# Knotting statistics for polygons in lattice tubes 

Nicholas Beaton ${ }^{1}$ Jeremy Eng ${ }^{2} \quad$ Chris Soteros ${ }^{2}$<br>${ }^{1}$ School of Mathematics and Statistics, University of Melbourne, Australia<br>${ }^{2}$ Department of Mathematics and Statistics, University of Saskatchewan, Canada

The Topology of Nucleic Acids: Research at the Interface of Low-Dimensional Topology, Polymer Physics and Molecular Biology

Banff International Research Station
March 25-29, 2019


THE UNIVERSITY OF MELBOURNE

## SAPs in $\mathbb{Z}^{3}$

A self-avoiding polygon (SAP) is a simple closed walk on the edges of $\mathbb{Z}^{3}$ :


## SAPs in $\mathbb{Z}^{3}$

A self-avoiding polygon (SAP) is a simple closed walk on the edges of $\mathbb{Z}^{3}$ :


Let $p_{n}$ be the number of $n$-edge polygons, defined up to translation ( $n$ must be even). Then

$$
\left(p_{2 m}\right)_{m \geq 1}=0,3,22,207,2412,31754,452640,6840774,108088232,1768560270, \ldots
$$

## SAPs in $\mathbb{Z}^{3}$

A self-avoiding polygon (SAP) is a simple closed walk on the edges of $\mathbb{Z}^{3}$ :


Let $p_{n}$ be the number of $n$-edge polygons, defined up to translation ( $n$ must be even). Then

$$
\left(p_{2 m}\right)_{m \geq 1}=0,3,22,207,2412,31754,452640,6840774,108088232,1768560270, \ldots
$$

(Known up to $n=32$. [Clisby et al 2007])

## Asymptotics of $p_{n}$

## Theorem (Hammersley 1961)

There exists $\kappa=\log \mu$ such that

$$
p_{n}=\exp \{\kappa n+o(n)\}=e^{o(n)} \mu^{n}
$$

Either $\kappa$ or $\mu$ (depending on who you ask) is known as the connective constant of the lattice.

## Asymptotics of $p_{n}$

## Theorem (Hammersley 1961)

There exists $\kappa=\log \mu$ such that

$$
p_{n}=\exp \{\kappa n+o(n)\}=e^{o(n)} \mu^{n} .
$$

Either $\kappa$ or $\mu$ (depending on who you ask) is known as the connective constant of the lattice.
For the cubic lattice [Clisby 2013]

$$
\mu \approx 4.684039931
$$

## Asymptotics of $p_{n}$

## Theorem (Hammersley 1961)

There exists $\kappa=\log \mu$ such that

$$
p_{n}=\exp \{\kappa n+o(n)\}=e^{o(n)} \mu^{n} .
$$

Either $\kappa$ or $\mu$ (depending on who you ask) is known as the connective constant of the lattice.

For the cubic lattice [Clisby 2013]

$$
\mu \approx 4.684039931
$$

The $e^{o(n)}$ is conjectured to follow a power law, so that

$$
p_{n} \sim A n^{\alpha-3} \mu^{n}
$$

for some constants $A$ and $\alpha$. The exponent $\alpha$ is expected to be universal (the same for any 3-dimensional lattice), while $A$ and $\mu$ are lattice-dependent. In 3D [Clisby \& Dünweg 2016]

$$
\alpha \approx 0.237209
$$

## Knotted polygons

In three dimensions SAPs can be knotted:


Let $p_{n}(K)$ be the number of $n$-edge polygons of knot type $K$.

## Knotted polygons cont'd

## Theorem (Sumners \& Whittington 1988)

There exists $\kappa_{0}=\log \mu_{0}$ such that

$$
p_{n}\left(0_{1}\right)=\exp \left\{\kappa_{0} n+o(n)\right\}=e^{o(n)} \mu_{0}^{n} .
$$

Moreover

$$
\mu_{0}<\mu .
$$

That is, the probability of a random n-edge polygon being unknotted decays exponentially:

$$
\mathbb{P}_{n}\left(0_{1}\right)=\frac{p_{n}\left(0_{1}\right)}{p_{n}}=e^{o(n)}\left(\frac{\mu_{0}}{\mu}\right)^{n}
$$

## Knotted polygons cont'd

## Theorem (Sumners \& Whittington 1988)

There exists $\kappa_{0}=\log \mu_{0}$ such that

$$
p_{n}\left(0_{1}\right)=\exp \left\{\kappa_{0} n+o(n)\right\}=e^{o(n)} \mu_{0}^{n} .
$$

Moreover

$$
\mu_{0}<\mu .
$$

That is, the probability of a random n-edge polygon being unknotted decays exponentially:

$$
\mathbb{P}_{n}\left(0_{1}\right)=\frac{p_{n}\left(0_{1}\right)}{p_{n}}=e^{o(n)}\left(\frac{\mu_{0}}{\mu}\right)^{n} .
$$

