

# Large-time asymptotic behavior of multi-cut solutions to Hele-Shaw flows

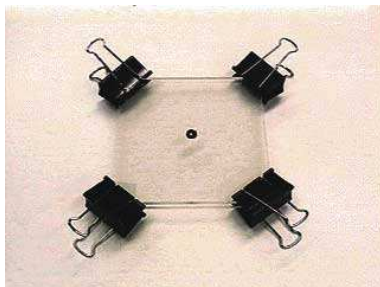
Yu-Lin Lin<sup>1</sup>

Joint work with B. Gustafsson

Integrable and stochastic Laplacian growth in modern mathematical physics

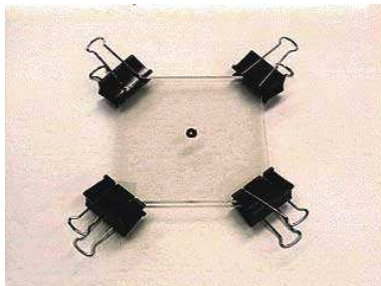
November 04, 2010

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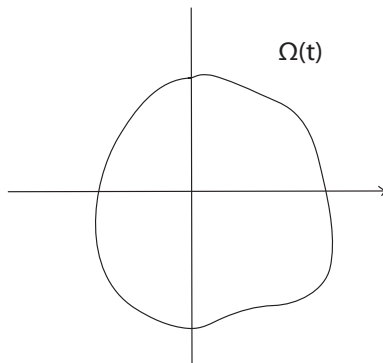
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# The case we consider

- ▶ Inject **viscid fluid (glycerol)** into **inviscous fluid (colored water)** slowly.
- ▶ The boundary is getting better.

# The moving domain

- $p(z, t)$  : pressure at  $z \in \Omega(t)$ .
- $\kappa(z, t)$  : curvature at  $z \in \partial\Omega(t)$ .
- $v_n(z, t)$  : normal velocity at  $z$  on  $\partial\Omega(t)$ .
- $\gamma$  : surface tension.
- $n$  : unit normal.
- $Q$ : injection rate.





# The free boundary problem

$$\begin{cases} \Delta p = -Q\delta_0 & \text{in } \Omega(t), \\ p = \gamma\kappa & \text{on } \partial\Omega(t), \\ v_n = -\frac{\partial p}{\partial n} & \text{on } \partial\Omega(t). \end{cases}$$

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- ▶ Consider the problem  $\gamma = 0$  now.

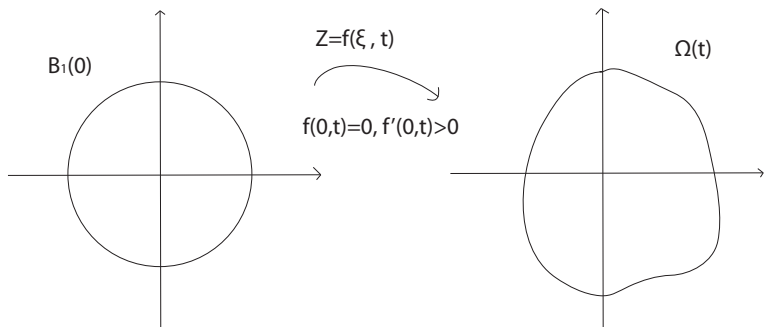
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- ▶ Consider the problem  $\gamma = 0$  now.
- ▶ Injection with speed  $Q$ ;  $\Omega(s) \subset \Omega(t)$  if  $s < t$ .

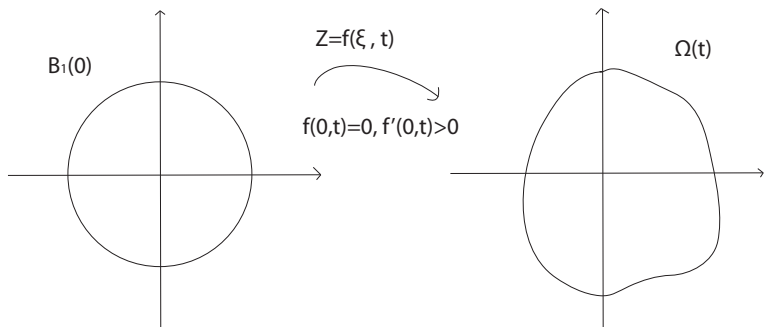
P. Ya. Polubarinova-Kochina and L. A. Galin (1945) gave a conformal formulation of the Hele-Shaw problem.

