Probabilistic Tools for Lattice Path Enumeration

KILIAN RASCHEL



Lattice walks at the Interface of Algebra, Analysis and Combinatorics September 18, 2017 BIRS

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Introduction

Dimension 1: examples & limits

Central idea in dimension ≥ 2 : approximation by Brownian motion

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Application #1: excursions

Application #2: walks with prescribed length

Discrete harmonic functions and critical exponents

First exit time from a cone C



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First exit time from a cone C

 $\succ \tau_C = \inf\{n \in \mathbf{N} : S(n) \notin C\} (S \text{ RW})$ $\triangleright T_C = \inf\{t \in \mathbf{R}_+ : B(t) \notin C\} (B \text{ BM})$



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Persistence probabilities → total number of walks

$$\triangleright \mathbf{P}_{x}[\tau_{\mathcal{C}} > n] \sim \kappa \cdot V(x) \cdot \rho^{n} \cdot n^{-\alpha}$$

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Local limit theorems ~> excursions

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Aim of the talk: understanding the critical exponents α

Random walk on Z^d

▷ A random walk $\{S(n)\}_{n \ge 0}$ is

$$S(n) = x + X(1) + \cdots + X(n),$$

where the X(i) are i.i.d.



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▷ Example (Dyck paths): simple random walk $X(i) \in \{-1, +1\}$

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The ubiquity of random walks



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$$\triangleright \ \#\{x \stackrel{n}{\longrightarrow} \mathbf{Z}\} = 2^n$$

Walk \rightsquigarrow Exponent 0

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▷ $\#\{x \xrightarrow{n} \mathbf{Z}\} = 2^n$ Walk \rightsquigarrow Exponent 0 ▷ $\#\{x \xrightarrow{n} y\} = {\binom{n}{\frac{n+(y-x)}{2}}} \sim \sqrt{\frac{2}{\pi}} \frac{2^n}{\sqrt{n}}$ Bridge \rightsquigarrow Exponent $\frac{1}{2}$

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▷ $\#\{x \xrightarrow{n} Z\} = 2^n$ Walk \rightsquigarrow Exponent 0 ▷ $\#\{x \xrightarrow{n} y\} = {n \choose \frac{n+(y-x)}{2}} \sim \sqrt{\frac{2}{\pi}} \frac{2^n}{\sqrt{n}}$ Bridge \rightsquigarrow Exponent $\frac{1}{2}$ ▷ $\sum \frac{1}{\sqrt{n}} = \infty$: recurrence of the simple random walk in Z



▷ #{x → Z} = 2ⁿ Walk → Exponent 0 ▷ #{x → y} = $\binom{n}{\frac{n+(y-x)}{2}} \sim \sqrt{\frac{2}{\pi}} \frac{2^n}{\sqrt{n}}$ Bridge → Exponent $\frac{1}{2}$ ▷ $\sum \frac{1}{\sqrt{n}} = \infty$: recurrence of the simple random walk in Z ▷ Constant $\sqrt{\frac{2}{\pi}}$ independent of x & y in the asymptotics



Constrained walk with $\mathfrak{S} = \{-1, +1\}$ (Dyck paths)

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Beyond the classical exponents 0, $\frac{1}{2}$ & $\frac{3}{2}$

Weighted models in dimension 1

Drift $\sum_{s\in\mathfrak{S}} s$ governs the exponents, which are still 0, $\frac{1}{2}$ & $\frac{3}{2}$

Beyond the classical exponents 0, $\frac{1}{2}$ & $\frac{3}{2}$

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Weighted models in dimension 1

Drift $\sum_{s \in \mathfrak{S}} s$ governs the exponents, which are still 0, $\frac{1}{2}$ & $\frac{3}{2}$ The simple walk in two-dimensional wedges

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Weighted models in dimension 1

Drift $\sum_{s \in \mathfrak{S}} s$ governs the exponents, which are still 0, $\frac{1}{2}$ & $\frac{3}{2}$ The simple walk in two-dimensional wedges



- Half-plane: one-dimensional case
- Dyck paths
- ▷ Total number of walks: \rightarrow Exponent $\frac{1}{2}$
- Excursions:

