

1 07w5095 The Mathematics of Knotting and  
2 Linking in Polymer Physics and Molecular  
3 Biology

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7 **1 Introduction**

8 The focus of this workshop was the mathematics associated with an array of cut-  
9 ting edge problems in polymer physics and molecular biology showing promise  
10 for immediate progress at the interfaces between mathematics and the physical  
11 and life sciences.

12 The first targeted area concerns the presence of knotting of DNA in living  
13 cells at a steady-state level lower than the thermodynamic equilibrium expected  
14 for a system in which inter-segmental passages within long DNA molecules oc-  
15 curs at random. Can one develop a systematic approach to understanding the  
16 wide range of potential topoisomerase mechanisms and their application in di-  
17 verse settings? Is there a selective topoisomerase mechanism by which knotting  
18 is kept below a topological equilibrium or are there specific constraining mecha-  
19 nisms promoting this relaxation of knots? The study of the characteristics of the  
20 equilibrium now include geometric, spatial, and topological aspects that may be  
21 implicated in these mechanisms as well as the characteristics of polymers, for  
22 example under theta conditions. Computational, experimental and theoretical  
23 aspects of this area were featured in many of the presentations and discussions.

24 The second targeted area concerns the mathematical, statistical, and compu-  
25 tational tools under development for the study of knotting and linking of  
26 open and closed macromolecules. One example is the collection of strategies  
27 developed to quantify and characterize the entanglement, e.g. knotting and  
28 linking, of open macromolecules which show promise for practical application of

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29 polymers. Another is the development of different methods for the selection of  
30 random equilateral polygons, with respect to the natural measure on the space  
31 of equilateral polygons. With efforts to quantify a wide range of new spatial fea-  
32 tures of these random equilateral polygons, greater care is necessary in order to  
33 demonstrate that the selection process is sufficient to provide statistically accu-  
34 rate estimations of critical quantities. Still another concerns the methods used  
35 to identify the topological type of the knotted polygons. Many of these methods  
36 are based on calculations of the Alexander polynomial or the more recent Jones  
37 and HOMFLY polynomials. While these have worked well to date, research  
38 questions are now moving into the range of 1500 edges (or Kuhn statistical  
39 segments) and, therefore, many thousands of crossings in a generic projection.  
40 Still another, distinct, computational thrust concerns efforts to achieve optimal  
41 spatial configurations when measured by the ropelength. With effort by several  
42 teams, this work faces challenging theoretical and computational obstacles.

43 The third focus is the application of the theory and methods above to the  
44 study of macromolecules in confined geometries, for example polymers between  
45 two parallel planes as in models of steric stabilization of dispersions or in DNA  
46 molecules contained in a capsid. Macromolecules so confined exhibit signifi-  
47 cantly different average and individual structures in comparison with those  
48 in free environments. Effective confining arises in the case of macromolecules  
49 that have specific hydrophobic and hydrophilic regions or when regions have re-  
50 stricted flexibility or torsion. While, in general, one might expect that much is  
51 now known concerning the knotting of macromolecules in such environments,  
52 in fact little is known rigorously and many fundamental questions appear to be  
53 beyond immediate reach, both theoretically or via numerical studies.

## 54 **2 Knotting in DNA and polymers**

55 One of the key themes of this workshop was the focus upon the implications of  
56 experimental results in the context of theoretical models to understand them  
57 and their physiological implications. Setting the theme, **Lynn Zechiedrich's**  
58 opening session described the role of knotting on gene function by leading to a  
59 significant increase in mutation. DNA must be long enough to encode for the  
60 complexity of an organism, yet thin and flexible enough to fit within the cell.  
61 The combination of these properties greatly favors DNA collisions, which can  
62 tangle the DNA. Despite the well-accepted propensity of cellular DNA to collide  
63 and react with itself, it is not clear what the physiological consequences are.  
64 When cells are broken open, the classified knots have all been found to be the  
65 mathematically interesting twist knots. These remarkable knots can have very  
66 high knotting node numbers (complexity), but can be untied in only one strand  
67 passage event. Zechiedrich's group used the *Hin* site-specific recombination  
68 system to tie twist knots in plasmids in *E. coli* cells to assess the effect of knots  
69 on the function of a gene. Knots block DNA replication and transcription. In  
70 addition, knots promote DNA rearrangements at a rate four orders of magnitude  
71 higher than an unknotted plasmid. These results show that knots are potentially

72 toxic, and may help drive genetic evolution. The enzymes that untie knots are  
73 the type II topoisomerases. How they carry out their function to unknot and not  
74 knot DNA is largely unknown. Although domains of type II topoisomerases have  
75 been crystallized and the atomic structures solved, no complete, intact, active  
76 enzyme structure is known and no co-crystals with DNA have been obtained.  
77 Zechiedrich's group used electron cryomicroscopy (CryoEM) to generate the  
78 first three-dimensional structure of any intact, active type II topoisomerase.  
79 The data suggest a simple one-gate mechanism for enzyme function.