# Reformulation by Riemann mapping



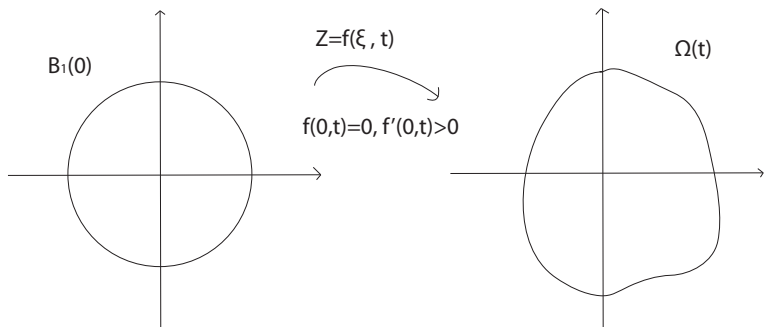
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- ▶ Definition:  
 $O(E) = \{f(\zeta) \mid f(\zeta) \text{ is univalent in } E, f(0) = 0 \text{ and } f'(0) > 0\}$ .

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- ▶  $f(\zeta, t) \in \overline{O(B_1(0))}$ .

# The Polubarinova-Galin equation

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$$\operatorname{Re}\left[\frac{d}{dt}f(\zeta,t)\overline{f'(\zeta,t)\zeta}\right] = \frac{Q}{2\pi}, \zeta \in \partial B_1(0), f(\zeta,t) \in O(\overline{B_1(0)}).$$



B. Gustafsson (1984) gave a new formulation of the P-G equation to be a Löwner-Kufarev type equation.

For  $f(\zeta, t) \in O(\overline{B_1(0)})$

$$f_t(\zeta, t) = \frac{Q}{2\pi} \frac{f'(\zeta, t)\zeta}{2\pi i} \int_{\partial B_1(0)} \frac{1}{|f'(z, t)|^2} \frac{z+\zeta}{z-\zeta} \frac{dz}{z}, |\zeta| < 1.$$

# Definition of a strong solution to the P-G equation

$$\operatorname{Re}\left[\frac{d}{dt}f(\zeta, t)\overline{f'(\zeta, t)\zeta}\right] = \frac{Q}{2\pi}, \zeta \in \partial B_1(0).$$

## ► Definition

A solution  $f(\zeta, t)$  is a strong solution of the P-G equation if  $f(\zeta, t) \in O(\overline{B_1(0)})$  is continuously differentiable with respect to  $t$  in  $[0, \epsilon)$ .

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## ▶ Definition

If a strong solution  $f(\zeta, t)$  fails to exist at  $t = T_0$ , we say the strong solution  $f(\zeta, t)$  blows up at  $t = T_0$ .

- ▶ B. Gustafsson (1984) found a general set of solutions

$$f(\zeta, t) = \sum_{j=1}^m d_j(t) \zeta^j + d_0(t) + \sum_{l=1}^n \sum_{k=1}^{s_l} \frac{a_{l,k}(t)}{(\zeta - \zeta_l(t))^k}$$

# Two special general solutions

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- ▶ Abanov, Ar. and Mineev-Weinstein, M. and Zabrodin, A. (2009) found multi-cut solutions

$$f(\zeta, t) = \sum_{j=1}^m d_j(t) \zeta^j + d_0(t) + \sum_{l=1}^n \sum_{k=1}^{s_l} \frac{a_{l,k}(t)}{(\zeta - \zeta_l(t))^k} + \sum_{l=1}^{m_0} e_l \ln(\zeta - \zeta_{-l}(t)),$$

where  $e_l$  are constant.