 \rightsquigarrow Exponent $2 = \frac{3}{2} + \frac{1}{2}$

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Weighted models in dimension 1

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- Quarter plane: product of two one-dimensional cases
- Reflection principle
- ▷ Total number of walks: \rightarrow Exponent $1 = \frac{1}{2} + \frac{1}{2}$
- Excursions:

 \rightsquigarrow Exponent $3 = \frac{3}{2} + \frac{3}{2}$

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Weighted models in dimension 1

Drift $\sum_{s \in \mathfrak{S}} s$ governs the exponents, which are still 0, $\frac{1}{2}$ & $\frac{3}{2}$ The simple walk in two-dimensional wedges



- Slit plane:
 Bousquet-Mélou & Schaeffer '00
 - Highly non-convex cone
 - $\triangleright \text{ Total number of walks:} \\ \rightsquigarrow \text{ Exponent } \frac{1}{4}$

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 $\triangleright \text{ Excursions:} \\ \rightsquigarrow \text{ Exponent } \frac{3}{2}$

Weighted models in dimension 1

Drift $\sum_{s \in \mathfrak{S}} s$ governs the exponents, which are still 0, $\frac{1}{2}$ & $\frac{3}{2}$ The simple walk in two-dimensional wedges



▷ 45°: Souyou-Beauchamps '86

▷ See

🖗 Bousquet-Mélou & Mishna '10

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- Excursions:
 - \rightsquigarrow Exponent 5

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- ▷ 135°: Gessel
- See See Kauers, Koutschan & Zeilberger '09; etc.
- ▷ Total number of walks: \rightarrow Exponent $\frac{2}{3}$

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 $\triangleright \text{ Excursions:} \\ \rightsquigarrow \text{ Exponent } \frac{7}{3}$

Weighted models in dimension 1

Drift $\sum_{s \in \mathfrak{S}} s$ governs the exponents, which are still 0, $\frac{1}{2} \& \frac{3}{2}$ The simple walk in two-dimensional wedges



- Walks avoiding a quadrant
- See See Bousquet-Mélou '15;
 Mustapha '15; Trotignon et al. '17+

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- ▷ Total number of walks: \rightarrow Exponent $\frac{1}{3}$
- $\triangleright \quad \text{Excursions:} \\ \rightsquigarrow \quad \text{Exponent } \frac{5}{3}$

Weighted models in dimension 1

Drift $\sum_{s \in \mathfrak{S}} s$ governs the exponents, which are still 0, $\frac{1}{2}$ & $\frac{3}{2}$ The simple walk in two-dimensional wedges



- \triangleright Arbitrary angular sector θ
- ▷ See [®] Varopoulos '99; Denisov & Wachtel '15

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- \triangleright Arbitrary angular sector θ
- ▷ See [®] Varopoulos '99; Denisov & Wachtel '15
- ▷ Total number of walks: \rightarrow Exponent $\frac{\pi}{2\theta}$
- Excursions:
 - \rightsquigarrow Exponent $\frac{\pi}{\theta} + 1$

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 - \rightsquigarrow Exponent $\frac{\pi}{\theta}+1$

Conclusion: 1D case not enough

- Dramatic change of behavior: every exponent is possible!
- Non-D-finite behaviors (first observed by Varopoulos '99)

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Law of large numbers

$$\frac{X(1) + \dots + X(n)}{n^1} \stackrel{\text{a.s.}}{\longrightarrow} \mathbf{E}[X(1)]$$

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Law of large numbers

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Central limit theorem

$$n^{\frac{1}{2}}\left\{\frac{X(1)+\cdots+X(n)}{n^{1}}-\mathsf{E}[X(1)]\right\}\stackrel{\mathsf{law}}{\longrightarrow}\mathcal{N}(0,\mathsf{V}[X(1)])$$

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Donsker's theorem (functional central limit theorem)



 $RW \longrightarrow BM$

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Denisov & Wachtel '15 (excursions for RW in cones $\subset Z^d$)

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- $\triangleright \ \mathsf{RW} \longrightarrow \mathsf{BM}$
- Mapping theorem: many asymptotic results concerning RW can be deduced from BM

