# A Hörmander type multiplier theorem for arbitrary discrete groups

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Noncommutative  $L_p$  spaces, Operator spaces and Applications Banff International Research Station. June 28, 2010

# The problem

Consider a Fourier multiplier on  $(\mathbb{T}^n,\mu)$ 

$$T_m\left(\sum_{k\in\mathbb{Z}^n}\widehat{f}(k)\,e^{2\pi i\,\langle k,\cdot\rangle}\right) = \sum_{k\in\mathbb{Z}^n}m_k\,\widehat{f}(k)\,e^{2\pi i\,\langle k,\cdot\rangle}.$$

A lifting multiplier for  $\boldsymbol{m}$  is an smooth function

 $\tilde{m}: \mathbb{R}^n \to \mathbb{C}$  satisfying  $\tilde{m}_{|_{\mathbb{Z}^n}} = m$ .

It is well-known that  $L_p$ -boundedness is preserved, so that we have

 $\left|\partial_{\xi}^{\beta}\tilde{m}(\xi)\right| \leq c_n \left|\xi\right|^{-\left|\beta\right|} \text{ for all } \left|\beta\right| \leq \left[\frac{n}{2}\right] + 1 \quad \Rightarrow \quad T_m : L_p(\mathbb{T}^n,\mu) \to L_p(\mathbb{T}^n,\mu).$ 

In the case of arbitrary discrete groups

- There is no canonical differential structure to work with.
- No sufficient differentiability conditions are known for  $L_p$ -boundedness.

Our main goals in this talk is to present

- A Hörmander multiplier theorem for arbitrary discrete groups.
- A noncommutative Calderón-Zygmund theory for von Neumann algebras.

## **Compact duals**

Let  $\boldsymbol{G}$  be a discrete group and

 $f\sim \sum_{g\in \mathbf{G}}\widehat{f}(g)\lambda(g)\in L_p(\widehat{\mathbb{G}},\tau) \quad \text{such that} \quad \widehat{f}(g)\,=\,\tau(f\lambda(g^{-1}))$ 

a Fourier series on its compact dual, where:

•  $\lambda : G \to \mathcal{B}(\ell_2(G))$  is the left regular representation  $\lambda(g)\delta_h = \delta_{gh}$  for  $\delta_g(h) = \delta_{g=h}$ , and  $\widehat{\mathbb{G}} = \mathcal{L}(G)$  is the weak operator closure of  $\operatorname{span}\lambda(G)$  in  $\mathcal{B}(\ell_2(G))$ . • We equip the group von Neumann algebra  $\mathcal{L}(G)$  with its natural trace  $\tau(f) = \langle f\delta_e, \delta_e \rangle_{\ell_2(G)} = \widehat{f}(e)$  where  $e = \operatorname{Identity}$  of G.  $L_p(\widehat{\mathbb{G}}, \tau)$  is the noncommutative  $L_p$  space associated to  $(\mathcal{L}(G), \tau)$ , with norm  $\|f\|_p = (\tau |f|^p)^{\frac{1}{p}}$  and  $\|f\|_{\infty} = \|f\|_{\mathcal{B}(\ell_2(G))}$ .

**Example.** If  $G = \mathbb{Z}$ , we have an isometry

$$\ell_2(\mathbf{G}) \ni \delta_m \mapsto \exp(2\pi i m \cdot) \in L_2(\mathbb{T}, \mu) \implies \lambda(m) \sim \exp(2\pi i m \cdot).$$

We have  $\tau = \int_{\mathbb{T}} \cdot d\mu$  for the normalized Haar measure  $\mu$  and  $L_p(\widehat{\mathbb{G}}, \tau) = L_p(\mathbb{T}, \mu)$ .

## Sketch of the talk

Given a discrete group  $\mathrm{G},$  the key points are:

• Cocycles. If  $\psi: \mathbf{G} \to \mathbb{R}_+$  is such that

 $e^{-t\psi}$  is positive definite for all t > 0,

we may find an inclusion  $b_{\psi} : G \to \mathcal{H}_{\psi}$  into a Hilbert space  $\mathcal{H}_{\psi}$ , which is isometric in the sense that

$$\operatorname{dist}(g,h) \equiv \sqrt{\psi(g^{-1}h)} = \left\| b_{\psi}(g) - b_{\psi}(h) \right\|_{\mathcal{H}_{\psi}}$$

defines a **pseudo-metric** on G, and  $\mathcal{H}_{\psi}$  plays the role of an **'ambient space'** for G.

