

"Toward a Construction of Non-classical String Solutions"

Savdeep Sethi

University of Chicago

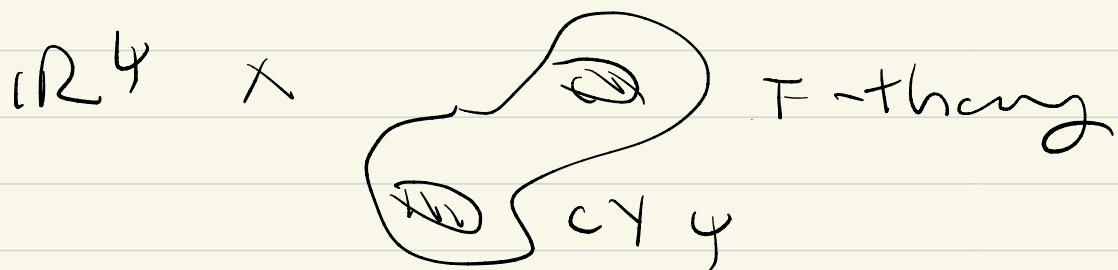
Outline:



I No-go results

II Status of the IR landscape

III Toward a construction of
Non-classical string solns.



AdS ? dS ? Swampland vs landscape

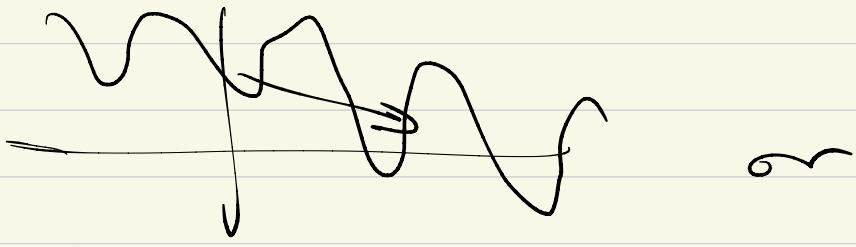
$\text{D}2^4 \times (\text{something})$ classical string
solns -

CFT

- Not supersymmetric
- $V_{\text{loop}} \sim \int_S 2\pi \text{loop}$
- $\Lambda \rightarrow \Lambda + \epsilon \Lambda$
- Scale separated AdS ?
- de Sitter ?
- Including string loops but not string non-perturbation.

I No-go results

- Dark energy $M_{\text{obs}} \sim 10^{-?} M_{\text{pl}}$
 $\gg 10^6$ of the obs critical energy.



time-dependent
dark energy

Landscape

Refined de Sitter
conjecture

$$M_p \frac{|V'|}{\sqrt{V}} \equiv \lambda \gtrsim o(1)$$

$$\vdash - M_p^2 \frac{V''}{\sqrt{V}} \equiv c \gtrsim o(1)$$

$$(i) V(\phi) = R c^{-\lambda} \phi$$

$$(ii) V(\phi) = R \cos(c\phi)$$

$c \approx 0.16$ at 68% CL. Numbers

to explain,

Data set	c	λ_{eff}	$ \Delta\phi [M_P]$
CMB	$c < 2.3 (3.1)$	$\lambda_{\text{eff}} < 1.4 (2.2)$	$ \Delta\phi < 0.51 (0.66)$
CMB + SN	$c < 0.25 (1.4)$	$\lambda_{\text{eff}} < 0.40 (0.71)$	$ \Delta\phi < 0.11 (0.19)$
CMB + H_0	$c < 0.17 (0.84)$	$\lambda_{\text{eff}} < 0.31 (0.58)$	$ \Delta\phi < 0.09 (0.16)$
ALL	$c < 0.16 (0.73)$	$\lambda_{\text{eff}} < 0.29 (0.53)$	$ \Delta\phi < 0.08 (0.15)$

No- y_0 results

Appearance 1 - Spacetime

Classic no- y_0 (Gibbons)