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▷ For excursions, α {RW} = α {BM} if $\begin{cases}
E[RW] = E[BM] = 0 \\
V[RW] = V[BM] = id
\end{cases}$

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 \mathbf{E}[RW] = \mathbf{E}[BM] = 0 \\
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 \end{cases}$
- ▷ If $\mathbf{V}[\mathsf{RW}] \neq \mathsf{id}$ then $\mathbf{V}[M \cdot \mathsf{RW}] = \mathsf{id}$ for some $M \in \mathbf{M}_d(\mathbf{R})$

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 \end{cases}$
- $\triangleright \ \, \mathsf{If} \ \, \mathbf{V}[\mathsf{RW}] \neq \mathsf{id} \ \mathsf{then} \ \, \mathbf{V}[M \cdot \mathsf{RW}] = \mathsf{id} \ \mathsf{for} \ \mathsf{some} \ \, M \in \mathbf{M}_d(\mathbf{R})$
- \triangleright Cone *C* becomes $M \cdot C$

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Remainder of this section: computing α {BM} (easier)

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Two derivations of the BM persistence probability in R

Reflection principle



$$\begin{aligned} \mathbf{P}_{x}[T_{(0,\infty)} > t] &= \mathbf{P}_{0}[\min_{0 \le u \le t} B(u) > -x] \\ &= \mathbf{P}_{0}[|B(t)| < x] \\ &= \frac{2}{\sqrt{2\pi t}} \int_{0}^{x} e^{-\frac{y^{2}}{2t}} dy \end{aligned}$$

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Two derivations of the BM persistence probability in R

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Heat equation

Function $g(t; x) = \mathbf{P}_x[T_{(0,\infty)} > t]$ satisfies

$$\begin{cases} \left(\frac{\partial}{\partial t} - \frac{1}{2}\Delta\right)g(t;x) = 0, & \forall x \in (0,\infty), \ \forall t \in (0,\infty) \\ g(0;x) = 1, & \forall x \in (0,\infty) \\ g(t;0) = 0, & \forall t \in (0,\infty) \end{cases}$$

Dimension *d*: explicit expression for $P_x[T_C > t]$

Heat equation

🔊 Doob '55

For essentially any domain C in any dimension d, $\mathbf{P}_x[T_C > t] \& p^C(t; x, y) (\mathbf{P}_x[T_C > t] = \int_C p^C(t; x, y) dy)$ satisfy heat equations

Dimension *d*: explicit expression for $P_x[T_C > t]$

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Dirichlet eigenvalues problem

🕲 Chavel '84

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$$\Delta_{\mathbf{S}^{d-1}}m = -\lambda m \quad \text{in } \mathbf{S}^{d-1} \cap C$$
$$m = 0 \qquad \text{in } \partial(\mathbf{S}^{d-1} \cap C)$$

Dimension d: explicit expression for $P_x[T_c > t]$ Heat equation Doob '55 For essentially any domain C in any dimension d, $\mathbf{P}_{x}[T_{C} > t]$ & $p^{C}(t; x, y)$ ($\mathbf{P}_{x}[T_{C} > t] = \int_{C} p^{C}(t; x, y) dy$) satisfy heat equations **Dirichlet eigenvalues problem** Chavel '84 $\begin{cases} \Delta_{\mathbf{S}^{d-1}}m = -\lambda m & \text{in } \mathbf{S}^{d-1} \cap C \\ m = 0 & \text{in } \partial(\mathbf{S}^{d-1} \cap C) \end{cases}$ Discrete eigenvalues $0 < \lambda_1 < \lambda_2 \leq \lambda_3 \leq \dots$ and eigenfunctions m_1, m_2, m_3, \dots