80 **Jennifer Mann** described how human topoisomerase II  $\alpha$  resolves DNA  
81 twist knots in a single step. Cellular DNA knotting is driven by DNA com-  
82 paction, topoisomerization, replication, supercoiling-promoted strand collision,  
83 and DNA self-interactions resulting from transposition, site-specific recombina-  
84 tion, and transcription. Type II topoisomerases are ubiquitous, essential en-  
85 zymes that inter-convert DNA topoisomers to resolve knots. These enzymes  
86 pass one DNA helix through another by creating an enzyme-bridged transient  
87 break. How type II topoisomerases accomplish their unknotting feat is a cen-  
88 tral question. Will a type II topoisomerase resolve a DNA twist knot in one  
89 cycle of action? Each crossing reversal performed by a type II topoisomerase  
90 requires energy. Within the cell, DNA knots might be pulled tight by forces  
91 such as those which accompany transcription, replication, and segregation, thus  
92 increasing the likelihood of DNA damage. The results show DNA knots can be  
93 lethal and promote mutations. Therefore, it would be advantageous for type  
94 II topoisomerases to resolve DNA knots in the most efficient manner. Mann's  
95 data show that purified five- and seven-noded twist knots are converted to the  
96 unknot by human topoisomerase II  $\alpha$  with no appearance of either trefoils or  
97 five-noded twist knots which are intermediates if the enzyme acted on one of  
98 the inter-wound nodes.

99 **Dorothy Buck** presented a topological model that predicts which knots  
100 and links are the products of site-specific recombination. Buck described the  
101 topology of how DNA knots and links are formed as a result of a single recombina-  
102 tion event, or multiple rounds of (processive) recombination events, starting  
103 with substrate(s) consisting of an unknot, an unlink, or a  $(2, n)$ -torus knot or  
104 link. The model relies on only three assumptions and Buck provided biological  
105 evidence for each of these assumptions. This talk presented the biological back-  
106 ground, evidence, and applications of the model that was further explored in the  
107 talk of Erica Flapan. The biological determination is accomplished by describ-  
108 ing the topology of how DNA knots and links are formed as a result of a single  
109 recombination event, or multiple rounds of (processive) recombination events,  
110 starting with substrate(s) consisting of an unknot, an unlink, or a  $(2, n)$ -torus  
111 knot or link.

112 **Giovanni Dietler** reported on the properties of knotted DNA in respect  
113 to the critical exponents and the localization of the knot crossings. He showed  
114 that probably two universality classes exist in this case and that localization of  
115 the knot crossings could explain the activity of the topoisomerases. Gel elec-  
116 trophoresis of DNA knots was discussed and simulations as well as experiments  
117 were presented in which the knot complexity and its topology play an essential

118 role. Some hydrodynamics experiments with knots were presented at the end.

### 119 **3 Mathematical, statistical, and computational** 120 **methods**

121 Discussing models employed in modeling DNA molecules, **Alexander**  
122 **Vologodskii** put the attention on the discrete worm-like chain, a carefully  
123 tested model that leads to a reliable analysis of enzymatic topological trans-  
124 formations. First, he described what exactly can be computed by the method,  
125 and how the computational results can be used to test a particular model of the  
126 enzyme action used in the simulation. He showed how two kinds of experimen-  
127 tal data can be compared with the simulation results and discussed the major  
128 assumptions and theoretical bases of the approach. Then the key elements of  
129 the simulation were briefly considered. This general description of the approach  
130 was illustrated by specific examples.