SUGRA $D = 10, 11$

$$S = \frac{1}{2k^2} \int \sqrt{g} R + \dots$$

$$R_{mn} \sim \left(T_{mn} - \frac{1}{\theta^2} T g_{mn} \right)$$

T obeys SEC (strong energy condition)

$$R_{\mu\nu} \geq 0 \quad (R_{\mu\nu} u^\mu u^\nu \geq 0)$$

$R_{\mu\nu} u^\mu u^\nu \geq 0$

SEC violation

Hilfstruktur $(g_{\mu\nu}, \phi, H, F)$

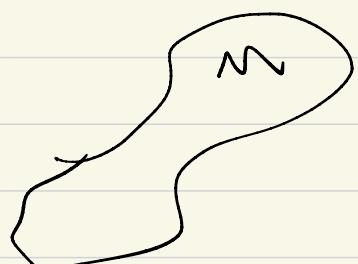
$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{g} e^{-2\phi} [R +$$

$$\frac{1}{2} (\partial\phi)^2 - \frac{1}{2} |H|^2 - \frac{\alpha'}{4} (\text{tr} F)^2$$

$$+ \dots [d^4]$$

$$R_{mn} = \frac{1}{2} \nabla_m \phi \nabla_n \phi + \dots$$

$$R_{\infty} = \frac{1}{2} (\dot{\phi})^2 + \dots \geq 0.$$



$$ds^2 = w^2(y) (g_{\mu\nu} dx^\mu dx^\nu + g_{ij}(y) dy^i dy^j)$$

$$R_{\infty}^{(D=10)} = R_{\infty}^{(D)} + \frac{1}{8w^4} \nabla^2 w^2 \geq 0$$

$$\Rightarrow R_{\infty}^{(D)} \geq 0.$$

Same argument rules out

$(n^4 \times CY_3)$ w/ flux.

Branes / anti-branes - conventional
stress-energy - don't work.

Orientifolds :



→ localized objects w/
negative tension

Violate SFC . Need higher derivative
interactions (4 derivatives)

Heterotic :

- (a) Leading higher deriv are known
- (b) Understand the w.c. description

$$\begin{aligned} \sigma(\omega^1) \\ \rightarrow -\frac{\omega^1}{4} \left(+_r |F|^2 - +_S |R_+|^2 \right) \\ + \sigma(\omega^1)^2 \end{aligned}$$

$$H = \partial R + \frac{\alpha'}{4} (\zeta(\omega_+ - \zeta(\Lambda))$$

(Bergshoeff, de Roo)

- Expansion term
- Dilaton term

$R_{\mu\nu}$ is not positive.

$\rightarrow \alpha' = 4$; $\alpha' = 4$ Planck mass is finite.

$$R_{\mu\nu} = \Lambda g_{\mu\nu} \quad (\text{const } \Lambda \text{ FLRW})$$

$$\Lambda = -\frac{\alpha'}{2\sqrt{1}} \left(\sum |J|^2 + o(\alpha'^2) \right)$$

$\Rightarrow \Lambda \leq 0$, Perfect square

No repulsion, no acceleration.

Approach 2 — worldsheet physics

Tree-level

$$S_{\text{het}} = \frac{1}{2k^*} \int d^{10}x \sqrt{-g} e^{-\psi} \sum R + \frac{\alpha'}{4} \left[R_+^2 + (\alpha')^3 R^4 + \dots \right]$$

Mink / AdS type. Also solutions

w flux (non-Kähler torional
backgrounds)

(Mansfield, Melnikov, Kutasov, V.S.)

$$D\mathcal{L}_4 \mapsto \text{isometry } SO(4,1)$$

Realized w symmetries
 $(\bar{\tau}^a, \bar{\bar{\tau}}^a)$

$$\Rightarrow \bar{\tau}\bar{\tau}^a + \partial\bar{\tau}^a = 0.$$

With state $x^a \rightarrow ix^a$, Typically
can't do this because of flux.