Dimension d: explicit expression for $P_x[T_c > t]$ Heat equation Doob '55 For essentially any domain C in any dimension d, $\mathbf{P}_{x}[T_{C} > t]$ & $p^{C}(t; x, y)$ ($\mathbf{P}_{x}[T_{C} > t] = \int_{C} p^{C}(t; x, y) dy$) satisfy heat equations **Dirichlet eigenvalues problem** Chavel '84 $\begin{cases} \Delta_{\mathbf{S}^{d-1}}m = -\lambda m & \text{in } \mathbf{S}^{d-1} \cap C \\ m = 0 & \text{in } \partial(\mathbf{S}^{d-1} \cap C) \end{cases}$ $S^{d-1} \cap C$ Discrete eigenvalues $0 < \lambda_1 < \lambda_2 \leq \lambda_3 \leq \dots$ and eigenfunctions m_1, m_2, m_3, \dots Series expansion DeBlassie '87; Bañuelos & Smits '97 ∞

$$\mathbf{P}_{x}[T_{C} > t] = \sum_{j=1}^{\infty} B_{j}(|x|^{2}/t)m_{j}(x/|x|)$$

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Series expansion

DeBlassie '87; Bañuelos & Smits '97

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$$\mathbf{P}_{x}[T_{C} > t] = \sum_{j=1}^{\infty} B_{j}(|x|^{2}/t)m_{j}(x/|x|),$$

with

- ▷ *B_j* hypergeometric
- ▷ series expansion very well suited for asymptotics

Series expansion

DeBlassie '87; Bañuelos & Smits '97

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Asymptotic result

DeBlassie '87; Bañuelos & Smits '97

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$$\mathbf{P}_{x}[\mathbf{T}_{C} > t] \sim \kappa \cdot u(x) \cdot t^{-\alpha},$$

Series expansion

🖄 DeBlassie '87; Bañuelos & Smits '97

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Asymptotic result

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 $\mathbf{P}_{x}[T_{C} > t] \sim \kappa \cdot u(x) \cdot t^{-\alpha},$ with $\alpha = \frac{1}{2} \left(\sqrt{\lambda_{1} + (\frac{d}{2} - 1)^{2}} - (\frac{d}{2} - 1) \right)$ linked to first eigenvalue

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 linked to first eigenvalue

Exercise

Recover the exponent $\frac{\pi}{2\theta}$ of the persistence probability for a simple random walk in a two-dimensional wedge of opening angle θ

Introduction

Dimension 1: examples & limits

Central idea in dimension ≥ 2 : approximation by Brownian motion

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Application #1: excursions

Application #2: walks with prescribed length

Discrete harmonic functions and critical exponents

In the quarter plane





In the quarter plane



Hypotheses on the *moments*:

$$\mathbf{E}[GB] = (1,0) + (1,-1) + (-1,0) + (-1,1)$$

= (0,0)

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In the quarter plane



Hypotheses on the *moments*:

$$\begin{aligned} \mathbf{E}[GB] &= (1,0) + (1,-1) + (-1,0) + (-1,1) \\ &= (0,0) \\ \mathbf{V}[GB] &= \begin{pmatrix} 4 & -2 \\ -2 & 2 \end{pmatrix} \neq \mathsf{id} \end{aligned}$$

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In the quarter plane



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Changing the cone



In the quarter plane



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Changing the cone



- \triangleright Wedge of angle $\theta = \frac{\pi}{4}$
- ▷ Total number of walks: \Rightarrow Exponent $\frac{\pi}{2\theta} = 2$

Excursions:

 \rightsquigarrow Exponent $\frac{\pi}{\theta} + 1 = 5$

Example #2: quadrant walks

A scarecrow



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$$\triangleright \mathbf{E} = (0,0) \& \mathbf{V} = \begin{pmatrix} 4 & -1 \\ -1 & 4 \end{pmatrix} \neq \mathsf{id}$$

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A scarecrow



$$\triangleright \mathbf{E} = (0,0) \& \mathbf{V} = \begin{pmatrix} 4 & -1 \\ -1 & 4 \end{pmatrix} \neq \mathrm{id}$$
$$\triangleright \theta = \arccos\left(-\frac{1}{4}\right) \Longrightarrow \alpha = \frac{\pi}{\theta} + 1 \notin \mathbf{Q}$$

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A scarecrow



► **E** = (0,0) & **V** =
$$\begin{pmatrix} 4 & -1 \\ -1 & 4 \end{pmatrix} \neq \text{id}$$
► θ = $\arccos\left(-\frac{1}{4}\right) \implies \alpha = \frac{\pi}{\theta} + 1 \notin \mathbf{Q}$
► $\sum_{n=0}^{\infty} \#_{\mathbf{N}^2}\{(0,0) \xrightarrow{n} (0,0)\}t^n$
non-D-finite