131 **Hue Sun Chan** described the statistical mechanics of how recognition of lo-  
132 cal DNA juxtaposition geometry may underlie the unknotting and decatenating  
133 actions of type II topoisomerases. Topoisomerases may unknot and decate-  
134 nate by recognizing specific DNA juxtapositions. The statistical mechanical  
135 viability of this hypothesis was investigated by considering lattice models of  
136 single-loop conformations and two-loop configurations of ring polymers. Using  
137 exact enumerations and Monte Carlo sampling, the statistical relationship be-  
138 tween the local geometry of a juxtaposition of two chain segments on one hand,  
139 and whether a single loop was knotted or whether two loops were linked glob-  
140 ally on the other was determined; and it was ascertained how the knot/unknot  
141 topology and global linking were altered by a topoisomerase-like segment pas-  
142 sage at the juxtaposition. Presented results showed that segment passages at  
143 a “free” juxtaposition tend to increase knot probability but segment passages  
144 at a “hooked” juxtaposition cause more transitions from knot to unknot than  
145 vice versa, resulting in a steady-state knot probability far lower than that at  
146 topological equilibrium. Similarly, the selective segment passage at hooked jux-  
147 tapositions can lower catenane populations significantly. A general exhaustive  
148 analysis of 6,000 different juxtaposition geometries showed that the ability of a  
149 segment passage to unknot and decatenate correlates strongly with a juxtaposi-  
150 tion’s “hookedness.” Most remarkably, and consistent with earlier experiments  
151 on type II topoisomerases from different organisms, the unknotting potential of a  
152 juxtaposition geometry in the presented model correlates almost perfectly with  
153 its corresponding decatenation potential. These quantitative findings suggest  
154 that it is possible for type II topoisomerases to disentangle by acting selectively  
155 on juxtapositions with hook-like geometries.

156 **Andrzej Stasiak** presented another perspective on a model of selective  
157 simplification of DNA topology by DNA topoisomerases. The presented model  
158 tested the hypothesis that type II DNA topoisomerases maintain the steady  
159 state level of DNA knotting below the thermodynamic equilibrium by acting

160 as topological filters that recognize preferentially certain geometrical arrange-  
161 ments of juxtaposed segments, “hooked relationships”. It was shown that such  
162 specificity can result in two interrelated topological consequences: maintaining  
163 the steady-state knot probability level below the topological equilibrium and  
164 selecting a specific way of relaxation of more complex knots. It was observed,  
165 in addition, that local structures in random configurations of a given knot sta-  
166 tistically behave as analogous local structures in ideal geometric configurations  
167 of the corresponding knot types.

168 **Mariel Vazquez** contributed to the theme of modeling DNA topology sim-  
169 plification. Random cyclization of linear DNA can result in knotted DNA circles.  
170 Experiments on DNA confined inside P4 viral capsids have found knotting prob-  
171 abilities as high as 0.95. A full description of the complicated knots remains  
172 unavailable. Type II topoisomerases unknot DNA very efficiently by perform-  
173 ing strand-passage on DNA strands. Motivated by these biological observations,  
174 Vazquez and colleagues studied random state transitions in knot space for all  
175 prime knots with 8 or fewer crossings and fixed length. The main goal was  
176 to quantify unknotting under different geometrical constraints. The long term  
177 goal is to understand the mechanism of action of type II topoisomerases, and  
178 to characterize the knots extracted from the P4 capsids. They used the Monte  
179 Carlo based BFACF algorithm to generate ensembles of self-avoiding polygons  
180 (SAP) in  $Z^3$  with identical knot type and fixed length. The BFACF algorithm  
181 produces a reducible Markov chain whose ergodicity classes are the knot types.  
182 They performed random strand-passage on these knots, computed state transi-  
183 tions between knot types, and steady-state distributions after repeated strand-  
184 passages. Introducing different topological biases resulted in various probability  
185 distributions. The large amount of knots used in their model made it possible to  
186 gather additional information regarding knots and their projections. They com-  
187 puted minimal lattice knots, and in some cases improve existing lower bounds.  
188 They also provided other physical measures such as the writhe and average  
189 crossing number. Finally, using an algorithm that removes Reidemeister I and  
190 II moves simultaneously, they computed the average number of crossings before  
191 and after Reidemeister removal.

192 **Christine Soteris** discussed the asymptotics of knotting after a local  
193 strand passage. On the macroscopic scale, circular DNA can be viewed simply  
194 as a ring polymer. Experimental evidence indicates that topoisomerases act  
195 locally in DNA allowing two strands of the DNA which are close together to  
196 pass through one another (i.e. enabling a “local” strand passage) in order to  
197 disentangle the DNA. This has inspired investigation of the following question  
198 about self-avoiding polygon (SAP) models: Given a SAP with a fixed knot type,  
199 how does the distribution of knots after a local strand passage depend on the  
200 initial knot type of the SAP, the length of the SAP, and on the specific details  
201 of the strand passage such as where the strand passage occurs and the number  
202 of edges altered in the strand passage? In 2000, graduate student M. Szafron  
203 introduced a model of unknotted ring polymers in dilute solution for which it is  
204 assumed that two segments of the polymer have already been brought close to-  
205 gether for the purposes of performing a local strand passage. The conformations