Only flux is $H_2 \in \Omega_{\mathbb{R}^2}$ is
"problematic".

$$\Rightarrow \bar{\mathcal{J}} \bar{\mathcal{J}}^a = \delta^{\bar{a}}_a \text{ KM symmetric}$$

$$\langle \mathcal{J} \mathcal{J} \rangle \sim \frac{b}{z^2} \quad \langle \bar{\mathcal{J}} \bar{\mathcal{J}} \rangle \sim \frac{\bar{b}}{\bar{z}^2}$$

by conf inv.

[Not what happens in $\mathbb{R}^{3,1}$ where
st. gen. are not promoted to KM]

Gravity $b > 0, \bar{b} > 0$

Left should be super KM

$$\Rightarrow b \geq 4.$$

DS length scale $L^2 \sim k l_s^2$

Critical string theory

$$(c, \bar{c}) = (15, 26)$$

The SIKM theory was w/o

$$C_{SIKM} = \left(\frac{k-4}{k} + \frac{1}{2} \right) 10 \quad (\text{ds-4})$$

Leaves $C_{SIKM} + C_{\text{corr}} = 15$
 $C_{\text{corr}} > 0$.

Once $k > 21$, could $\angle 31^\circ$

\Rightarrow must be on $N=1$ minimal model

$$\text{e)} \quad C_{\text{corr}} = \frac{20}{k-1} = \frac{3}{2} \left(1 - \frac{\epsilon}{\beta(\beta+1)} \right)$$

Discrete solutions:

$$(k-4, \beta) \in \{(15, 12), (21, 6), (23, 4)\}$$

No microscopic de Jitter.

- dilation doesn't have to be
jittered
- can have touch zones
- no stable or unstable zones.

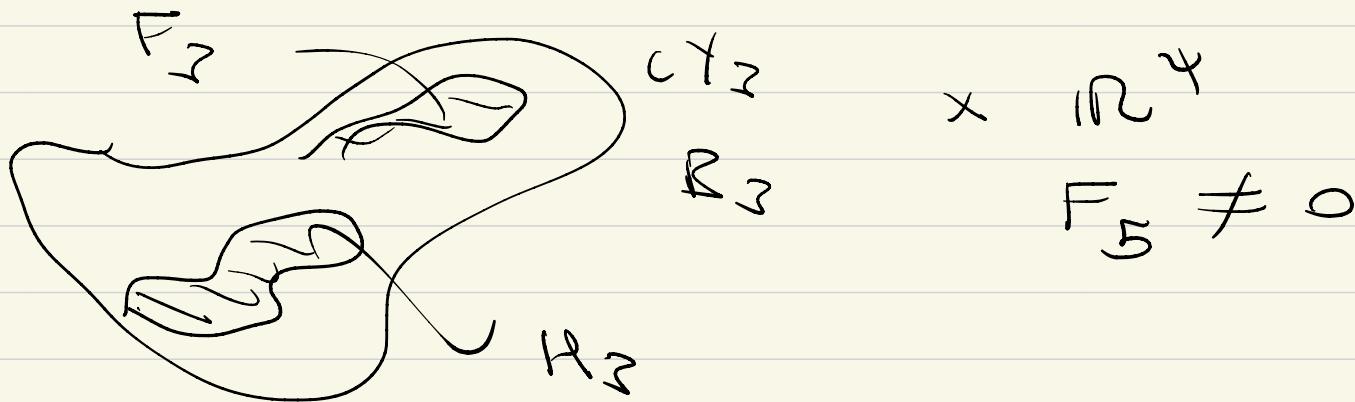
Misery:

- strong loops
- strong non-pert.

II. States of Mink IIB Landscape

In IIB or M-theory

\mathcal{P}_{IIB}



flux background. No sugra soln.
(period).

$N=1$ effective theory $D=4$

- characterized by K, w

$$\text{w/ potential } V = c^K \left(K^{ij} D_i w D_j \bar{w} - \beta(w) \right)$$

$$(*) \quad K = -3 \log(\varphi + \bar{\varphi})$$

\mathbb{R} no scale

$$\omega = \omega_0 = \int G_3 \wedge \omega_3$$

"Add" instanton $\omega = \omega_0 + A e^{-\rho}$

+ uplift.