# The Richardson complex moments

- ▶ Given  $\Omega(t)$  which solves the problem, then the Richardson complex moments are

$$M_k(t) = \frac{1}{\pi} \int_{\Omega(t)} z^k dx dy, z = x + iy, k \geq 0.$$

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- ▶ If  $\Omega(t) = f_{k_0}(B_1(0), t)$  where  $f_{k_0}(\zeta, t) = a_1(t)\zeta + \dots + a_{k_0}(t)\zeta^{k_0}$  is a polynomial strong solution,

$$\begin{aligned} M_k(f_{k_0}(\zeta, t)) &= \frac{1}{2\pi i} \int_{\partial B_1(0)} f_{k_0}^k(\zeta, t) f'_{k_0}(\zeta, t) \overline{f_{k_0}(\zeta, t)} d\zeta \\ &= \sum_{i_1, \dots, i_{k+1}} i_1 a_{i_1}(t) a_{i_2}(t) \cdots a_{i_{k+1}}(t) \overline{a_{i_1 + \dots + i_{k+1}}(t)}. \quad (1) \end{aligned}$$

$$M_{k_0}, M_{k_0+1}, \dots = 0.$$



- ▶ Some solutions blow up

# The strong global solutions

- ▶ Some solutions blow up
- ▶ Some solutions are global.

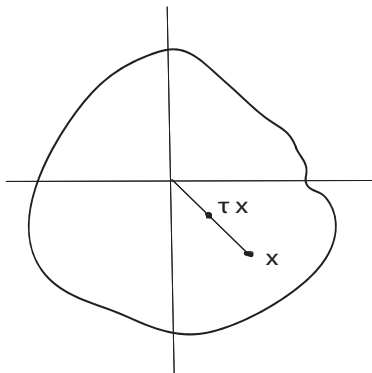
Now assume  $Q = 2\pi$

# The global dynamics

B. Gustafsson, D. Prokhorov, and A. Vasilev(2004)

## Theorem

*If  $f(\zeta, 0)$  is a starlike function and  $f(\zeta, 0) \in O(\overline{B_1(0)})$ , then the strong solution is global and  $f(\zeta, t)$  is starlike forever.*



M.Sakai (1998)

## Theorem

*If  $\Omega(0) \subset B_R(0)$ , and  $t$  is large, then*

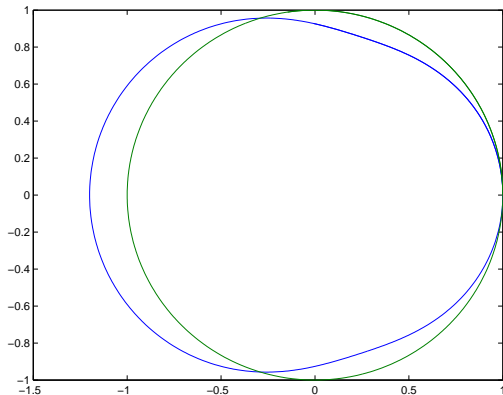
$$B_{\sqrt{(|\Omega(0)|/\pi+2t)}-R} \subset \Omega(t) \subset B_{\sqrt{(|\Omega(0)|/\pi+2t)}+R}.$$

E. Vondenhoff(2008):

$\Omega(0)$  is a small perturbation of  $B_R(0)$  where  $|\Omega(0)| = |B_R(0)|$ .

Then

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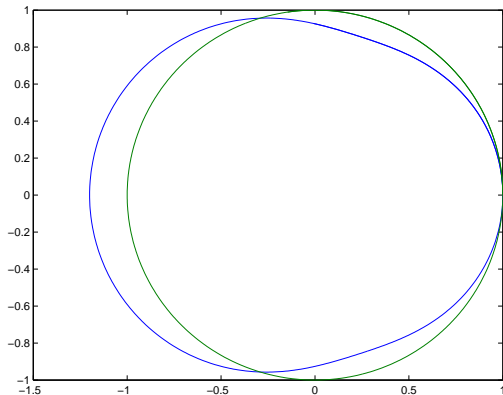


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- ▶ The solution  $\Omega(t)$  is global.
- ▶ A rescaling behavior is described in terms of moments.