206 of the ring polymer are represented by  $n$ -edge unknotted polygons containing  
207 a specific pattern (designed to facilitate a strand passage in which exactly two  
208 segments of the polygon pass through each other) on the simple cubic lattice.  
209 Based on the assumption that each such SAP conformation is equally likely,  
210 Soteris and Szafron investigated, both theoretically and numerically, the distri-  
211 bution of knots after a strand passage has been performed at the location of the  
212 special pattern. The talk reviewed the theoretical and numerical (via Markov  
213 Chain Monte Carlo) results for this model with emphasis on the asymptotic  
214 properties as  $n$  increases. In addition, results for the extension of the model to  
215 other knot types such as the figure-eight knot were presented.

216 **Enzo Orlandini** discussed the topological effects of knotting on the dynam-  
217 ics of polymers. Knots are frequent in long polymer rings at equilibrium and  
218 it is now well established that their presence can affect the static properties of  
219 the polymer. On the other hand, topological constraints (knots) influence also  
220 the dynamical properties of a polymer. This has been shown in recent exper-  
221 iments where the motion of a single knotted DNA has been followed within a  
222 viscous solution and in the presence of a stretching force. These experiments  
223 raise interesting challenges to the theoretical understanding of the problem, an  
224 issue that is still in its infancy. As a first step towards the understanding of  
225 the mechanism underlying the mobility of a knot, the relaxation and diffusion  
226 dynamics of flexible knotted rings in equilibrium under good solvent conditions  
227 was investigated by Monte Carlo simulations. By focusing on prime knots and  
228 using a knot detection algorithm it was possible to monitor the diffusion in  
229 space of the knotted part of the ring, and observe in time the fluctuations of its  
230 length along the backbone. This identified a novel, slow topological time-scale,  
231 and to show that it is related to a self-reptation of the knotted region. For open  
232 chains, knotted configurations do not represent an equilibrium state any more.  
233 However, under suitable conditions (for example very tight knots or quite rigid  
234 chains), knotted metastable states persist for a very long time and a statistical  
235 description of their dynamical properties is then possible. By performing off  
236 lattice molecular dynamic simulations of a semiflexible polymer, an estimate  
237 was obtained of the average living time and the stability of these states as a  
238 function of the initial conditions (size of the initial knot) and of the rigidity of  
239 the chain.

240 **Carla Tesi** discussed the probability of knotting of polygons under a stretch-  
241 ing force. Knots are practically unavoidable in long polymer rings and influence  
242 their properties. This has been witnessed by an increasing number of exper-  
243 iments that can nowadays probe the detailed properties of knotted molecules.  
244 In particular micro-manipulation techniques enable direct measurements of me-  
245 chanical properties of a single molecule, and it is also possible to probe the  
246 behavior of artificially knotted DNA. It is becoming important to study theo-  
247 retically how, for example, the presence of topological constraints (knots) can  
248 affect the mechanical or elastic responses of knotted molecules under external  
249 forces. As a first step in this direction Tesi and colleagues considered first the  
250 problem of looking at how the entanglement complexity in ring polymers can  
251 be affected by the presence of a tensile or contractile force. A possible experi-

252 mental realization of this problem could be bacterial (or mitochondrial) DNA in  
253 solution with topoisomerases that are subjected to an external force (AFM or  
254 optical tweezers) or to flow files (shear flow for example). In this work stretched  
255 ring polymers are modeled by polygons in the cubic lattice weighted by a fugac-  
256 ity coupled to its span along a given direction. By performing extensive Monte  
257 Carlo simulations on this system they have been able to estimate how the knot-  
258 ting probability and the knot spectra depends on the force strength, both in  
259 the extensile and in the contractile regime. These findings were compared with  
260 recent rigorous results on similar models of stretched polygons.