But the decay is still slow [Janse van Rensburg 2008]:

$$
\frac{\mu_{0}}{\mu} \approx 0.999996
$$

## Knotted polygons cont'd

For $K \neq 0_{1}$ very little has been proved. It is straightforward to show that

$$
\limsup _{n \rightarrow \infty} \frac{1}{n} \log p_{n}(K) \geq \kappa_{0} .
$$

## Knotted polygons cont'd

For $K \neq 0_{1}$ very little has been proved. It is straightforward to show that

$$
\limsup _{n \rightarrow \infty} \frac{1}{n} \log p_{n}(K) \geq \kappa_{0} .
$$

But it is generally believed that

$$
\lim _{n \rightarrow \infty} \frac{1}{n} \log p_{n}(K)=\kappa_{0},
$$

and moreover

$$
p_{n}(K) \sim A_{K} n^{\alpha-3+f(K)} \mu_{0}^{n}
$$

where $\alpha$ is the same exponent as for all polygons, $A_{K}$ is a constant and $f(K)$ is the number of prime knot factors in the knot decomposition of $K$.

## Knotted polygons cont'd

For $K \neq 0_{1}$ very little has been proved. It is straightforward to show that

$$
\limsup _{n \rightarrow \infty} \frac{1}{n} \log p_{n}(K) \geq \kappa_{0}
$$

But it is generally believed that

$$
\lim _{n \rightarrow \infty} \frac{1}{n} \log p_{n}(K)=\kappa_{0}
$$

and moreover

$$
p_{n}(K) \sim A_{K} n^{\alpha-3+f(K)} \mu_{0}^{n}
$$

where $\alpha$ is the same exponent as for all polygons, $A_{K}$ is a constant and $f(K)$ is the number of prime knot factors in the knot decomposition of $K$.

This is essentially because the "knotted parts" of a long polygon are expected to be small:


## Polygons in lattice tubes

Let $\mathbb{T}_{L, M} \equiv \mathbb{T}$ be the $L \times M$ infinite tube of $\mathbb{Z}^{3}$ :

$$
\mathbb{T}_{L, M}=\{(x, y, z) \mid 0 \leq y \leq L, 0 \leq z \leq M\}
$$

and let $p_{\mathbb{T}, n}$ be the number of SAPs in $\mathbb{T}$, defined up to translation in the $x$-direction only.


## Polygons in lattice tubes

Let $\mathbb{T}_{L, M} \equiv \mathbb{T}$ be the $L \times M$ infinite tube of $\mathbb{Z}^{3}$ :

$$
\mathbb{T}_{L, M}=\{(x, y, z) \mid 0 \leq y \leq L, 0 \leq z \leq M\}
$$

and let $p_{\mathbb{T}, n}$ be the number of SAPs in $\mathbb{T}$, defined up to translation in the $x$-direction only.


In $\mathbb{T}$, polygons are characterised by a finite transfer matrix $\Rightarrow$ growth rates, exponents, etc. can be computed exactly (in theory).

## Theorem (Soteros 1998)

There exist constants $A_{\mathbb{T}}$ and $\kappa_{\mathbb{T}}=\log \mu_{\mathbb{T}}$ such that

$$
p_{\mathbb{T}, n} \sim A_{\mathbb{T}} \mu_{\mathbb{T}}^{n}
$$

$A_{\mathbb{T}}$ and $\mu_{\mathbb{T}}$ are algebraic numbers.

## Knotted polygons in tubes

If $L, M>0$ and $(L, M) \neq(1,1)$ then polygons in $\mathbb{T}$ can be knotted. (Only unknots in $1 \times 1$.) Define $p_{\mathbb{T}, n}(K)$ to count polygons of knot type $K$.

Similarly to $\mathbb{Z}^{3}$ :

## Theorem (Soteros 1998)

There exists $\kappa_{\mathbb{T}, 0}=\log \mu_{\mathbb{T}, 0}$ such that

$$
p_{\mathbb{T}, n}\left(0_{1}\right)=\exp \left\{\kappa_{\mathbb{T}, 0} n+o(n)\right\}=e^{o(n)} \mu_{\mathbb{T}, 0}^{n} .
$$

If $L, M>0$ and $(L, M) \neq(1,1)$ then

$$
\mu_{\mathbb{T}, 0}<\mu_{\mathbb{T}} .
$$

That is, in a 3-dimensional tube other than $1 \times 1$, the probably of a polygon being unknotted decays exponentially.