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▷ Systematic computation of $\alpha = \arccos\{algebraic number\}$





In dimension 2 (excursions only) Sostan, R. & Salvy '14

▷ Systematic computation of $\alpha = \arccos{algebraic number}$

▷ Walks with small steps:

 $\triangleright \frac{\pi}{\alpha} \in \mathbf{Q}$ iff

- generating function of the excursions is D-finite iff
- ▷ the group of the model is finite





In dimension 2 (excursions only) 🔊 Bostan, R. & Salvy '14

▷ Systematic computation of $\alpha = \arccos{algebraic number}$

▷ Walks with small steps:

 $\triangleright \ \frac{\pi}{\alpha} \in \mathbf{Q}$ iff

▷ generating function of the excursions is D-finite iff

▷ the group of the model is finite

▷ If $\sum_{s \in \mathfrak{S}} s \neq 0$, first perform a *Cramér transform*

Example: Kreweras 3D

Model with jumps:





Example: Kreweras 3D

Model with jumps:

Exponent $\alpha = \frac{1}{2}\sqrt{\lambda_1 + \frac{1}{4}} - \frac{1}{4}$





Example: Kreweras 3D

Model with jumps:







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Example: Kreweras 3D

Model with jumps:







Value of λ_1 ? $\lambda_1 \in \mathbf{Q}$?



Model with jumps:







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Value of λ_1 ? $\lambda_1 \in \mathbf{Q}$?

General theory (still to be done!)

▷ Classification & resolution of some finite group models

🕾 Bostan, Bousquet-Mélou, Kauers & Melczer '16

- ▷ Asymptotic simulation Sector, Kauers & Yatchak '16; Guttmann '16
 → Conjectured Kreweras exponent ≈ 3.3257569
- Equivalence finite group iff D-finite generating functions?

Eigenvalues of spherical triangles and 3D models A (the?) soluble case





$$\triangleright$$
 SRW in 3D: $eta=rac{\pi}{2}$ and $\lambda_1=12$

Generic case



No closed-form formula known

 \triangleright Even for a flat triangle in $R^2,$ no closed-form formula for smallest eigenvalue...

▷ Is there a miracle for Kreweras? $(\beta = \delta = \varepsilon = \frac{2\pi}{3})$ \rightsquigarrow Tetrahedral tiling of the sphere

Central weightings and stability of the exponent Critical exponents for weighted GB model



 ≈ 4.9042377

Central weightings and stability of the exponent Critical exponents for weighted GB model



Central weightings

- \triangleright Replace the initial weight 1 of jump (i, j) by $c \cdot a^i \cdot b^j$
- ▷ *Critical exponent* for the excursions *unchanged*
- ▷ Second example above: $a = \frac{1}{\sqrt{6}}$, $b = \frac{1}{3}$, $c = \sqrt{6}$
- > Third example is not a central weighting

Much more in Julien Courtiel's talk!

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Excursions: formula for α independent of the drift $\sum_{s \in \mathfrak{S}} s$

Excursions: formula for α independent of the drift $\sum_{s \in \mathfrak{S}} s$

Case #1: interior drift



Non-universal exponents: six cases **Excursions:** formula for α independent of the drift $\sum_{s \in \mathfrak{S}} s$

Case #1: interior drift



▷ Law of large numbers: $\mathbf{P}[\forall n, S(n) \in C] > 0$

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 $\triangleright \ \mathsf{Exponent} \ \alpha = \mathbf{0}$

Excursions: formula for α independent of the drift $\sum_{s \in \mathfrak{S}} s$

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- ▷ Half-plane case
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- Cannot be used as a filter to detect non-D-finiteness



- Half-plane case
- ▷ Exponent $\alpha = \frac{1}{2}$
- Cannot be used as a filter to detect non-D-finiteness
- $\triangleright \text{ Exponent } \alpha = \frac{i}{2} \text{ for non-smooth}$ boundary