261 **Isabel Darcy** described the modeling of protein-DNA complexes in three  
262 dimensions using TopoICE (Topological Interactive Construction Engine). Protein-  
263 DNA complexes have been modeled using tangles. A tangle consists of arcs  
264 properly embedded in a 3-dimensional ball. The protein is modeled by the 3D  
265 ball while the segments of DNA bound by the protein can be thought of as  
266 arcs embedded within the protein ball. This is a very simple model of protein-  
267 DNA binding, but from this simple model, much information can be gained.  
268 The main idea is that when modeling protein-DNA reactions, one would like  
269 to know how to draw the DNA. For example, are there any crossings trapped  
270 by the protein complex? How do the DNA strands exit the complex? Is there  
271 significant bending? Tangle analysis cannot determine the exact geometry of  
272 the protein-bound DNA, but it can determine the overall entanglement of this  
273 DNA, after which other techniques may be used to more precisely determine  
274 the geometry. KnotPlot, developed by Rob Scharein, is an interactive 3D pro-  
275 gram for visualizing and manipulating knots. TopoICE-X is a subroutine within  
276 KnotPlot for solving tangle equations modeling topoisomerase reactions.

277 **Eric Flapan** described the topological faces of the model for DNA knotting  
278 and linking developed jointly with Dorothy Buck. Flapan presented a topologi-  
279 cal model that predicts which knots and links can be the products of site-specific  
280 recombination. This is done by describing the topology of how DNA knots and  
281 links are formed as a result of a single recombination event, or multiple rounds  
282 of (processive) recombination events, starting with substrate(s) consisting of an  
283 unknot, an unlink, or a  $(2, n)$ -torus knot or link. The model relies on only three  
284 assumptions and we give biological evidence for each of these assumptions.

285 **Alexander Grosberg** described metastable tight knots as a worm-like poly-  
286 mer. Based on an estimate of the knot entropy of a worm-like chain. Grosberg  
287 and colleagues predict that the interplay of bending energy and confinement  
288 entropy will result in a compact metastable configuration of the knot that will  
289 diffuse, without spreading, along the contour of the semi-flexible polymer un-  
290 til it reaches one of the chain ends. The estimate of the size of the knot as a  
291 function of its topological invariant (ideal aspect ratio) agrees with recent ex-  
292 perimental results of knotted dsDNA. Further experimental tests of these ideas  
293 were proposed.

294 **Bertrand Duplantier** discussed random linking of curves and manifolds.  
295 Duplantier proposed a formalism for evaluating random linking integrals of  
296 closed curves in  $\mathbb{R}^3$  or, more generally, manifolds in  $\mathbb{R}^n$ , all in relative motions.  
297 It is based on the existence of universal geometric characteristic functions for

298 each closed curve or manifold separately. It allows further averaging over the  
299 possible random shapes of those curves and manifolds.

300 **Tetsuo Deguchi** discussed the dynamics and statistical mechanics of knot-  
301 ted ring polymers in solution using a simulations approach toward an experimen-  
302 tal confirmation of topological effects. Deguchi described how topological effects  
303 may give nontrivial results on the macroscopic behavior of ring polymers in so-  
304 lution and how one can confirm them experimentally. Numerical evaluations  
305 of some characteristic physical quantities of the solution that can be measured  
306 in polymer experiments were presented. This study was strongly motivated  
307 by recent experimental developments for synthesizing ring polymers with large  
308 molecular weights. Numerical results on dynamical and statistical properties of  
309 a dilute solution of ring polymers where topological constraints play a central  
310 role were presented. Dynamical quantities such as the diffusion constants of  
311 ring polymers in solution and the viscosity of the ring-polymer solution were  
312 discussed. These show their difference from those of the corresponding linear  
313 polymers with the same molecular weights. Secondly, the osmotic pressure of  
314 the ring-polymer solution reflects the topological interaction among ring poly-  
315 mers. It was numerically evaluated in terms of the random linking probability.  
316 Thirdly, the mean square radius of gyration of ring polymers under a topolog-  
317 ical constraint, which is one of the most fundamental quantities in the physics  
318 of knotted ring polymers, can be measured in the scattering experiment. The  
319 single-chain static structure factor, i.e. the scattering function, can be obtained  
320 experimentally for ring polymers with fixed topology, from which one derives  
321 the mean square radius of gyration. It is therefore important to evaluate nu-  
322 merically the scattering function of a knotted ring polymer in solution. Some  
323 theoretical and simulational results on the scattering functions were discussed.

324 **Kenneth Millett** discussed the problem of estimating the number of dis-  
325 tinct topological knot types and their proportion in the space of (equilateral)  
326 polygonal knots with a fixed number of edges. For very small numbers of edges,  
327 one knows the number of knot types and can estimate their proportion but, for  
328 larger numbers of edges, only rough estimates are available. Estimates derive  
329 from Monte Carlo explorations of the (equilateral) polygonal knot space and an  
330 analysis using the HOMFLY polynomial as a surrogate for the topological knot  
331 type. As a consequence, one is interested in knowing how large a sample of knots  
332 is needed to give a good estimate of the number of topological knot types as  
333 detected by distinct HOMFLY polynomials. Some theoretical and experimental  
334 efforts concerning this question were discussed.

