right) & \left(A^{*} ;\left\{V_{d-i}^{*}\right\}_{i=0}^{d} ; A ;\left\{V_{i}\right\}_{i=0}^{d}\right) \\
\left(A ;\left\{V_{i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{d-i}^{*}\right\}_{i=0}^{d}\right) & \left(A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d} ; A ;\left\{V_{d-i}\right\}_{i=0}^{d}\right) \\
\left(A ;\left\{V_{d-i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{d-i}^{*}\right\}_{i=0}^{d}\right) & \left(A^{*} ;\left\{V_{d-i}^{*}\right\}_{i=0}^{d} ; A ;\left\{V_{d-i}\right\}_{i=0}^{d}\right)
\end{array}
$$

These eight TD systems are said to be relatives of one another.

Big Goal: Better understand the relationship between these relatives!

## Relatives of a TD system

A given TD system can be modified in a number of ways to get a new TD system.

$$
\begin{array}{rlrl}
\longrightarrow & \left(A ;\left\{V_{i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d}\right) & & \left(A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d} ; A ;\left\{V_{i}\right\}_{i=0}^{d}\right) \\
\longrightarrow & \left(A ;\left\{V_{d-i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d}\right) & \left(A^{*} ;\left\{V_{d-i}^{*}\right\}_{i=0}^{d} ; A ;\left\{V_{i}\right\}_{i=0}^{d}\right) \\
& \left(A ;\left\{V_{i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{d-i}^{*}\right\}_{i=0}^{d}\right) & \left(A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d} ; A ;\left\{V_{d-i}\right\}_{i=0}^{d}\right) \\
& \left(A ;\left\{V_{d-i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{d-i}^{*}\right\}_{i=0}^{d}\right) & \left(A^{*} ;\left\{V_{d-i}^{*}\right\}_{i=0}^{d} ; A ;\left\{V_{d-i}\right\}_{i=0}^{d}\right)
\end{array}
$$

These eight TD systems are said to be relatives of one another.

Big Goal: Better understand the relationship between these relatives!

Smaller Goal: Better understand the relationship between these 2 relatives.

## Assumptions/Notation

- Fix a TD system $\Phi=\left(A ;\left\{V_{i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d}\right)$ on $V$. Let $\Phi^{\Downarrow}=\left(A ;\left\{V_{d-i}\right\}_{i=0}^{d} ; A^{*} ;\left\{V_{i}^{*}\right\}_{i=0}^{d}\right)$ denote the second inversion of $\Phi$.
- For $0 \leq i \leq d$, we let $\theta_{i}$ (resp. $\theta_{i}^{*}$ ) denote the eigenvalue of $A$ (resp. $\left.A^{*}\right)$ corresponding to the eigenspace $V_{i}$ (resp. $V_{i}^{*}$ ).


## $q$-Racah case

## Definition

We say that the TD system $\Phi$ has $q$-Racah type whenever there exist nonzero scalars $q, a, b \in \overline{\mathbb{K}}$ such that $q^{4} \neq 1$ and

$$
\begin{aligned}
\theta_{i} & =a q^{d-2 i}+a^{-1} q^{2 i-d}, \\
\theta_{i}^{*} & =b q^{d-2 i}+b^{-1} q^{2 i-d}
\end{aligned}
$$

for $0 \leq i \leq d$.

## Assumption

Throughout this talk, we assume that $\Phi$ has $q$-Racah type. For simplicity, we also assume that $\mathbb{K}$ is algebraically closed.

## The split decompositions of $V$

## Definition

For $0 \leq i \leq d$, define

$$
\begin{aligned}
& U_{i}=\left(V_{0}^{*}+V_{1}^{*}+\cdots+V_{i}^{*}\right) \cap\left(V_{i}+V_{i+1}+\cdots+V_{d}\right), \\
& U_{i}^{\Downarrow}=\left(V_{0}^{*}+V_{1}^{*}+\cdots+V_{i}^{*}\right) \cap\left(V_{0}+V_{1}+\cdots+V_{d-i}\right) .
\end{aligned}
$$

We refer to $\left\{U_{i}\right\}_{i=0}^{d}$ as the first split decomposition of $V$.

We refer to $\left\{U_{i}^{\Downarrow}\right\}_{i=0}^{d}$ as the second split decomposition of $V$.

## The maps $K, B$

## Definition

Let $K: V \rightarrow V$ denote the linear transformation such that for $0 \leq i \leq d$, $U_{i}$ is an eigenspace of $K$ with eigenvalue $q^{d-2 i}$. That is,

$$
\left(K-q^{d-2 i} I\right) U_{i}=0
$$

for $0 \leq i \leq d$.