- ▷ Half-plane case
- ▷ Exponent $\alpha = \frac{3}{2}$
- Cannot be used as a filter to detect non-D-finiteness

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- ▷ Half-plane case
- ▷ Exponent $\alpha = \frac{3}{2}$
- Cannot be used as a filter to detect non-D-finiteness

Case #4: zero drift



- ▷ See [©] Varopoulos '99; Denisov & Wachtel '15
- Exponent

$$\alpha_1 = \frac{1}{2} \left(\sqrt{\lambda_1 + (\frac{d}{2} - 1)^2} - (\frac{d}{2} - 1) \right)$$

Can be used as a filter to detect non-D-finiteness

Case #5: polar interior drift



- 🕞 See 🥯 Duraj '14
 - \triangleright Exponent $2\alpha_1 + 1$
 - Can be used as a filter to detect non-D-finiteness

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Case #5: polar interior drift



- 🕞 👂 🔊 Duraj '14
- \triangleright Exponent $2\alpha_1 + 1$
 - Can be used as a filter to detect non-D-finiteness

Case #6: polar boundary drift



- \triangleright Exponent $\alpha_1 + 1$
- Can be used as a filter to detect non-D-finiteness

Case #5: polar interior drift



- ▷ See [©] Duraj '14
- \triangleright Exponent $2\alpha_1 + 1$
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Case #6: polar boundary drift



- \triangleright Exponent $\alpha_1 + 1$
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Weighted GB model: with J. Courtiel, S. Melczer & M. Mishna

Case #5: polar interior drift



- ▷ See [©] Duraj '14
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Case #6: polar boundary drift



- \triangleright Exponent $\alpha_1 + 1$
- Can be used as a filter to detect non-D-finiteness

Six-exponents-result: joint with R. Garbit & S. Mustapha

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Discrete harmonic functions and critical exponents
Absorption probabilities for the SRW on N





Absorption probabilities for the SRW on N



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Probability $a(i) := \mathbf{P}_i[\exists n \ge 0 : \text{SRW } S(n) = 0]$ satisfies $\triangleright a(0) = 1 \rightsquigarrow \text{initial condition}$ $\triangleright a(i) = \mathbf{p} \cdot a(i+1) + (1-\mathbf{p}) \cdot a(i-1) \rightsquigarrow \text{recurrence}$

Absorption probabilities for the SRW on N



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Absorption probabilities for the SRW on N



Probability $a(i) := \mathbf{P}_i[\exists n \ge 0 : \text{SRW } S(n) = 0]$ satisfies $\triangleright a(0) = 1 \rightsquigarrow \text{initial condition}$ $\triangleright a(i) = p \cdot a(i+1) + (1-p) \cdot a(i-1) \rightsquigarrow \text{recurrence}$ Solution $a(i) = \begin{cases} 1 & \text{if } p \le \frac{1}{2} \\ \left(\frac{1-p}{p}\right)^i & \text{if } p > \frac{1}{2} \end{cases}$

Definition: f harmonic if L[f](x) = 0 for all x in a region $\subset \mathbf{Z}^d$

$$L[f](x) = \sum_{y \in N_x} p(y) \{ f(x+y) - f(x) \},\$$

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with set of neighbors $N_{\times} \subset \mathbf{Z}^d$ and weights $p = \{p(y)\}_{y \in \mathbf{Z}^d}$

Absorption probabilities for the SRW on N



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Multivariate linear recurrences with constant coefficients

Bousquet-Mélou & Petkovšek '00

520



$$P \quad q(n; i, j) = \#_{\mathbb{N}^2} \{ (0, 0) \xrightarrow{n} (i, j) \}$$

$$P \quad q(n+1; i, j) = q(n; i-1, j) + q(n; i+1, j) + q(n; i, j-1) + q(n; i, j+1)$$

$$(Caloric functions)$$

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> q(n; i, j) = #_{N²}{(0, 0) → (i, j)}
> q(n + 1; i, j) =
q(n; i - 1, j) + q(n; i + 1, j) + q(n; i, j - 1) + q(n; i, j + 1)
(Caloric functions)
> f(i, j) =
$$\frac{1}{4}$$
{f(i - 1, j) + f(i + 1, j) + f(i, j - 1) + f(i, j + 1)}
(Preharmonic functions)