• Hörmander-Mihlin condition. Given a Fourier multiplier

$$T_m: \sum_g \widehat{f}(g)\lambda(g) \mapsto \sum_g m_g \widehat{f}(g)\lambda(g),$$

a lifting multiplier is given by  $\tilde{m} : \mathcal{H}_{\psi} \to \mathbb{C}$  with  $m = \tilde{m} \circ b_{\psi}$ . If  $\dim \mathcal{H}_{\psi} = n$ 

$$\left|\partial_{\xi}^{\beta}\widetilde{m}(\xi)\right| \leq c_n \left|\xi\right|^{-\beta}$$
 for all  $\left|\beta\right| \leq \left[\frac{n}{2}\right] + 1 \xrightarrow{?} T_m : L_p(\widehat{\mathbb{G}}, \tau) \to L_p(\widehat{\mathbb{G}}, \tau)$ 

• Calderón-Zygmund theory. This naturally leads to a form of

$$\operatorname{ess\,sup}_{x\in\mathbb{R}^n} \int_{|s|>2|x|} \left| k_{\tilde{m}}(s-x) - k_{\tilde{m}}(s) \right| ds < \infty \quad \Rightarrow \quad T_{\tilde{m}} : L_{\infty} \to \operatorname{BMO}$$

the Hörmander condition for the kernel, valid for arbitrary von Neumann algebras.

# Length functions

We will be working with functions  $\psi : G \to \mathbb{R}_+$  such that:

i)  $\psi(e) = 0$ ,

- ii)  $\psi(g) = \psi(g^{-1})$  for all  $g \in \mathbf{G}$ ,
- iii)  $\psi$  is a conditionally negative function

$$\sum_{g\in\Lambda\subset\mathbf{G}}\gamma_g=0\quad\text{and}\quad |\Lambda|<\infty\quad\Rightarrow\quad \sum_{g,h\in\Lambda}\overline{\gamma}_g\gamma_h\psi(g^{-1}h)\leq 0.$$

We will call such a  $\psi$  a length function. By Schoenberg theorem

 $\psi$  length function  $\Leftrightarrow e^{-t\psi}$  positive definite for all t > 0.

#### **Examples.** Two standard cases:

• If  $G = \mathbb{Z}^n$ , we may take

$$\psi_1 = | |^2$$
 and  $\psi_2 = | |$ .

The heat and Poisson kernels are positive with Fourier transforms  $e^{-t\psi_j}$ , j = 1, 2.

• If  $G = \mathbb{F}_n$  is the free group with n generators, we may use  $\psi(g) = |g| = \text{ standard length function},$ 

because the associated Poisson semigroup is formed of completely positive maps.

### **Construction of the cocycle**

Given a length function  $\psi$ 

$$K_{\psi}(g,h) = \frac{\psi(g) + \psi(h) - \psi(g^{-1}h)}{2} \rightsquigarrow \left\langle \sum_{g \in \mathcal{G}} \gamma_g \delta_g, \sum_{h \in \mathcal{G}} \gamma'_h \delta_h \right\rangle_{\psi} = \sum_{g,h \in \mathcal{G}} \gamma_g K_{\psi}(g,h) \gamma'_h.$$

This is an  $\mathbb{R}$ -product on the group algebra  $\mathbb{R}[G]$  of finitely supported real functions on G. If we set  $N_{\psi}$  to be the null space of  $\langle \cdot, \cdot \rangle_{\psi}$ , we may define the Hilbert space  $\mathcal{H}_{\psi}$  as the completion of  $(\mathbb{R}[G]/N_{\psi}, \langle \cdot, \cdot \rangle_{\psi})$  and define the natural inclusion