## Knotted polygons in tubes cont'd

As with $\mathbb{Z}^{3}$, it is easy to show that

$$
\limsup _{n \rightarrow \infty} \frac{1}{n} \log p_{\mathbb{T}, n}(K) \geq \kappa_{\mathbb{T}, 0} .
$$

## Knotted polygons in tubes cont'd

As with $\mathbb{Z}^{3}$, it is easy to show that

$$
\limsup _{n \rightarrow \infty} \frac{1}{n} \log p_{\mathbb{T}, n}(K) \geq \kappa_{\mathbb{T}, 0} .
$$

And it is believed that

$$
p_{\mathbb{T}, n}(K) \sim A_{\mathbb{T}, K} n^{f(K)} \mu_{\mathbb{T}, 0}^{n} .
$$

## Knotted polygons in tubes cont'd

As with $\mathbb{Z}^{3}$, it is easy to show that

$$
\limsup _{n \rightarrow \infty} \frac{1}{n} \log p_{\mathbb{T}, n}(K) \geq \kappa_{\mathbb{T}, 0} .
$$

And it is believed that

$$
p_{\mathbb{T}, n}(K) \sim A_{\mathbb{T}, K} n^{f(K)} \mu_{\mathbb{T}, 0}^{n} .
$$

This is again because the "knotted parts" of a long polygon are expected to be small:


## Counting by span

The $x$-span of a polygon is the maximal difference in $x$-coordinates between any two vertices.

## Counting by span

The $x$-span of a polygon is the maximal difference in $x$-coordinates between any two vertices.
In $\mathbb{T}$, we can count polygons by $x$-span instead of length: $q_{\mathbb{T}, s}$ and $q_{\mathbb{T}, s}(K)$. The same theorems hold:

## Theorem (Atapour 2008)

There exist constants $B_{\mathbb{T}}$ and $\chi_{\mathbb{T}}=\log \nu_{\mathbb{T}}$ such that

$$
q_{\mathbb{T}, s} \sim B_{\mathbb{T}} \nu_{\mathbb{T}}^{s}
$$

and constants $\chi_{\mathbb{T}, 0}=\log \nu_{\mathbb{T}, 0}$ such that

$$
q_{\mathbb{T}, s}\left(0_{1}\right)=\exp \left\{\chi_{\mathbb{T}, 0} s+o(s)\right\}=e^{o(s)} \nu_{\mathbb{T}, 0}^{s} .
$$

If $L, M>0$ and $(L, M) \neq(1,1)$ then

$$
\nu_{\mathbb{T}, 0}<\nu_{\mathbb{T}}
$$

## Counting by span

The $x$-span of a polygon is the maximal difference in $x$-coordinates between any two vertices.
In $\mathbb{T}$, we can count polygons by $x$-span instead of length: $q_{\mathbb{T}, s}$ and $q_{\mathbb{T}, s}(K)$. The same theorems hold:

## Theorem (Atapour 2008)

There exist constants $B_{\mathbb{T}}$ and $\chi_{\mathbb{T}}=\log \nu_{\mathbb{T}}$ such that

$$
q_{\mathbb{T}, s} \sim B_{\mathbb{T}} \nu_{\mathbb{T}}^{s}
$$

and constants $\chi_{\mathbb{T}, 0}=\log \nu_{\mathbb{T}, 0}$ such that

$$
q_{\mathbb{T}, s}\left(0_{1}\right)=\exp \left\{\chi_{\mathbb{T}, 0} s+o(s)\right\}=e^{o(s)} \nu_{\mathbb{T}, 0}^{s} .
$$

If $L, M>0$ and $(L, M) \neq(1,1)$ then

$$
\nu_{\mathbb{T}, 0}<\nu_{\mathbb{T}}
$$

Again expect that

$$
q_{\mathbb{T}, s}(K) \sim A_{\mathbb{T}, K} s^{f(K)} \nu_{\mathbb{T}, 0}^{s}
$$

$B_{\mathbb{T}}$ and $\nu_{\mathbb{T}}$ are algebraic numbers.

## Computing probabilities

Polygons in $\mathbb{T}$ are characterised by a finite transfer matrix, but polygons of fixed knot type (including unknots) are not. This is because the state depends on the previous entanglement of the "strands", and there are infinitely many possibilities:


## Computing probabilities

Polygons in $\mathbb{T}$ are characterised by a finite transfer matrix, but polygons of fixed knot type (including unknots) are not. This is because the state depends on the previous entanglement of the "strands", and there are infinitely many possibilities:


In general can only estimate the behaviour of $p_{\mathbb{T}, n}(K)$ and $q_{\mathbb{T}, s}(K)$.

## Computing probabilities

Polygons in $\mathbb{T}$ are characterised by a finite transfer matrix, but polygons of fixed knot type (including unknots) are not. This is because the state depends on the previous entanglement of the "strands", and there are infinitely many possibilities:


In general can only estimate the behaviour of $p_{\mathbb{T}, n}(K)$ and $q_{\mathbb{T}, s}(K)$.
Analysing series data is hopeless because knots are not common until very big lengths / spans.