335 **Rob Kusner** discussed the geometric problems for embedded bands in  
336 space. Just as one can minimize the ropelength for knotted or linked space  
337 curves, one can also minimize the analogous “bandlength” for smoothly framed  
338 curves, either within a framed isotopy class, or with a pointwise constraint on  
339 the framing (which we view as a normal vector field along the corresponding  
340 bands). As a limiting case where the framing for the bands is constant, one gets  
341 knotted or linked “raceways” in the plane, a flattened analogue of knotted or  
342 linked “ropes” in space. Kusner showed that the bandlength of raceways grows  
343 at least as fast as the square root of crossing number (recall that for ropes one

344 had instead the three-fourths power) and that this power is sharp. Kusner also  
345 commented on the shapes of length minimizing raceways, and speculated on  
346 bands or raceways as models for folded or packed proteins.

347 **Atilio Stella** discussed how the probability of realization of configurations  
348 with specific knots in closed random chains play a major role in topological poly-  
349 mer statistics and in its applications to macromolecular and biological physics.  
350 A problem of considerable current interest is that of comparing the knot spectra  
351 obtained for random models with those analyzed by electrophoresis for the DNA  
352 extracted from viral capsids. This comparison should help in identifying specific  
353 mechanisms of knot formation in the biological context. In the case of collapsed  
354 polymer rings, interest in the knot spectrum is also enhanced by the recent  
355 discovery that knots are fully delocalized along the backbone. Understanding  
356 if, and up to what extent, topological invariants can affect the globular state  
357 in such conditions is an intriguing fundamental issue. An analysis of extensive  
358 Monte Carlo simulations of interacting self-avoiding polygons on cubic lattice  
359 was presented. The results showed that the frequencies of different knots real-  
360 ized in a random, collapsed polymer ring decrease as a power (about -0.6) of  
361 the ranking order. This Zipf type of law also suggests that the total number  
362 of different knots realized grows exponentially with the chain length. Relative  
363 frequencies of specific knots converge to definite ratios for long chains, because  
364 of the free energy per monomer and its leading finite size corrections do not  
365 depend on the ring topology, while a subleading correction only depends on the  
366 minimal crossing number of the knots. This topological invariant appears to  
367 play a fundamental role in the statistics of collapsed polymers.

368 **Jon Simon** discussed the problem of measuring tangling in a large filament  
369 system. Imagine a protein or other polymer filament (or several) entangled in  
370 some complicated way, perhaps with tens or hundreds of crossings. Now imag-  
371 ine a second example with similarly large entanglement. Can one say something  
372 useful to distinguish the tangling in the two examples? For relatively small sys-  
373 tems, topological knotting and linking is a powerful tool, witness the success  
374 of “topological enzymology”. But for large systems, calculating exact knotting  
375 and linking may be computationally impractical; there are uncertainties in how  
376 to deal with open filaments; and knowing that one is knot 10.156 and the other  
377 10.157 might not tell us much about the physical properties of the given sys-  
378 tem. Simon proposed that describing and quantifying tangling in large filament  
379 systems should be one of the important next-stage problems for the field of  
380 physical knots. To describe shapes of proteins (in static conformations), several  
381 researchers have developed numerical descriptors based on variations of Gausss  
382 linking-number integrals; these are related to average crossing number. Simon  
383 has begun studying another modification of average crossing number called the  
384 average bridging number. This is a simple idea, but when taken together with  
385 average crossing number, it seems to distinguish nicely between different kinds  
386 of packings for long filaments. And there appears to be reasonable stability of  
387 the relationship under random perturbations, so this approach may be useful  
388 for statistical ensembles as well as for individual conformations.

389 **Jason Cantarella** gave a talk intended as an (mostly expository) invitation  
390 to the community interested in modeling large molecules to consider an alter-  
391 nate mathematical framework for their work: modeling large macromolecules  
392 as divergence-free vector fields instead of as curves, polygons, chains, or tubes.  
393 From this point of view, the actual topological knot type of a very large and  
394 complicated curve will be seen as less important than its average entanglement  
395 complexity. The talk introduced this framework, reviewed some older results  
396 about the helicity of vector fields (which measures a kind of average linking  
397 number of integral curves), outlined some speculative applications to macro-  
398 molecules, and introduced some work in progress reformulating the helicity of  
399 vector fields from a more modern perspective. Cantarella’s reformulation of  
400 helicity opens the possibility of constructing a family of “generalized helicity”  
401 integrals analogous to finite-type invariants for knots.