 $b_{\psi}: g \in \mathcal{G} \mapsto \delta_g + N_{\psi} \in \mathcal{H}_{\psi}$ 

which satisfies the isometric identity  $\|b_{\psi}(g) - b_{\psi}(h)\|_{\mathcal{H}_{\psi}} = \sqrt{\psi(g^{-1}h)} = \operatorname{dist}(g,h).$ There exists a natural action  $\alpha_{\psi} : \mathbf{G} \to \operatorname{Aut}(\mathcal{H}_{\psi})$ 

$$\alpha_{\psi,g}(b_{\psi}(h)) = b_{\psi}(gh) - b_{\psi}(g)$$

which is isometric in the sense that we have  $\langle \alpha_{\psi,g}(\xi_1), \alpha_{\psi,g}(\xi_2) \rangle_{\psi} = \langle \xi_1, \xi_2 \rangle_{\psi}$ . This allows us to construct a **semidirect product embedding**  $g \mapsto b_{\psi}(g) \rtimes g$  which extends to the group von Neumann algebras as follows

 $\pi_{\psi} : \lambda(g) \in \mathcal{L}(\mathbf{G}) \mapsto \exp b_{\psi}(g) \rtimes \lambda(g) \in \mathcal{L}(\mathcal{H}_{\psi}) \rtimes_{\alpha_{\psi}} \mathbf{G}.$ 

The key is to show  $T_{\tilde{m}}: L_{\infty}(\mathcal{H}_{\psi}) \to BMO \Rightarrow T_{\tilde{m}} \rtimes id_{G}$  is still bounded (NCCZ theory).

## Hörmander-Mihlin multipliers for discrete groups

**Theorem** [JMP]. Let G be a discrete group and

$$T_m: \sum_g \widehat{f}(g)\lambda(g) \mapsto \sum_g m_g \widehat{f}(g)\lambda(g)$$

a Fourier multiplier on its compact dual. Assume that G is equipped with a length function  $\psi$ , with associated cocycle  $b_{\psi} : G \to \mathcal{H}_{\psi}$  such that  $\dim \mathcal{H}_{\psi} = n$ . Let  $\tilde{m} : \mathbb{R}^n \to \mathbb{C}$  be a lifting multiplier for m, so that  $m = \tilde{m} \circ b_{\psi}$ . Then

 $T_m: L_p(\widehat{\mathbb{G}}, \tau) \to L_p(\widehat{\mathbb{G}}, \tau)$  is cb-bounded for 1

provided the condition below holds for some  $\varepsilon > 0$ 

 $\left|\partial_{\xi}^{\beta}\tilde{m}(\xi)\right| \leq c_n \left|\xi\right|^{-\left|\beta\right|-\varepsilon}$  for all multi-indices  $\beta$  s.t.  $\left|\beta\right| \leq n+2$ .

The classical hypotheses with

$$|\beta| \leq \left[\frac{n}{2}\right] + 1 \text{ and } \varepsilon = 0$$

suffice in the following particular cases:

- If  $b_{\psi}(G)$  is a lattice in  $\mathcal{H}_{\psi}$ .
- Radial Fourier multipliers  $m_g = h(\psi(g))$ .

We may also prove  $L_{\infty} \to BMO$  type inequalities and some free-dimensional estimates.

## Some comments

- If  $G = \mathbb{Z}^n$  and  $\psi_1(k) = |k|^2$ , we easily get  $\mathcal{H}_{\psi_1} \simeq \mathbb{R}^n$  and  $b_{\psi_1}(k) = k \Rightarrow$  Classical Hörmander multiplier theorem. However,  $\psi_2(k) = |k|$  gives  $\dim \mathcal{H}_{\psi_2} = \infty$ ! Highly non canonical choice of cocycle.
- The are two problems to solve

 $\circ$  Interpolation problem. Given  $\psi : \mathbf{G} \to \mathbb{R}_+$ , estimate

 $\inf \Big\{ \sup_{\xi \in \mathbb{R}^n} \sup_{|\beta| \le d_n} |\xi|^{-|\beta|} \Big| \partial_\beta \tilde{m}(\xi) \Big| \quad \text{s.t.} \quad \tilde{m} \circ b_\psi(g) = m_g \Big\}.$ 

Related to Fefferman's recent work on 'smooth interpolation of data'. • Dimensional problem. Given G, find  $\inf_{\psi} \dim \mathcal{H}_{\psi}$  for  $b_{\psi} : G \to \mathcal{H}_{\psi}$  injective.