## Computing probabilities

Polygons in $\mathbb{T}$ are characterised by a finite transfer matrix, but polygons of fixed knot type (including unknots) are not. This is because the state depends on the previous entanglement of the "strands", and there are infinitely many possibilities:


In general can only estimate the behaviour of $p_{\mathbb{T}, n}(K)$ and $q_{\mathbb{T}, s}(K)$.
Analysing series data is hopeless because knots are not common until very big lengths / spans.
Must use Monte Carlo methods!

## Transfer matrix Monte Carlo method

A 1-pattern is any configuration of vertices and (half)-edges between $x=k \pm \frac{1}{2}$ which can form part of a polygon, together with a pairing of the open half-edges on the left.


## Transfer matrix Monte Carlo method

A 1-pattern is any configuration of vertices and (half)-edges between $x=k \pm \frac{1}{2}$ which can form part of a polygon, together with a pairing of the open half-edges on the left.


A 1-pattern occurring at the very left end of a polygon is a starting 1-pattern (set $\mathcal{A}$ ); at the very right is an ending 1-pattern (set $\mathcal{B}$ ); otherwise it is internal (set $\mathcal{I}$ ).

## Transfer matrix Monte Carlo method

A 1-pattern is any configuration of vertices and (half)-edges between $x=k \pm \frac{1}{2}$ which can form part of a polygon, together with a pairing of the open half-edges on the left.


A 1-pattern occurring at the very left end of a polygon is a starting 1-pattern (set $\mathcal{A}$ ); at the very right is an ending 1 -pattern (set $\mathcal{B}$ ); otherwise it is internal (set $\mathcal{I}$ ).

Let $\mathbf{M}$ be the transfer matrix for internal 1-patterns, ie. $\mathbf{M}_{i j}=1$ if $i$ can be immediately followed by $j$ and 0 otherwise. Since $\mathbf{M}$ is irreducible and aperiodic, it has a unique dominant eigenvalue $\lambda \in \mathbb{R}^{>0}$ with right eigenvector $\zeta$.

## Transfer matrix Monte Carlo method

A 1-pattern is any configuration of vertices and (half)-edges between $x=k \pm \frac{1}{2}$ which can form part of a polygon, together with a pairing of the open half-edges on the left.


A 1-pattern occurring at the very left end of a polygon is a starting 1-pattern (set $\mathcal{A}$ ); at the very right is an ending 1-pattern (set $\mathcal{B}$ ); otherwise it is internal (set $\mathcal{I}$ ).

Let $\mathbf{M}$ be the transfer matrix for internal 1-patterns, ie. $\mathbf{M}_{i j}=1$ if $i$ can be immediately followed by $j$ and 0 otherwise. Since $\mathbf{M}$ is irreducible and aperiodic, it has a unique dominant eigenvalue $\lambda \in \mathbb{R}^{>0}$ with right eigenvector $\zeta$.

## Lemma (adapted from Alm \& Janson 1990)

If $i$ and $j$ are internal 1-patterns such that $j$ can follow $i$, let $p_{i j}^{\text {sp }}(s)$ be the probability that an occurrence of $i$ in a uniformly random polygon of span $s$ is followed by $j$. Then as $s \rightarrow \infty$,

$$
p_{i j}^{\mathrm{sp}}(s) \rightarrow p_{i j}^{\mathrm{sp}}=\lambda^{-1} \frac{\zeta_{j}}{\zeta_{i}}
$$

## Transfer matrix Monte Carlo method cont'd

Can likewise define matrices $\mathbf{A}$ with rows indexed by $\mathcal{A}$ and columns indexed by $\mathcal{I}$, and $\mathbf{B}$ indexed by $\mathcal{I}$ and $\mathcal{B}$.

## Transfer matrix Monte Carlo method cont'd

Can likewise define matrices $\mathbf{A}$ with rows indexed by $\mathcal{A}$ and columns indexed by $\mathcal{I}$, and $\mathbf{B}$ indexed by $\mathcal{I}$ and $\mathcal{B}$.

We then use an algorithm for generating a polygon $\pi=\pi_{0} \pi_{1} \cdots \pi_{s}$ uniformly at random. For $a \in \mathcal{A}$ and $i \in \mathcal{I}$, define

$$
\mathcal{F}_{\text {start }}(a)=\left\{j \in \mathcal{I}: \mathbf{A}_{a j}=1\right\} \quad \text { and } \quad \mathcal{F}_{\text {end }}(i)=\left\{b \in \mathcal{B}: \mathbf{B}_{i b}=1\right\}
$$

and then

$$
t_{1}(a)=\sum_{j \in \mathcal{F}_{\text {start }}(a)} \zeta_{j} \quad \text { and } \quad t_{s}(i)=\frac{\left|\mathcal{F}_{\text {end }}(i)\right|}{\zeta_{i}}
$$