## Definition

Let $B: V \rightarrow V$ denote the linear transformation such that for $0 \leq i \leq d$, $U_{i}^{\Downarrow}$ is an eigenspace of $B$ with eigenvalue $q^{d-2 i}$. That is,

$$
\left(B-q^{d-2 i} I\right) U_{i}^{\Downarrow}=0
$$

for $0 \leq i \leq d$.

## The linear transformation $\psi$

There is a linear transformation $\psi: V \rightarrow V$ associated with the TD system $\Phi$. The exact definition is somewhat technical. One key feature of $\Psi$ is given below.

Lemma (B. 2012)
For $0 \leq i \leq d$, both

$$
\begin{aligned}
\psi U_{i} & \subseteq U_{i-1}, \\
\psi U_{i}^{\Downarrow} & \subseteq U_{i-1}^{\Downarrow} .
\end{aligned}
$$

Moreover, $\psi^{d+1}=0$.

In light of the above result, we refer to $\psi$ as the double lowering operator.

We see that both $K \psi=q^{2} \psi K$ and $B \psi=q^{2} \psi B$.

## The linear transformation $\Delta$

We now introduce a linear transformation $\Delta: V \rightarrow V$ which sends the first split decomposition to the second split decomposition.

Lemma (B. 2012)
There exists a unique linear transformation $\Delta: V \rightarrow V$ which satisfies

$$
\begin{aligned}
& \Delta\left(U_{i}\right) \subseteq U_{i}^{\Downarrow} \\
& (\Delta-I) U_{i} \subseteq U_{0}+U_{1}+\cdots+U_{i-1}
\end{aligned}
$$

for $0 \leq i \leq d$.

## $\triangle$ as a polynomial in $\psi$

Theorem (B. 2014)
Both

$$
\begin{aligned}
\Delta & =\sum_{i=0}^{d}\left(\prod_{j=1}^{i} \frac{a q^{j-1}-a^{-1} q^{1-j}}{q^{j}-q^{-j}}\right) \psi^{i}, \\
\Delta^{-1} & =\sum_{i=0}^{d}\left(\prod_{j=1}^{i} \frac{a^{-1} q^{j-1}-a q^{1-j}}{q^{j}-q^{-j}}\right) \psi^{i} .
\end{aligned}
$$

## $\Delta$ as a polynomial in $\psi$

Theorem (B. 2014)
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$$
\begin{aligned}
\Delta & =\sum_{i=0}^{d}\left(\prod_{j=1}^{i} \frac{a q^{j-1}-a^{-1} q^{1-j}}{q^{j}-q^{-j}}\right) \psi^{i}, \\
\Delta^{-1} & =\sum_{i=0}^{d}\left(\prod_{j=1}^{i} \frac{a^{-1} q^{j-1}-a q^{1-j}}{q^{j}-q^{-j}}\right) \psi^{i} .
\end{aligned}
$$

## Question

Does this polynomial factor nicely?
If it does, what does that factorization mean?

## The linear transformation $\mathcal{M}$

## Definition

Define a linear transformation $\mathcal{M}: V \rightarrow V$ by

$$
\mathcal{M}=\frac{a K-a^{-1} B}{a-a^{-1}}
$$

We will use this map $\mathcal{M}$ to find a factorization of $\Delta$.

## The $q$-exponential function

We now recall the $q$-exponential function. For nilpotent $T \in \operatorname{End}(V)$,

$$
\exp _{q}(T)=\sum_{n=0}^{\infty} \frac{q^{\binom{n}{2}}}{[n]_{q}^{!}} T^{n} .
$$

Here

$$
[n]_{q}^{!}=[n]_{q}[n-1]_{q} \cdots[1]_{q}
$$

and

$$
[n]_{q}=\frac{q^{n}-q^{-n}}{q-q^{-1}} .
$$

Recall that the map $\exp _{q}(T)$ is invertible and its inverse is given by

$$
\exp _{q^{-1}}(-T)=\sum_{n=0}^{\infty} \frac{(-1)^{n} q^{-\binom{n}{2}}}{[n]_{q}^{!}} T^{n}
$$

## Lemma

Both

$$
\begin{aligned}
K \exp _{q}\left(\frac{a^{-1}}{q-q^{-1}} \psi\right) & =\exp _{q}\left(\frac{a^{-1}}{q-q^{-1}} \psi\right) \mathcal{M} \\
B \exp _{q}\left(\frac{a}{q-q^{-1}} \psi\right) & =\exp _{q}\left(\frac{a}{q-q^{-1}} \psi\right) \mathcal{M}
\end{aligned}
$$

These results turns out to be the key to being able to factor the polynomial in $\psi$ for $\Delta$.