**Remark.** We have  $H-\dim(\mathbb{Z}^n) = 1!!$  Both problems are 'incompatible'.

• The negative generator of the semigroup

 $\lambda(g) \mapsto \exp(-t\psi(g))\lambda(g)$ 

is the map  $A(\lambda(g)) = \psi(g)\lambda(g)$ . In particular, we find

**Radial Fourier multipliers**  $\subset$  **McIntosh's**  $H_{\infty}$ -calculus.

However, our Hörmander-Mihlin type condition above is considerably weaker.

• If  $\dim \mathcal{H}_{\psi} = n \quad \rightarrow \quad \mathcal{H}_{\psi} \simeq \mathbb{R}^n_{\text{disc}} \text{ and } \mathcal{L}(\mathcal{H}_{\psi}) \simeq L_{\infty}(\widehat{\mathbb{R}}^n_{\text{disc}}, \mu) \quad \rightarrow \quad \text{de Leeuw.}$ 

#### **Riesz transforms**

Given a discrete group  $\mathrm{G},$  we have seen

$$\begin{array}{ccc} \psi: \mathcal{G} \to \mathbb{R}_+ \\ \text{length function} \end{array} \Rightarrow \begin{array}{ccc} (\mathcal{H}_{\psi}, \langle \cdot, \cdot \rangle_{\psi}) \\ \text{Hilbert space} \end{array} \Rightarrow \begin{array}{ccc} b_{\psi}: \mathcal{G} \to \mathcal{H}_{\psi} \\ \text{cocyle map.} \end{array}$$

Thus, we consider the  $\eta$ -th Riesz  $\psi$ -transform for  $\eta \in \mathcal{H}_{\psi}$  as

$$R_{\eta}\Big(\sum_{g\in \mathcal{G}}\widehat{f}(g)\lambda(g)\Big) = -i\sum_{g\in \mathcal{G}}\frac{\langle b_{\psi}(g),\eta\rangle_{\psi}}{\sqrt{\psi(g)}}\widehat{f}(g)\lambda(g).$$

The lifting multiplier  $\tilde{m}_{\eta}(\xi) = -i \frac{\langle \xi, \eta \rangle_{\psi}}{\sqrt{\langle \xi, \xi \rangle_{\psi}}}$  only satisfies the classical condition

 $\left|\partial_{\xi}^{\beta} \tilde{m}_{\eta}(\xi)\right| \leq c_n |\xi|^{-|\beta|}$  for any multi-index  $\beta \Rightarrow \varepsilon = 0$ .

**Theorem** [JMP]. If dim  $\mathcal{H}_{\psi} < \infty$ , any operator in

 $\mathcal{R} = \operatorname{span} \left\{ \prod_{\eta \in \Gamma} R_{\eta} \mid \Gamma \text{ finite set in } \mathcal{H}_{\psi} \right\}$ defines a cb-map  $\mathcal{L}(G) \to \operatorname{BMO}_{\mathcal{S}_{\psi}}$  and  $L_p(\widehat{\mathbb{G}}, \tau) \to L_p(\widehat{\mathbb{G}}, \tau)$  for all 1 .

#### Noncommutative tori

Given 
$$n \ge 1$$
 and  $\Theta = (\theta_{jk})_{n \times n}$  antisymmetric  
 $\mathcal{A}_{\Theta} = \left\langle u_1, u_2, \dots, u_n \mid \text{unitaries with } u_j u_k = e^{2\pi i \theta_{jk}} u_k u_j \right\rangle$   
 $= \left\{ f \sim \sum_{k \in \mathbb{Z}^n} \widehat{f}(k) w_k \mid w_k = u_1^{k_1} u_2^{k_2} \cdots u_n^{k_n} \text{ with } k = (k_1, k_2, \dots, k_n) \right\}.$ 

We also need the trace  $\tau(f) = \widehat{f}(0)$  and the heat semigroup  $S_{\Theta,t}(f) = \sum_k \widehat{f}(k)e^{-t|k|^2}w_k$ .