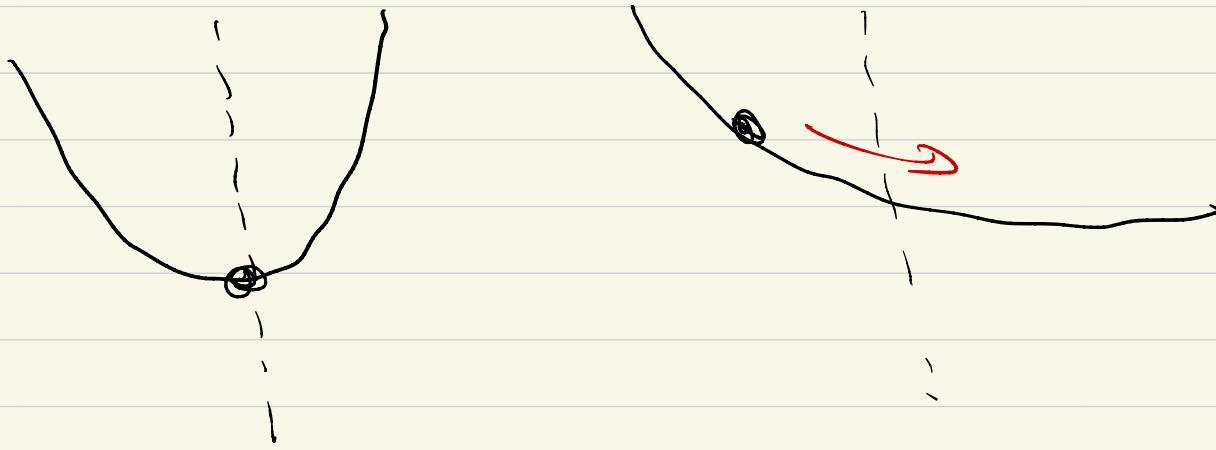
If $\omega_0 \neq 0 \Rightarrow$ sym broken
classically

and these backgrounds are not ext
of the Euclidean effective action in
any approx.

off-shell configuration

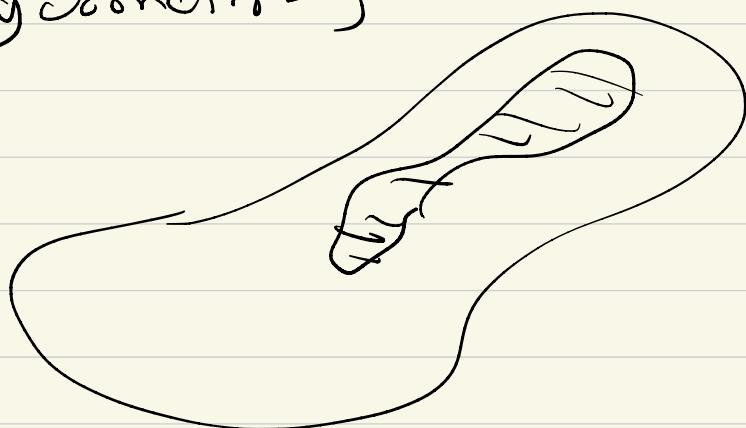
ω_0 is an obstruction to solving

The Question com.



Landscape $\mathcal{O}(L^{2000})$ CY 4

(geometric →)



CY 4

$$w_0 = 0$$

$\mathcal{O}(0)$ or no solution

per CY

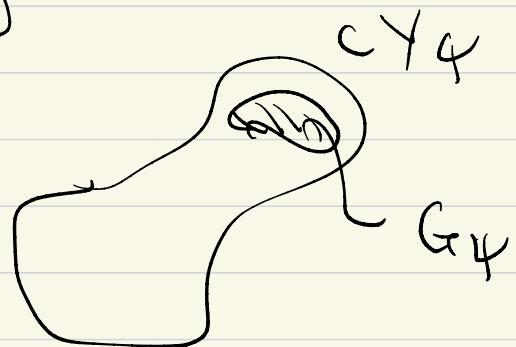
$$\begin{aligned} h^{21} &\rightarrow D_i w = 0 \\ +) \quad w &= 0. \end{aligned}$$

Intuitively \rightarrow Euclidean space
but this is complx.