402 **Claus Ernst** gave a summary of what is currently known about the topo-  
403 logical aspects of lattice knots such as their length and curvature. The length  
404 as braids is also considered.

405 **Eric Rawdon** presented computer simulations to examine the equilibrium  
406 length of random equilateral polygons with respect to different spatial quanti-  
407 ties, in particular with respect to the total curvature and total torsion of the  
408 polygons. Rawdon and colleagues use Markov Chain Monte Carlo methods to  
409 determine likely scaling profiles and error bars for the equilibrium length calcu-  
410 lations

411 **John Maddocks** discussed the optimal packing of tubes in  $\mathbb{R}^3$  and  $\mathbb{S}^3$ ,  
412 contact sets in  $\mathbb{R}^3$ , and connections with sedimentation dynamics.

413 **Henryk Gerlach** described the optimal packing of curves on  $S^2$ , both fam-  
414 ilies of circles and open curves.

415 **Stuart Whittington** reviewed some results about lattice models of ring  
416 polymers, focusing on rigorous asymptotic results about the knot probability as  
417 a function of length, the topological and geometrical entanglement complexity  
418 and the relative frequency of occurrence of different link types. He discussed a  
419 number of open questions. For instance, we know that the knot probability goes  
420 to unity exponentially rapidly as the size of the lattice polygon goes to infinity  
421 but we know almost nothing (rigorously) about the constant appearing in the  
422 exponential term. Similarly, although we know that all non-trivial link types  
423 where both polygons are knotted grow at the same exponential rate, we know  
424 nothing about the sub-exponential terms.

## 425 4 Macromolecules in confined geometries

426 **Javier Arsuaga** discussed the topological considerations of the interphase nu-  
427 cleus. During the early phase of the cell cycle (G0/G1) chromosomes are con-  
428 fined to spherical regions within the nucleus called chromosome territories. The  
429 position of these territories is important in a number of biological processes  
430 (e.g. transcription, replication and DNA repair) and has important implica-  
431 tions in human genetic diseases, in cancer and in the formation of chromosome

432 aberrations after exposure to DNA damaging agents. Recently, a model has  
 433 been proposed for the interface region between territories in which chromo-  
 434 somes overlap and intermingle. This new model naturally raises the question  
 435 of whether chromosomes are linked or not. Motivated by this problem Arsuaga  
 436 and colleagues investigated the linking of curves in confined volumes. Arsuaga  
 437 presented recent results using the uniform random polygon model. First, ana-  
 438 lytically, they showed that the linking probability between a fixed closed curve  
 439 and a random polygon of length  $n$  increases as  $1 - O((\frac{1}{n})^{\frac{1}{2}})$ . Next, numerically  
 440 that the linking probability between two polygons of lengths  $n$  and  $m$  increase  
 441 as  $1 - O((\frac{1}{nm})^{\frac{1}{2}})$ . They extended these results to the case when two polygons  
 442 have a predetermined overlapping volume (as is the case in experimental obser-  
 443 vations). Arsuaga concluded with a discussion of potential extensions to other  
 444 polymer models and biological implications.

445 **Buks Janse van Rensburg** discussed the properties of lattice polygons of  
 446 fixed knot types in a slab of width,  $w$ , by using scaling arguments and presented  
 447 numerical results from Monte Carlo simulations using the BFACF algorithm. If  
 448  $p_n(K)$  is the number of polygons of length  $n$  and of knot type  $K$  in the cubic  
 449 lattice, then it is known that  $\lim_{n \rightarrow \infty} \frac{[\log(p_n(\emptyset))]}{n} = \log(\mu_\emptyset)$  exists, where  $K = \emptyset$   
 450 is the unknot, and  $\mu_\emptyset$  is the growth constant of unknotted polygons in the  
 451 cubic lattice. Suppose that  $p_n(K, w)$  is the number of knotted polygons of  
 452 length  $n$  and of knot type  $K$  in a slab of width  $w$  in the cubic lattice. The  
 453 generating function of this model is given by  $g_K(w; t) = \sum p_n(K, w) t^n$ , where  $t$  is  
 454 a generating variable conjugate to the length of the polygons. The mean length  
 455  $\langle n \rangle_{K, w}$  of polygons of knot type  $K$  in a slab of width  $w$  may be estimated  
 456 from  $g_K(w; t)$  using the BFACF algorithm. The dependence of  $\langle n \rangle_{K, w}$  on  $w$   
 457 was estimated for  $t = \mu_\emptyset^{-1}$ , and the results were compared to predictions of  
 458 scaling arguments. In addition, numerical results for the metric properties of  
 459 knotted polygons in this ensemble were presented.