#### Theorem [JMP]. Let

$$T_m: \sum_k \widehat{f}(k) w_k \mapsto \sum_k m_k \widehat{f}(k) w_k.$$

If a lifting multiplier  $\tilde{m}:\mathbb{R}^n\to\mathbb{C}$  with  $\tilde{m}_{\mid_{\mathbb{Z}^n}}=m$  satisfies

$$\left|\partial_{\xi}^{\beta}\tilde{m}(\xi)\right| \leq c_{n}|\xi|^{-\beta}$$
 for all  $|\beta| \leq \left[\frac{n}{2}\right] + 1$ ,

then we find that  $T_m : L_{\infty}(\mathcal{A}_{\Theta}, \tau) \to \text{BMO}_{\mathcal{S}_{\Theta}}$  and  $L_p(\mathcal{A}_{\Theta}) \to L_p(\mathcal{A}_{\Theta})$  for all 1 .

**Proof 1.** Noncommutative form of Calderón's transference from  $\mathbb{T}^n$ . **Proof 2.** We have  $\mathcal{L}(H_{\Theta}) = \int_{\mathbb{R}}^{\oplus} \mathcal{A}_{x\Theta} dx$  and apply our multiplier theorem to  $H_{\Theta}$ .

## Noncommutative Calderón-Zygmund theory

We are interested in a noncommutative form of

for any Calderón-Zygmund T with kernel k and with  $\delta_{\mathbb{R}^n}(f) = f \otimes 1_{\mathbb{R}^n} - 1_{\mathbb{R}^n} \otimes f$ .

**Major difficulty:** Construct projections playing the role of the Euclidean balls  $B_s(0)$ .

The key ingredients are

- Noncommutative BMO's over semigroups.
- An associated 'metric' on the von Neumann algebra.

# Semigroup type BMO's

Duong and Yan recently extended BMO theory to certain semigroups on homogeneous spaces assuming certain regularity. This theory, however, *still imposes the existence of a metric in the underlying space*. We may not assume the existence of a metric.

Given a noncommutative measure space  $(\mathcal{M},\tau)$  and

$$\mathcal{S} = (S_t)_{t \ge 0}$$
 with  $S_t : f \in L_p(\mathcal{M}, \tau) \to d_t * f \in L_p(\mathcal{M}, \tau)$ ,

a noncommutative diffusion semigroup of convolution type, we define

$$\|f\|_{\mathrm{BMO}_{\mathcal{S}}^{c}} = \sup_{t \ge 0} \left\| \left( S_{t} |f|^{2} - |S_{t}f|^{2} \right)^{\frac{1}{2}} \right\|_{\infty}$$

for the column semigroup type BMO. The row analog and the general form are

$$\|f\|_{\mathrm{BMO}_{\mathcal{S}}^{r}} = \|f^{*}\|_{\mathrm{BMO}_{\mathcal{S}}^{c}}$$
 and  $\|f\|_{\mathrm{BMO}_{\mathcal{S}}} = \max\left\{\|f\|_{\mathrm{BMO}_{\mathcal{S}}^{r}}, \|f\|_{\mathrm{BMO}_{\mathcal{S}}^{c}}\right\}.$ 

**Theorem** [Junge/Mei]. If S admits a 'nice enough' Markov dilation  $[BMO_{S}, L_{p}(\mathcal{M}, \tau)]_{p/q} \simeq L_{q}(\mathcal{M}, \tau).$ 

Remark. The regularity assumed in the result above holds in our main examples.

## Noncommutative 'metrics'

#### A weighted spectral decomposition for

 $\mathcal{S} = (S_t)_{t \geq 0}$  with  $S_t f = d_t * f$ 

is a family of projections  $(q_{k,t})$  in  $\mathcal{M}$  —indexed by  $(k,t) \in \mathbb{N} \times \mathbb{R}_+$ — which are increasing in k for t fixed, together with a family of positive numbers  $\beta_{k,t} \in \mathbb{R}_+$  such that the following conditions hold for absolute constants  $c_w, c_s, c_d$ 

$$\begin{array}{l} \textbf{i)} \sum_{k \geq 1} \beta_{k,t} \tau(q_{k,t}) \leq c_s, \\ \textbf{ii)} \ d_t \leq c_d \sum_{k \geq 1} \beta_{k,t} (q_{k,t} - q_{k-1,t}), \\ \textbf{iii)} \ \sum_{k \geq 1} \beta_{k,t} w_{k,t} \tau(q_{k,t} - q_{k-1,t}) \leq c_w \quad \text{for} \quad w_{k,t} = \big( \sum_{j \leq k} \sqrt{\frac{\tau(q_{j+1,t})}{\tau(q_{j,t})}} \big)^2. \end{array}$$