More like a theory than YM.

Effective field theory?

M-theory: $R^3 \times$



$$S_2 = \int_{\text{hyp}} R - \int_{C \cap G \cap G}$$

No solutions.

$$S_8 = \int_{\text{hyp}} R^4 + \dots + \int_{C \cap X_8} C + \dots$$

↑
concrete flux
stress-energy $\sim \frac{\chi}{24}$ (lattice flux)

↑ tadpole

$$R^4 \rightarrow K_{\text{no-scale}} + \delta K$$

Cannot ignore.

No solution if $w \neq 0$.

III Non-classical stringy solutions

- de Sitter
- scale-separated AdS
- Non-susy AdS / CFT? (originally Fringuel-Kleban)

Idea: $V_{1-1-\beta} \sim \int d^{1-\beta} x$

balanced against a shift
in α and/or fluxes.

$D=10$ tachyon free non-susy
strings; 3 examples

$$V_{1-\beta\phi} > 0$$

"Strong island" (Dabholkar +
Honey)

Susy symmetric orbifolds w/

only the dilaton left.

T^4 $\mathcal{D}^{4,4}(A_4)$ asymmetric
orbifold by $2S$ left + shift
right
16 twists.

Type II on T^4/G

- no tachyons
- no moduli

w/ Zihni Baybars, Daniel Robbins

All attempts on Type have led to
tachyons.

Heterotic $T^4/(2S + \text{shift})$.

Wilson line moduli $E_8 \times \bar{E}_8$.

$\boxed{\quad} \rightarrow T^3$

$E_8 \times E_8$

\vdots
 $\boxed{\quad} \rightarrow T^3$

"triple" flat connection

asymmetric
orbifolds
25
26

$$\int_{T^3} CS = \frac{1}{5}, \frac{2}{5}, \frac{3}{5}, \frac{4}{5}$$

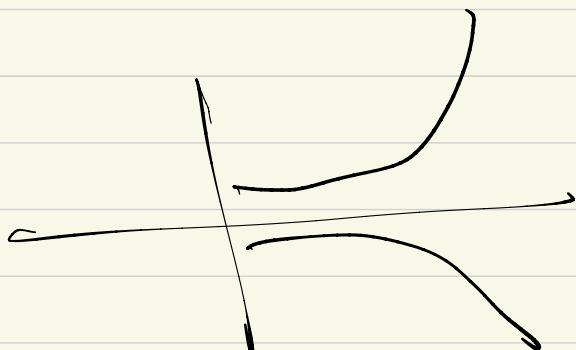
$$\int_{T^3} \omega = \frac{1}{6}, \frac{5}{6}$$

$(\frac{1}{5}, \frac{4}{5})$: complete rank reduction

$$P^{4,20} \rightarrow P^{4,4}$$

In progress.