460 **De Witt Summers** discussed why DNA knots reveal chiral packing of DNA  
 461 in phage capsids. Bacteriophages are viruses that infect bacteria. They pack  
 462 their double-stranded DNA genomes to near-crystalline density in viral capsids  
 463 and achieve one of the highest levels of DNA condensation found in nature.  
 464 Despite numerous studies, some essential properties of the packaging geometry  
 465 of the DNA inside the phage capsid are still unknown. Although viral DNA is  
 466 linear doublestranded with sticky ends, the linear viral DNA quickly becomes  
 467 cyclic when removed from the capsid, and for some viral DNA the observed  
 468 knot probability is an astounding 95%. Summers discussed comparison of the  
 469 observed viral knot spectrum with the simulated knot spectrum, concluding  
 470 that the packing geometry of the DNA inside the capsid is non-random and  
 471 writhe-directed.

472 **Cristian Micheletti** discussed the knotting of ring polymers in confined  
 473 spaces. Stochastic simulations were used to characterize the knotting distri-  
 474 butions of random ring polymers confined in spheres of various radii. The  
 475 approach was based on the use of multiple Markov chains and reweighting tech-  
 476 niques, combined with effective strategies for simplifying the geometrical com-

477 plexity of ring conformations without altering their knot type. By these means,  
478 Micheletti and colleagues extended previous studies and characterized in detail  
479 how the probability to form a given prime or composite knot behaves in terms  
480 of the number of ring segments  $n$  and confining radius  $R$ . For  $50 \leq n \leq 450$  they  
481 showed that the probability of forming a composite knot rises significantly with  
482 the confinement, while the occurrence probability of prime knots are, in general,  
483 nonmonotonic functions of  $\frac{1}{R}$ . The dependence of other geometrical indicators,  
484 such as writhe and chirality, in terms of  $R$  and  $n$  was also characterized. It was  
485 found that the writhe distribution broadens as the confining sphere narrows

486 **Yuanan Diao** discussed the sampling of large random knots in a confined  
487 space. Diao proposed 2-dimensional uniform random polygons as an alternative  
488 method of sampling large random knot diagrams. In fact, the 2-dimensional  
489 uniform random polygons allow one to sample knot diagrams with large crossing  
490 numbers that are diagrammatically prime since one can rigorously prove that the  
491 probability that a randomly selected 2D uniform random polygon of  $n$  vertices  
492 is almost diagrammatically prime (in the sense that the diagram becomes a  
493 reduced prime diagram after a few third Reidemeister moves) goes to one as  $n$   
494 goes to infinity, and that the average number of crossings in such a diagram is  
495 on the order of  $O(n^2)$ . This strongly suggests that the 2-dimensional uniform  
496 random polygons are good candidates if one is interested in sampling large  
497 (prime) knots. Numerical studies on the 3D uniform random polygons show  
498 that these polygons for complicated knots even when they have relatively small  
499 number of vertices.

500 **Andrew Rechnitzer** talked about the mean unknotting times of random  
501 knots and knot embeddings by crossing reversals, in a problem motivated by  
502 DNA entanglement. Using self-avoiding polygons (SAPs) and self-avoiding poly-  
503 gon trails (SAPTs) Rechnitzer and colleagues proved that the mean unknotting  
504 time grows exponentially in the length of the SAPT and at least exponentially  
505 with the length of the SAP. The proof uses Kesten's pattern theorem, together  
506 with results for mean first-passage times in the two-parameter Ehrenfest urn  
507 model. They used the pivot algorithm to generate random SAPTs of up to  
508 3000 steps, calculated the corresponding unknotting times, and found that the  
509 mean unknotting time grows very slowly even at moderate lengths. These meth-  
510 ods are quite general—for example the lower bound on the mean unknotting  
511 time applies also to Gaussian random polygons. This work was accomplished  
512 in collaboration with Aleks Owcarek and Yao-ban Chan at the University of  
513 Melbourne, and Gord Slade at the University of British Columbia.