(1) Select $\pi_{0}$ uniformly at random from $\mathcal{S}$.
(2) With probability $r_{1}\left(\pi_{0}\right)$ (see below) reject the sample and start from (1) again. Else select $\pi_{1}$ from $\mathcal{F}_{\text {start }}\left(\pi_{0}\right)$ with probability proportional to $\zeta_{\pi_{1}}$.
(3) For $k=2,3, \ldots, s-1$, choose $\pi_{k}$ with probability $p_{\pi_{k-1}, \pi_{k}}^{\mathrm{sp}}$.
(9) With probability $r_{s}\left(\pi_{s-1}\right)$ (see below), reject the sample and start from (1) again. Otherwise select $\pi_{s}$ uniformly from $\mathcal{F}_{\text {end }}\left(\pi_{s-1}\right)$.
The rejection probabilities are chosen to make the sampling uniformly random:

$$
r_{1}\left(\pi_{0}\right)=1-\frac{t_{1}\left(\pi_{0}\right)}{\max _{a \in \mathcal{S}}\left\{t_{1}(a)\right\}} \quad \text { and } \quad r_{s}\left(\pi_{s-1}\right)=1-\frac{t_{s}\left(\pi_{s-1}\right)}{\max _{j \in \mathcal{I}}\left\{t_{s}(j)\right\}}
$$

## Computing probabilities cont'd

For polygons of thousands of edges, determining knot type directly is very difficult!

Computing probabilities cont'd
For polygons of thousands of edges, determining knot type directly is very difficult!
But we can slice them up at " 2 -sections" - $x$-coordinates with only two edges crossing - and close up the pieces to form a sequence of smaller polygons.


## Computing probabilities cont'd

For polygons of thousands of edges, determining knot type directly is very difficult!
But we can slice them up at " 2 -sections" - $x$-coordinates with only two edges crossing - and close up the pieces to form a sequence of smaller polygons.


Then we compute the knot type of each, and the overall knot type is the connect-sum of the parts. Since long polygons have a positive density of 2 -sections (guaranteed by pattern theorems), the pieces are almost always very small.

## Fixed-length vs. fixed-span vs. Hamiltonian

The transfer matrix Monte Carlo method described above was for fixed-span s, but it can be adapted to fixed-length as well.

## Fixed-length vs. fixed-span vs. Hamiltonian

The transfer matrix Monte Carlo method described above was for fixed-span s, but it can be adapted to fixed-length as well.

In $\mathbb{T}$ it may seem that the choice is arbitrary, but in fact there is a big difference.

## Fixed-length vs. fixed-span vs. Hamiltonian

The transfer matrix Monte Carlo method described above was for fixed-span s, but it can be adapted to fixed-length as well.

In $\mathbb{T}$ it may seem that the choice is arbitrary, but in fact there is a big difference.
Fixed-span polygons are much denser. On average (edges per unit span):

|  | fixed-length | fixed-span |
| :---: | :---: | :---: |
| $2 \times 1$ | 3.6214 | 4.8865 |
| $3 \times 1$ | 4.1052 | 6.5244 |

## Fixed-length vs. fixed-span vs. Hamiltonian

The transfer matrix Monte Carlo method described above was for fixed-span s, but it can be adapted to fixed-length as well.

In $\mathbb{T}$ it may seem that the choice is arbitrary, but in fact there is a big difference.
Fixed-span polygons are much denser. On average (edges per unit span):

|  | fixed-length | fixed-span |
| :---: | :---: | :---: |
| $2 \times 1$ | 3.6214 | 4.8865 |
| $3 \times 1$ | 4.1052 | 6.5244 |

Denser $\Rightarrow$ more "entanglement" $\Rightarrow$ higher knotting probability. Here I will focus on fixed-span.

## Fixed-length vs. fixed-span vs. Hamiltonian

The transfer matrix Monte Carlo method described above was for fixed-span s, but it can be adapted to fixed-length as well.

In $\mathbb{T}$ it may seem that the choice is arbitrary, but in fact there is a big difference.
Fixed-span polygons are much denser. On average (edges per unit span):

|  | fixed-length | fixed-span |
| :---: | :---: | :---: |
| $2 \times 1$ | 3.6214 | 4.8865 |
| $3 \times 1$ | 4.1052 | 6.5244 |

Denser $\Rightarrow$ more "entanglement" $\Rightarrow$ higher knotting probability. Here I will focus on fixed-span.
The method also works for Hamiltonian polygons, which are as dense as possible ( 6 and 8 edges per unit span respectively), so we sample those too. (All the previous theorems about asymptotics also apply.)

## Fixed-length vs. fixed-span vs. Hamiltonian

The transfer matrix Monte Carlo method described above was for fixed-span s, but it can be adapted to fixed-length as well.

In $\mathbb{T}$ it may seem that the choice is arbitrary, but in fact there is a big difference.
Fixed-span polygons are much denser. On average (edges per unit span):

|  | fixed-length | fixed-span |
| :---: | :---: | :---: |
| $2 \times 1$ | 3.6214 | 4.8865 |
| $3 \times 1$ | 4.1052 | 6.5244 |

Denser $\Rightarrow$ more "entanglement" $\Rightarrow$ higher knotting probability. Here I will focus on fixed-span.
The method also works for Hamiltonian polygons, which are as dense as possible ( 6 and 8 edges per unit span respectively), so we sample those too. (All the previous theorems about asymptotics also apply.)