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- ▶ But I only assume solutions are global and more details result about coefficients of solutions are obtained.

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## (Yu-Lin Lin) Large-time rescaling behaviors of Stokes and Hele-Shaw flows driven by injection

- ▶ Assume  $f_{k_0}(t)$  is a global **polynomial solution**.
- ▶ Understand how each coefficient decays and grows in terms of moments.
- ▶ Obtain precise large-time rescaling behavior in terms of moments.

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# The decay of polynomial coefficients

Let  $f_{k_0} = a_1(t)\zeta + \cdots + a_{k_0}(t)\zeta^{k_0}$  be a global polynomial solution.

▶ •  $M_{k-1}(t) = M_{k-1}(0), k \geq 2$

$$\Leftrightarrow \sum_{i_1, \dots, i_k} i_1 a_{i_1}(t) \cdots a_{i_k}(t) \overline{a_{i_1 + \dots + i_k}(t)} = M_{k-1}(0), k \geq 2.$$



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$$a_1(t) \approx \sqrt{2t}, \quad |a_k(t)|, k \geq 2 \quad \text{are bounded.}$$

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$$M_{k-1} = a_1^k(t) \overline{a_k(t)} + \sum_{(i_1, \dots, i_k) \neq (1, \dots, 1)} i_1 a_{i_1}(t) \cdots a_{i_k}(t) \overline{a_{i_1 + \dots + i_k}(t)}$$

# The decay of polynomial coefficients

- ▶ Then by induction, we can get

$$a_k(t) \left( a_1^k(t) \right) = \overline{M_{k-1}} + O\left( \frac{1}{a_1^4(t)} \right), 2 \leq k \leq k_0.$$

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$$a_k(t) \left( a_1^{(n_0+1)}(t) \right) = o(1), 2 \leq k \leq n_0.$$

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$$\lim_{t \rightarrow \infty} \left[ f(\zeta, t) - \sqrt{2t + M_0(0)\zeta} \right] (\sqrt{2t})^{n_0+1} = \overline{M_{n_0}} \zeta^{n_0+1} \neq 0 \quad (2)$$

## Theorem

Let  $f(\zeta, t) = \frac{\sum_{j=1}^m b_j \zeta^j}{\prod_{l=1}^n (\zeta - \zeta_l)^{s_l}} + \sum_{l=1}^{m_0} e_l \ln(\zeta - \zeta_{-l}) + b_0 = \sum_{j=1}^{\infty} a_j \zeta^j$  be a global solution. Assume  $n_0 = \min\{k \geq 1 | M_k \neq 0\}$ . Then we have

1.

$$a_k(t) \left( a_1^k(t) \right) = \overline{M_{k-1}} + O\left(\frac{1}{a_1^4(t)}\right), 2 \leq k < \infty.$$

$$a_k(t) \left( a_1^{(n_0+1)}(t) \right) = o(1), 2 \leq k \leq n_0.$$

2.

$$\lim_{t \rightarrow \infty} \max_{\zeta \in \partial B_1(0)} \left| \left[ f(\zeta, t) - \sqrt{2t + M_0(0)\zeta} \right] \left( \sqrt{2t} \right)^{n_0+1} - \overline{M_{n_0}} \zeta^{n_0+1} \right| = 0 \quad (3)$$

3. Let  $n_1 = \{\zeta_j | \zeta_j \text{ singularity of } f(\zeta, 0)\}$ . Then  $n_0 \leq m + n_1 - 1$ .

Denote  $f = \sum_{j=1}^{\infty} a_j \zeta^j$  and  $f_k = \sum_{j=1}^k a_j \zeta^j$ .

$$\max_{\zeta \in B_1(0)} |f_k(\zeta, t) - f(\zeta, t)| = O(a_1(t)^{-i_k}) \quad (4)$$

and

$$\max_{\zeta \in B_1(0)} |f'_k(\zeta, t) - f'(\zeta, t)| = O(a_1(t)^{-i_k}) \quad (5)$$

where  $i_k \rightarrow \infty$  as  $k \rightarrow \infty$ .