## $\Delta$ as a product of $q$-exponentials

## Theorem

Both

$$
\begin{aligned}
\Delta & =\exp _{q}\left(\frac{a}{q-q^{-1}} \psi\right) \exp _{q^{-1}}\left(-\frac{a^{-1}}{q-q^{-1}} \psi\right), \\
\Delta^{-1} & =\exp _{q}\left(\frac{a^{-1}}{q-q^{-1}} \psi\right) \exp _{q^{-1}}\left(-\frac{a}{q-q^{-1}} \psi\right) .
\end{aligned}
$$

If we multiply out the right-hand side of the above product and use the $q$-binomial theorem to simplify the coefficients, we will obtain the polynomial for $\Delta$ given earlier in the talk.

## $\Delta$ as a transition matrix

We view $\Delta$ as a transition matrix from the first split decomposition of $V$ to the second. Consequently, we view

$$
\exp _{q^{-1}}\left(-\frac{a^{-1}}{q-q^{-1}} \psi\right)
$$

as a transition matrix from the first split decomposition to a decomposition of $V$ which we interpret as a kind of half-way point.

We will describe this new decomposition of $V$ using the linear transformation $\mathcal{M}$.

## The eigenspaces of $\mathcal{M}$

## Lemma

The map $\mathcal{M}$ is diagonalizable with eigenvalues $q^{d}, q^{d-2}, q^{d-4}, \ldots, q^{-d}$.

## Definition

For $0 \leq i \leq d$ let $W_{i}$ denote the eigenspace of $\mathcal{M}$ corresponding to the eigenvalue $q^{d-2 i}$. Note that $\left\{W_{i}\right\}_{i=0}^{d}$ is a decomposition of $V$.

## The eigenspaces of $\mathcal{M}$ as a half-way point

## Lemma

For $0 \leq i \leq d$,

$$
\begin{array}{ll}
U_{i}=\exp _{q}\left(\frac{a^{-1}}{q-q^{-1}} \psi\right) W_{i}, & W_{i}=\exp _{q^{-1}}\left(-\frac{a^{-1}}{q-q^{-1}} \psi\right) U_{i}, \\
U_{i}^{\Downarrow}=\exp _{q}\left(\frac{a}{q-q^{-1}} \psi\right) W_{i}, & W_{i}=\exp _{q^{-1}}\left(-\frac{a}{q-q^{-1}} \psi\right) U_{i}^{\Downarrow .}
\end{array}
$$

## The eigenspaces of $\mathcal{M}$ as a half-way point

## Lemma

For $0 \leq i \leq d$,

$$
\begin{array}{ll}
U_{i}=\exp _{q}\left(\frac{a^{-1}}{q-q^{-1}} \psi\right) W_{i}, & W_{i}=\exp _{q^{-1}}\left(-\frac{a^{-1}}{q-q^{-1}} \psi\right) U_{i}, \\
U_{i}^{\Downarrow}=\exp _{q}\left(\frac{a}{q-q^{-1}} \psi\right) W_{i}, & W_{i}=\exp _{q^{-1}}\left(-\frac{a}{q-q^{-1}} \psi\right) U_{i}^{\Downarrow .}
\end{array}
$$

We see that $W_{i}$ is the image of $U_{i}$ under our $q^{-1}$-exponential in $\psi$ !
We now know that the decomposition that we have regarded as a half-way point between the two split decompositions is the eigenspace decomposition for $\mathcal{M}$.

## The actions of various linear transformations on $W_{i}$

Now that we know how to describe this half-way point, we can investigate the actions of our other linear transformations on this decomposition.

## The actions of $\psi, K$, and $B$

Lemma
For $0 \leq i \leq d$,

$$
\psi W_{i} \subseteq W_{i-1}
$$

Lemma
For $0 \leq i \leq d$,

$$
\begin{aligned}
& \left(K-q^{d-2 i} I\right) W_{i} \subseteq W_{i-1}, \\
& \left(B-q^{d-2 i} I\right) W_{i} \subseteq W_{i-1} .
\end{aligned}
$$

## The action of $\Delta$

Lemma
For $0 \leq i \leq d$,

$$
\begin{aligned}
(\Delta-I) W_{i} & \subseteq W_{0}+W_{1}+\cdots+W_{i-1}, \\
\left(\Delta^{-1}-I\right) W_{i} & \subseteq W_{0}+W_{1}+\cdots+W_{i-1} .
\end{aligned}
$$

Thank you for your attention!