$$M^6 \times T^4/G$$

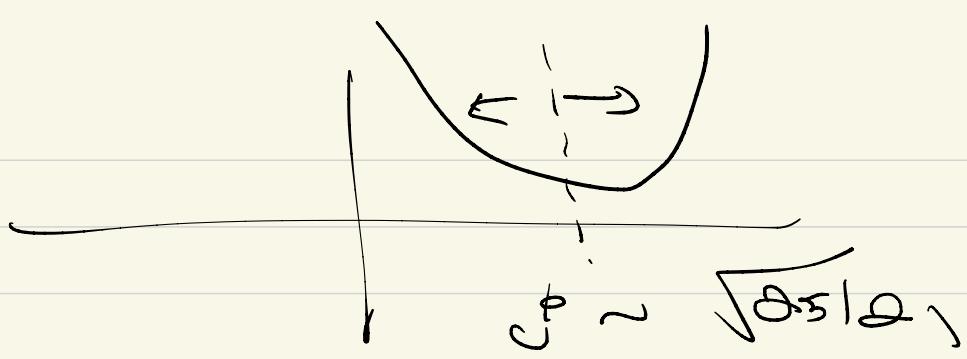


$$AdS_3 \times S^3 \times T^4/G$$

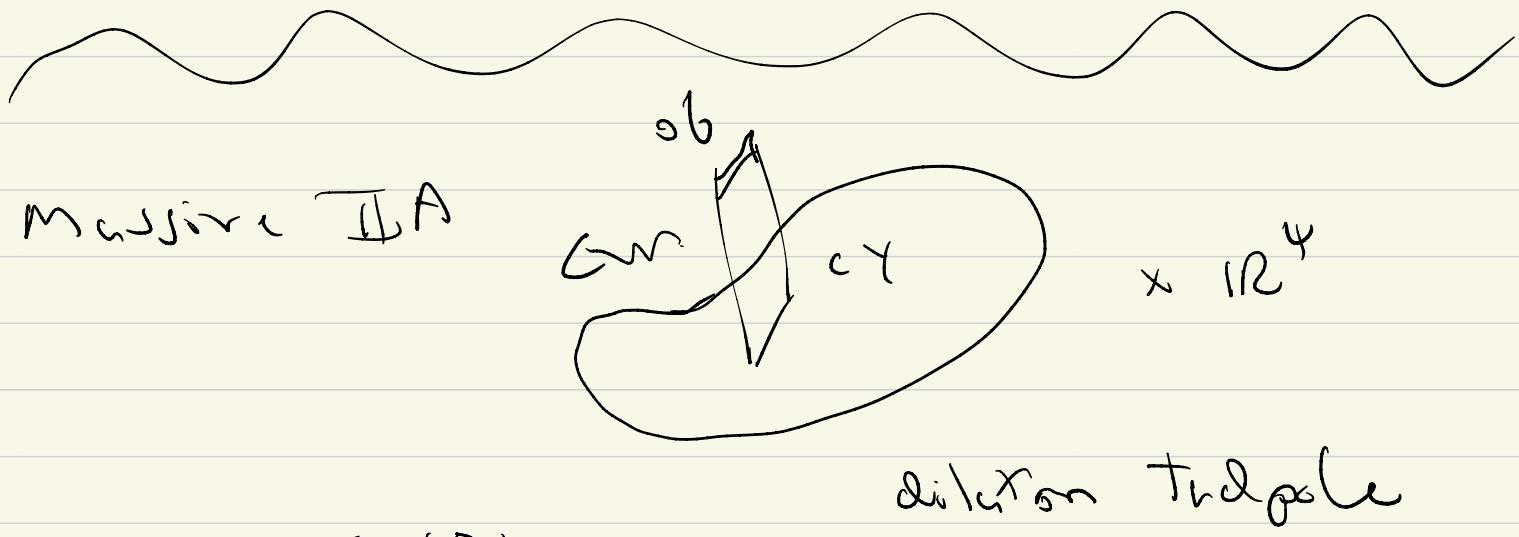
NS-flux

	Λ	Λ^\perp	Δr	E_t
Z_1	$\Gamma_{3,3} \oplus E_8 \oplus E_8$	\emptyset	0	-1
Z_2	$\Gamma_{3,3} \oplus D_4 \oplus D_4$	$D_4 \oplus D_4$	8	-1/2
Z_3	$\Gamma_{3,3} \oplus A_2 \oplus A_2$	$E_6 \oplus E_6$	12	-1/3
Z_4	$\Gamma_{3,3} \oplus A_1 \oplus A_1$	$E_7 \oplus E_7$	14	-1/4
Z_5	$\Gamma_{3,3}$	$E_8 \oplus E_8$	16	-1/5
Z_6	$\Gamma_{3,3}$	$E_8 \oplus E_8$	16	-1/6

Table 5: Lattices Λ , complements Λ^\perp , rank reduction Δr and zero-point energies in the twisted sector E_t for the Z_m asymmetric orbifolds corresponding to triples.



$$L^2 \sim l_s^2 \partial_5.$$



$$\sin(\gamma) \propto \sin(\beta)$$