Define

$$
\begin{aligned}
\mathbb{P}_{\mathbb{T}, s}^{(\mathrm{sp})}(K) & =\text { probability of knot-type } K \text { among all knots of span } s \text { in } \mathbb{T} \\
\mathbb{P}_{\mathbb{T}, s}^{\mathrm{H}}(K) & =\text { probability of knot-type } K \text { among all Hamiltonian knots of span } s \text { in } \mathbb{T} \\
\mathbb{P}_{\mathbb{T}, s}^{*}(K) & =\text { one of the above }
\end{aligned}
$$

## Results: Growth rates

Since $\mathbb{P}_{\mathbb{T}, s}^{*}\left(0_{1}\right)$ decays exponentially, plot the log and take a linear best fit.

## Results: Growth rates

Since $\mathbb{P}_{\mathbb{T}, s}^{*}\left(O_{1}\right)$ decays exponentially, plot the log and take a linear best fit.
$2 \times 1$ tube:


So

$$
\chi_{\mathbb{T}, 0}-\chi_{\mathbb{T}}=-8.12115 \times 10^{-7} \quad \Rightarrow \quad \frac{\nu_{\mathbb{T}, 0}}{\nu_{\mathbb{T}}}=0.99999919
$$

## Results: Growth rates

Since $\mathbb{P}_{\mathbb{T}, \mathrm{s}}^{*}\left(0_{1}\right)$ decays exponentially, plot the log and take a linear best fit.
$3 \times 1$ tube:


So

$$
\chi_{\mathbb{T}, 0}-\chi_{\mathbb{T}}=-0.000140089 \quad \Rightarrow \quad \frac{\nu_{\mathbb{T}, 0}}{\nu_{\mathbb{T}}}=0.99986
$$

## Results: Growth rates

Since $\mathbb{P}_{\mathbb{T}, s}^{*}\left(0_{1}\right)$ decays exponentially, plot the log and take a linear best fit.
$2 \times 1$ tube Hamiltonian:


So

$$
\chi_{\mathbb{T}, 0}^{\mathrm{H}}-\chi_{\mathbb{T}}^{\mathrm{H}}=-2.86675 \times 10^{-5} \Rightarrow \frac{\nu_{\mathbb{T}, 0}^{\mathrm{H}}}{\nu_{\mathbb{T}}^{\mathrm{H}}}=0.999971
$$

## Results: Growth rates

Since $\mathbb{P}_{\mathbb{T}, s}^{*}\left(0_{1}\right)$ decays exponentially, plot the log and take a linear best fit.
$3 \times 1$ tube Hamiltonian:


So

$$
\chi_{\mathbb{T}, 0}^{\mathrm{H}}-\chi_{\mathbb{T}}^{\mathrm{H}}=-0.00713521 \quad \Rightarrow \quad \frac{\nu_{\mathbb{T}, 0}^{\mathrm{H}}}{\nu_{\mathbb{T}}^{\mathrm{H}}}=0.99289
$$

## Results: Exponents

No evidence of logarithmic corrections: strongly suggests that

$$
q_{\mathbb{T}, s}\left(0_{1}\right) \sim B_{\mathbb{T}, 0}\left(\nu_{\mathbb{T}, 0}\right)^{s} \quad \text { and } \quad q_{\mathbb{T}, s}^{\mathrm{H}}\left(0_{1}\right) \sim B_{\mathbb{T}, 0}^{\mathrm{H}}\left(\nu_{\mathbb{T}, 0}^{\mathrm{H}}\right)^{s} .
$$

## Results: Exponents

No evidence of logarithmic corrections: strongly suggests that

$$
q_{\mathbb{T}, s}\left(0_{1}\right) \sim B_{\mathbb{T}, 0}\left(\nu_{\mathbb{T}, 0}\right)^{s} \quad \text { and } \quad q_{\mathbb{T}, s}^{\mathrm{H}}\left(0_{1}\right) \sim B_{\mathbb{T}, 0}^{\mathrm{H}}\left(\nu_{\mathbb{T}, 0}^{\mathrm{H}}\right)^{s} .
$$

For other knot types, examine ratio $\mathbb{P}_{\mathbb{T}, s}^{*}(K) / \mathbb{P}_{\mathbb{T}, s}^{*}\left(0_{1}\right)$.