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Main differences & difficulties

 $\triangleright\,$ A unique solution vs. an unknown ($\leqslant\infty)$ number of solutions

Consequence: guess and prove techniques do not work



▷ q(n; i, j) = #_{N²}{(0, 0)
$$\xrightarrow{n}$$
 (i, j)}
▷ q(n + 1; i, j) =
q(n + 1; i, j) + q(n; i + 1, j) + q(n; i, j - 1) + q(n; i, j + 1)
(Caloric functions)
▷ f(i, j) =
 $\frac{1}{4}$ {f(i - 1, j) + f(i + 1, j) + f(i, j - 1) + f(i, j + 1)}
(Preharmonic functions)

Main differences & difficulties

- \triangleright A unique solution vs. an unknown ($\leqslant \infty$) number of solutions
- Consequence: guess and prove techniques do not work
- Generating functions of preharmonic functions satisfy kernel functional equations
- \triangleright Preharmonic functions \approx homogenized enumeration problem:

K(x, y)Q(x, y) = K(x, 0)Q(x, 0) + K(0, y)Q(0, y) - K(0, 0)Q(0, 0) - xyK'(x, y)F(x, y) = K'(x, 0)F(x, 0) + K'(0, y)F(0, y) - K'(0, 0)F(0, 0)



▷ q(n; i, j) = #_{N²}{(0, 0)
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▷ Preharmonic functions → counting numbers asymptotics

Two examples of rational harmonic functions

The simple walk



- \triangleright Uniform weights $\frac{1}{4}$
- $\triangleright f(i,j) = i \cdot j$
- Unique preharmonic function (up to multiplicative factors)
- ▷ Product form Second Picardello & Woess '92





- \triangleright Uniform weights $\frac{1}{3}$
- $\triangleright f(i,j) = i \cdot j \cdot (i+j)$
- Unique preharmonic function (up to multiplicative factors) Siane '92

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Asymptotic statements

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Total number of walks starting at
$$(k, \ell)$$
:
 $q(n; k, \ell; \mathbb{N}^2) = \#_{\mathbb{N}^2}\{(k, \ell) \xrightarrow{n} \mathbb{N}^2\}$
 $\sim f_1(k, \ell) \cdot \rho_1^n \cdot n^{\alpha_1}$

Not proved yet!

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$$Excursions starting at (k, ℓ) :
 $q(n; k, \ell; i, j) = \#_{\mathbb{N}^2}\{(k, \ell) \xrightarrow{n} (i, j)\}$
 $\sim f_2(k, \ell) \cdot f'_2(i, j) \cdot \rho_2^n \cdot n^{\alpha_2}$
 $\square Denison & Wachtel '15$$$

Asymptotic statements



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🖄 Denisov & Wachtel '15

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Preharmonicity of the prefactors

 \triangleright f_1 is ρ_1 -harmonic & f_2 is ρ_2 -harmonic

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▷ f_1 is ρ_1 -harmonic & f_2 is ρ_2 -harmonic: replace $q(n; k, \ell; \mathbb{N}^2)$ by its asymptotic expansion in the step-by-step construction $q(n+1; k, \ell; \mathbb{N}^2) = \sum_{(i,j) \in S} q(n; k-i, \ell-j; \mathbb{N}^2)$

Asymptotic statements



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Asymptotic statements



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Denisov & Wachtel '15

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Preharmonicity of the prefactors

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- \triangleright f'_2 is ρ_2 -harmonic for the *reversed step set* S' = -S
- \triangleright Drift zero: unique harmonic function \Longrightarrow f_1 , f_2 and f'_2

A functional equation reminiscent of the enumeration



$$F(x, y) = \sum_{i,j \ge 1} f(i,j) x^{i-1} y^{j-1}$$

$$F'(x, y) = xy \{ \sum_{-1 \le k, \ell \le 1} p(k, \ell) x^{-k} y^{-\ell} - 1 \}$$

$$F'(x, y) = xy \{ \sum_{-1 \le k, \ell \le 1} p(k, \ell) x^{-k} y^{-\ell} - 1 \}$$