In this case, there exists  $f_{l_k}$  such that

$$\begin{aligned}
 M_{k-1} &= \frac{1}{2\pi i} \int_{\partial B_1(0)} f^k(\zeta, t) f'(\zeta, t) \overline{f(\zeta, t)} d\zeta \\
 &= \frac{1}{2\pi i} \int_{\partial B_1(0)} f_{l_k}^k(\zeta, t) f'_{l_k}(\zeta, t) \overline{f_{l_k}(\zeta, t)} d\zeta + O\left(\frac{1}{a_1^4}\right) \\
 &= \sum_{\substack{\sum_{j=1}^k l_j \leq l_k \\ i_1, \dots, i_k}} i_1 a_{i_1}(t) a_{i_2}(t) \cdots a_{i_k}(t) \overline{a_{i_1+\dots+i_k}(t)} + O\left(\frac{1}{a_1^4}\right). \tag{6}
 \end{aligned}$$



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▶

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# Why Multi-cut solutions behave like polynomial

- ▶ Understand precise large-time behavior of singularity  $\zeta_l, \zeta_{-l}$ .

$$f(\zeta, t) = \frac{\sum_{j=1}^m b_j \zeta^j}{\prod_{l=1}^n (\zeta - \zeta_l)^{s_l}} + \sum_{l=1}^{m_0} e_l \ln(\zeta - \zeta_{-l}) + b_0 = \sum_{j=1}^{\infty} a_j \zeta^j,$$

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- ▶ Understand large-time behavior of  $b_j$ .

Tool:  $|a_1(t)| \approx \sqrt{2t}$ ,  $|a_k(t)|, k \geq 2$  are uniformly bounded

# Derive the equation for singularity

$$g = f'(\zeta, t) = \frac{P_1(\zeta, t)}{\prod(\zeta - \theta_j)^{l_j}}, \theta_j \text{ are distinct singularity of } f(\zeta, t).$$



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- ▶ Multiply  $(\zeta - \theta_j)$  and let  $\zeta - \theta_j \rightarrow 0$ . We can obtain  $\dot{\theta}_j$ .



$$f(\zeta, t) = \frac{\sum_{j=1}^m b_j \zeta^j}{\prod_{l=1}^n (\zeta - \zeta_l)^{s_l}} + \sum_{l=1}^{m_0} e_l \ln(\zeta - \zeta_{-l}) + b_0,$$

$$\frac{d}{dt}(\ln |\zeta_k|) = \frac{1}{2\pi} \int_0^{2\pi} \frac{1}{|f'(e^{i\alpha}, t)|^2} \frac{|\zeta_k|^2 - 1}{|\zeta_k - e^{i\alpha}|^2} d\alpha > 0.$$



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- ▶ There exist  $c_1, c_2 > 0$  such that

$$c_1 \leq \frac{|\zeta_k(t)|}{a_1(t)} \leq c_2, \quad t \geq 0$$

# Geometric meaning of the rescaling behavior

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$$\max_{z \in \partial\Omega'(t)} \|z\| - 1 = O\left(\frac{1}{t^{1+\frac{n_0}{2}}}\right)$$

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- ▶ The value  $1 + \frac{n_0}{2}$  is the best rate we can get.

$$\limsup_{t \rightarrow \infty} \max_{z \in \partial\Omega'(t)} \|z\|^{-1} (2t)^{1+\frac{n_0}{2}} = |M_{n_0}|.$$

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# Why $n_0 \leq m + n_1 - 1$ , $n_1 = \text{total singularity}$ ?



$$f(\zeta, t) = \frac{\sum_{j=1}^m b_j \zeta^j}{\prod_{l=1}^n (\zeta - \zeta_l)^{s_l}} + \sum_{l=1}^{m_0} e_l \ln(\zeta - \zeta_{-l}) + b_0 = \sum_{j=1}^{\infty} a_j \zeta^j.$$

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- ▶ If  $f(\zeta, t)$  is rational,  $n_0 \leq m - 1$ .