## Results: Exponents

No evidence of logarithmic corrections: strongly suggests that

$$
q_{\mathbb{T}, s}\left(0_{1}\right) \sim B_{\mathbb{T}, 0}\left(\nu_{\mathbb{T}, 0}\right)^{s} \quad \text { and } \quad q_{\mathbb{T}, s}^{\mathrm{H}}\left(0_{1}\right) \sim B_{\mathbb{T}, 0}^{\mathrm{H}}\left(\nu_{\mathbb{T}, 0}^{\mathrm{H}}\right)^{s} .
$$

For other knot types, examine ratio $\mathbb{P}_{\mathbb{T}, s}^{*}(K) / \mathbb{P}_{\mathbb{T}, s}^{*}\left(0_{1}\right)$.
$2 \times 1$ tube:


## Results: Exponents

No evidence of logarithmic corrections: strongly suggests that

$$
q_{\mathbb{T}, s}\left(0_{1}\right) \sim B_{\mathbb{T}, 0}\left(\nu_{\mathbb{T}, 0}\right)^{s} \quad \text { and } \quad q_{\mathbb{T}, s}^{\mathrm{H}}\left(0_{1}\right) \sim B_{\mathbb{T}, 0}^{\mathrm{H}}\left(\nu_{\mathbb{T}, 0}^{\mathrm{H}}\right)^{s} .
$$

For other knot types, examine ratio $\mathbb{P}_{\mathbb{T}, s}^{*}(K) / \mathbb{P}_{\mathbb{T}, s}^{*}\left(0_{1}\right)$.
$3 \times 1$ tube:


## Results: Exponents

No evidence of logarithmic corrections: strongly suggests that

$$
q_{\mathbb{T}, s}\left(0_{1}\right) \sim B_{\mathbb{T}, 0}\left(\nu_{\mathbb{T}, 0}\right)^{s} \quad \text { and } \quad q_{\mathbb{T}, s}^{\mathrm{H}}\left(0_{1}\right) \sim B_{\mathbb{T}, 0}^{\mathrm{H}}\left(\nu_{\mathbb{T}, 0}^{\mathrm{H}}\right)^{s} .
$$

For other knot types, examine ratio $\mathbb{P}_{\mathbb{T}, s}^{*}(K) / \mathbb{P}_{\mathbb{T}, s}^{*}\left(0_{1}\right)$.
$2 \times 1$ tube Hamiltonian:


## Results: Exponents

No evidence of logarithmic corrections: strongly suggests that

$$
q_{\mathbb{T}, s}\left(0_{1}\right) \sim B_{\mathbb{T}, 0}\left(\nu_{\mathbb{T}, 0}\right)^{s} \quad \text { and } \quad q_{\mathbb{T}, s}^{\mathrm{H}}\left(0_{1}\right) \sim B_{\mathbb{T}, 0}^{\mathrm{H}}\left(\nu_{\mathbb{T}, 0}^{\mathrm{H}}\right)^{s} .
$$

For other knot types, examine ratio $\mathbb{P}_{\mathbb{T}, s}^{*}(K) / \mathbb{P}_{\mathbb{T}, s}^{*}\left(0_{1}\right)$.
$3 \times 1$ tube Hamiltonian:


## Results: Exponents cont'd

For composite knots, take the ratio $\mathbb{P}_{\mathbb{T}, s}^{*}(K) / \mathbb{P}_{\mathbb{T}, s}^{*}\left(0_{1}\right)$ and then look at log-log plot.

## Results: Exponents cont'd

For composite knots, take the ratio $\mathbb{P}_{\mathbb{T}, s}^{*}(K) / \mathbb{P}_{\mathbb{T}, s}^{*}\left(0_{1}\right)$ and then look at log-log plot.
$3 \times 1$ tube:


## Results: Exponents cont'd

For composite knots, take the ratio $\mathbb{P}_{\mathbb{T}, s}^{*}(K) / \mathbb{P}_{\mathbb{T}, s}^{*}\left(0_{1}\right)$ and then look at log-log plot.
$3 \times 1$ tube Hamiltonian:


## Results: Exponents cont'd

For composite knots, take the ratio $\mathbb{P}_{\mathbb{T}, s}^{*}(K) / \mathbb{P}_{\mathbb{T}, s}^{*}\left(0_{1}\right)$ and then look at log-log plot.
$3 \times 1$ tube Hamiltonian:


Altogether, strongly implies

$$
q_{\mathbb{T}, s}(K) \sim B_{\mathbb{T}, K} s^{f(K)}\left(\nu_{\mathbb{T}, 0}\right)^{s} \quad \text { and } \quad q_{\mathbb{T}, s}^{\mathrm{H}}(K) \sim B_{\mathbb{T}, K}^{\mathrm{H}} s^{f(K)}\left(\nu_{\mathbb{T}, 0}^{\mathrm{H}}\right)^{s} .
$$

## Results: Probability maxima

$\mathbb{P}_{\mathbb{T}, s}^{*}(K)$ decays exponentially for any fixed knot type $K$ or set of knot types. But if $K \neq 0_{1}$ then $\mathbb{P}_{\mathbb{T}, s}^{*}(K)$ initially increases, reaches some maximum, then decreases to 0 .