This notion is somehow related to

- Tolsa's notion of RBMO space for nondoubling measures.
- Blunck/Kunstmann's analysis of non-integral Calderón-Zygmund operators.

We will however require a doubling property of the trace  $\tau$ 

 $\tau(q_{\alpha(k),t}) \leq c_{\alpha}\tau(q_{k,t})$  for some strictly increasing function  $\alpha: \mathbb{N} \to \mathbb{N}$ .

### Boundedness of noncommutative CZO's

Taking  $Q_{k,t}(f) = \frac{1}{\tau(q_{k,t})} q_{k,t} * f$  yields a metric type BMO, called BMO<sub>Q</sub>.

**Theorem** [JMP]. Let  $(\mathcal{M}, \tau)$  be a noncommutative measure space and S a semigroup acting on it equipped with an  $\alpha$ -doubling weighted decomposition with associated metric  $\mathcal{Q} = (Q_{k,t})$ . Let  $T : \mathcal{A} \to \mathcal{M}$  defined on a weakly dense \*-subalgebra of  $\mathcal{M}$ . If we consider the derivation  $\delta_{\mathcal{M}}(f) = f \otimes \mathbf{1}_{\mathcal{M}} - \mathbf{1}_{\mathcal{M}} \otimes f$ , the conditions

a) 
$$T: L_2(\mathcal{M}, \tau) \rightarrow L_2(\mathcal{M}, \tau)$$
 is bounded by  $c_{22}$ ,

**b1)**  $\|\mathcal{R}_{q_{k,t}\otimes q_{k,t}}\delta_{\mathcal{M}}(T\otimes id_{\mathcal{M}})\mathcal{R}_{q_{\alpha(k),t}^{\perp}}:\mathcal{M}\bar{\otimes}\mathcal{M}\to\mathcal{M}\bar{\otimes}\mathcal{M}\bar{\otimes}\mathcal{M}\| \leq c_{h} \text{ for all } k,t,$  **b2)**  $\|\mathcal{L}_{q_{k,t}\otimes q_{k,t}}\delta_{\mathcal{M}}(T\otimes id_{\mathcal{M}})\mathcal{L}_{q_{\alpha(k),t}^{\perp}}:\mathcal{M}\bar{\otimes}\mathcal{M}\to\mathcal{M}\bar{\otimes}\mathcal{M}\bar{\otimes}\mathcal{M}\| \leq c_{h} \text{ for all } k,t,$ imply that  $T:\mathcal{A}\to BMO_{\mathcal{Q}}.$  More concretely, we obtain  $\|Tf\|_{BMO_{\mathcal{Q}}}\leq (2c_{22}\sqrt{c_{\alpha}}+c_{h})\|f\|_{\infty}.$  $\|Tf\|_{BMO_{\mathcal{S}}}\leq 2\sqrt{2}\sqrt{c_{d}(c_{s}+c_{w})}(2c_{22}\sqrt{c_{\alpha}}+c_{h})\|f\|_{\infty}.$ 

**Corollary** [JMP]. Additionally, if S has a nice Markov dilation, we obtain  $L_p$ -boundedness.

**Remark.** The heat semigroup reconstructs the classical  $\mathbb{R}^n$ -theory from Theorem above.

#### **Applications and examples**

- New  $L_{\infty} \to BMO$  Schur multipliers.
- Analysis of some **concrete groups**:  $\mathbb{Z}_n, \mathcal{S}_n, \mathbb{F}_n$ ...
- Burnside groups:  $\operatorname{H-dim}(B(n,m)) = \infty$  for  $n \ge 2$  and  $m \ge 665$  odd.
- Calderón's transference method for quantum groups.
- An adapted Littlewood-Paley theory.

# Thanks for listening!!