 $\begin{aligned} & K'(x, y)F(x, y) = \\ & K'(x, 0)F(x, 0) + K'(0, y)F(0, y) - K'(0, 0)F(0, 0) \end{aligned}$

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A functional equation reminiscent of the enumeration





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Definition of Tutte's invariants

- Introduced to count q-colored triangulations & planar maps Tutte '73; Bernardi & Bousquet-Mélou '11
- ▷ Define X_0 & X_1 by $K'(X_0, y) = K'(X_1, y) = 0$
- \triangleright Tutte's invariant: function $I \in {\sf Q}[[x]]$ such that $I(X_0) = I(X_1)$

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The sections K'(x, 0)F(x, 0) & K'(0, y)F(0, y) are invariants

- \triangleright Evaluate the functional equation at $X_0 \& X_1$
- Make the difference of the two identities

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- Make the difference of the two identities

Does this characterize the sections?

A product-form generating function

$$f(i,j) = i \cdot j \implies \left[F(x,y) = \sum_{i,j \ge 1} i \cdot j \cdot x^{i-1} y^{j-1} = \frac{1}{(1-x)^2 (1-y)^2} \right]$$

Kernel: $K'(x,y) = xy\{\frac{x}{4} + \frac{1}{4x} + \frac{y}{4} + \frac{1}{4y} - 1\} = \frac{y(x-1)^2}{4} + \frac{x(y-1)^2}{4}$

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Verification of the functional equation

K'(x, y)F(x, y) = K'(x, 0)F(x, 0) + K'(0, y)F(0, y) - K'(0, 0)F(0, 0)

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Verification of the functional equation

$$K'(x, y)F(x, y) = \frac{K'(x, 0)F(x, 0) + K'(0, y)F(0, y) - K'(0, 0)F(0, 0)}{\frac{x}{4} \times \frac{1}{(1-x)^2} + \frac{y}{4} \times \frac{1}{(1-y)^2} - \frac{1}{0} \times 1$$

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Tutte's invariants

$$I(X_0) = I(X_1) \xrightarrow{X_0 X_1 = 1} I(x) = I(\frac{1}{x}) \Longrightarrow I \text{ function of } x + \frac{1}{x}$$

$$K'(x, 0)F(x, 0) = \frac{x}{4} \frac{1}{(1-x)^2} = \frac{1}{4} \frac{1}{x + \frac{1}{x} - 2} \text{ is an invariant}$$

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A product-form generating function

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Verification of the functional equation

$$K'(x,y)F(x,y) = \frac{K'(x,0)F(x,0) + K'(0,y)F(0,y) - K'(0,0)F(0,0)}{\frac{x}{4} \times \frac{1}{(1-x)^2} + \frac{y}{4} \times \frac{1}{(1-y)^2} - 0 \times 1$$

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$$K'(x,0)F(x,0) = \frac{x}{4}\frac{1}{(1-x)^2} = \frac{1}{4}\frac{1}{x+\frac{1}{x}-2} \text{ is an invariant}$$

Why this function of $x + \frac{1}{x}$?

▷ Of order 1 in $x + \frac{1}{x} \rightarrow Minimality$ (conformal mappings) ▷ $F(1, 0) = \infty \rightarrow Liouville's theorem$

Tutte's invariants & conformal mappings A general theorem

K'(x,0)F(x,0) = w(x), characterized by

Conformal mapping of a quartic

$$\triangleright w(x) = w(\overline{x})$$

- $\triangleright w(x) = \frac{c+o(1)}{(1-x)^{\alpha-1}}, \alpha = \text{crit. exponent}$
- ▷ Same for K'(0, y)F(0, y)



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Going back to the SRW

$$K'(x,0)F(x,0) = \frac{x}{4(1-x)^2}$$
, characterized by

Conformal mapping of the unit disc

$$\triangleright w(e^{i\theta}) = w(e^{-i\theta})$$

$$\triangleright w(1) = \infty$$

▷ Same for
$$K'(0, y)F(0, y) = \frac{y}{4(1-y)^2}$$





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Question

How deep is this connection conformal maps/harmonic functions?