## Results: Probability maxima

$\mathbb{P}_{\mathbb{T}, s}^{*}(K)$ decays exponentially for any fixed knot type $K$ or set of knot types. But if $K \neq 0_{1}$ then $\mathbb{P}_{\mathbb{T}, s}^{*, s}(K)$ initially increases, reaches some maximum, then decreases to 0 .

If

$$
\mathbb{P}_{\mathbb{T}, s}^{*}(K) \sim C_{\mathbb{T}, K}^{*} s^{f(K)}\left(\frac{\nu_{\mathbb{T}, 0}^{*}}{\nu_{\mathbb{T}}^{*}}\right)^{s}
$$

then the maximum should be at roughly

$$
s^{*} \approx M_{\mathbb{T}}^{*}(K)=\frac{f(K)}{\chi_{\mathbb{T}}^{*}-\chi_{\mathbb{T}, 0}^{*}}= \begin{cases}\left(1.23 \times 10^{6}\right) f(K) & 2 \times 1 \text { tube } \\ 7140 f(K) & 3 \times 1 \text { tube } \\ 34900 f(K) & 2 \times 1 \text { tube Hamiltonian } \\ 1400 f(K) & 3 \times 1 \text { tube Hamiltonian }\end{cases}
$$

## Results: Probability maxima cont'd

Prime knots in the $3 \times 1$ tube:


## Results: Probability maxima cont'd

Prime Hamiltonian knots in the $3 \times 1$ tube:


## Results: Probability maxima cont'd

2-factor knots in the $3 \times 1$ tube:


## Results: Probability maxima cont'd

2-factor Hamiltonian knots in the $3 \times 1$ tube:


## Results: Probability maxima cont'd

Multi-factor knots in the $3 \times 1$ tube:


## Results: Probability maxima cont'd

Multi-factor Hamiltonian knots in the $3 \times 1$ tube:


## Closing comments

- Believe all the same to be true for fixed-length, but knots are very very rare in narrow tubes.


## Closing comments

- Believe all the same to be true for fixed-length, but knots are very very rare in narrow tubes.
- Transfer matrix method very memory intensive. Can (just barely) do Hamiltonian polygons in $2 \times 2$ and $4 \times 1$. For bigger tubes, other methods like PERM, Wang-Landau, etc may be needed.


## Closing comments

- Believe all the same to be true for fixed-length, but knots are very very rare in narrow tubes.
- Transfer matrix method very memory intensive. Can (just barely) do Hamiltonian polygons in $2 \times 2$ and $4 \times 1$. For bigger tubes, other methods like PERM, Wang-Landau, etc may be needed.
- In the bulk, numerical evidence [Janse van Rensburg \& Rechnitzer 2011] that the amplitude ratios $A_{K_{1}} / A_{K_{2}}$ are universal. Does not appear to be the case in tubes.


## Closing comments

- Believe all the same to be true for fixed-length, but knots are very very rare in narrow tubes.
- Transfer matrix method very memory intensive. Can (just barely) do Hamiltonian polygons in $2 \times 2$ and $4 \times 1$. For bigger tubes, other methods like PERM, Wang-Landau, etc may be needed.
- In the bulk, numerical evidence [Janse van Rensburg \& Rechnitzer 2011] that the amplitude ratios $A_{K_{1}} / A_{K_{2}}$ are universal. Does not appear to be the case in tubes.
- In $2 \times 1$ only, a stronger result can be proved [Atapour, NRB, Eng, Ishihara, Shimokawa, Soteros \& Vazquez, in preparation]:

$$
p_{\mathbb{T}, n}(K) \sim D_{\mathbb{T}, K} n^{f(K)} p_{n, \mathbb{T}}\left(0_{1}\right)
$$

if all of $K$ 's factors have unknotting number one.

## Closing comments

- Believe all the same to be true for fixed-length, but knots are very very rare in narrow tubes.
- Transfer matrix method very memory intensive. Can (just barely) do Hamiltonian polygons in $2 \times 2$ and $4 \times 1$. For bigger tubes, other methods like PERM, Wang-Landau, etc may be needed.
- In the bulk, numerical evidence [Janse van Rensburg \& Rechnitzer 2011] that the amplitude ratios $A_{K_{1}} / A_{K_{2}}$ are universal. Does not appear to be the case in tubes.
- In $2 \times 1$ only, a stronger result can be proved [Atapour, NRB, Eng, Ishihara, Shimokawa, Soteros \& Vazquez, in preparation]:

$$
p_{\mathbb{T}, n}(K) \sim D_{\mathbb{T}, K} n^{f(K)} p_{n, \mathbb{T}}\left(0_{1}\right)
$$

if all of $K$ 's factors have unknotting number one.

> NRB, Eng \& Soteros
> Knotting statistics for polygons in lattice tubes J. Phys. A: Math. Theor. 52 (2019), 144003 .

Thank you!

