

RANDOM NORMAL MATRICES BY RIEMANN-HILBERT PROBLEM

Seung Yeop Lee (Caltech)

(Work in progress with Ferenc Balogh, Marco Bertola and Ken
McLaughlin)

Banff 2010: Integrable and stochastic Laplacian growth in modern
mathematical physics (10w5019)

November 1, 2010

NORMAL MATRICES \rightarrow COULOMB GAS

The eigenvalues of $n \times n$ normal matrices with the probability distribution

$$\text{Prob}(M)dM = \frac{1}{\mathcal{Z}} e^{-N \text{Tr}[MM^\dagger + V(M) + V(M)^\dagger]} dM,$$

distributes by the probability density

$$P(\lambda_1, \dots, \lambda_n) = \frac{1}{\mathcal{Z}} \left| \prod_{j < k} (\lambda_j - \lambda_k) \right|^2 \exp \left(-N \sum_{j=1}^n Q(\lambda_j) \right),$$

where the potential is given by

$$Q(z) = |z|^2 + V(z) + \overline{V(z)},$$

i.e. **Gaussian** plus **harmonic** function when V is holomorphic.

Eigenvalues of random normal matrices are Coulomb gas in 2-dimension.

CONTINUUM LIMIT OF COULOMB GAS: $n, N \rightarrow \infty$

Define $t = \lim_{n, N \rightarrow \infty} n/N$.

For real analytic Q ,

$$\frac{1}{n} \sum_j \delta(z - \lambda_j) \xrightarrow{\text{weak}} \rho(z) + \frac{1}{n} \rho_{1/2}(z) + (\text{fluctuation})$$

where [Wiegmann-Zabrodin, Ameur-Hedenmalm-Makarov, ...]

$$\rho = \frac{\Delta Q}{4\pi t} \mathbf{1}_S,$$

$$\rho_{1/2} = \frac{2-\beta}{8\pi\beta} \Delta((\log \Delta Q)^H + \mathbf{1}_S) = 0.$$

In our case, $\beta = 2$ and $\Delta Q = 4$.

The density is asymptotically constant on S , and $\text{Area}(S) = \pi t$.

S is determined by the **support of σ** that minimizes

$$\int_{\mathbb{C}} Q(w) \sigma(w) d^2 w - \frac{n}{N} \iint_{\mathbb{C}^2} \sigma(z) \sigma(w) \log |z - w| d^2 z d^2 w.$$

Correspondence to Hele-Shaw flow: S is the domain of non-viscous fluid in ideal Hele-Shaw flow when t is the Hele-Shaw time.

ORTHOGONAL POLYNOMIALS ON \mathbb{C}

(Joint) Probability densities are given by (for $k \leq n$)

$$P(\lambda_1, \dots, \lambda_k) \propto \det \left(K_n(\lambda_i, \lambda_j) e^{-\frac{N}{2}(Q(\lambda_i) + Q(\lambda_j))} \right)_{i,j=1}^k,$$

where the (reproducing) kernel K_n is defined by

$$K_n(z, w) = \sum_{j=0}^{n-1} \frac{p_j(z) \overline{p_j(w)}}{h_j},$$

and $p_j = x^j + \dots$ is a polynomial of degree j defined by

$$h_j \delta_{ij} = \int_{\mathbb{C}} p_i(z) \overline{p_j(z)} e^{-NQ(z)} dA(z).$$

Therefore, the kernel K_n is **the most wanted** in our analysis.
For example, the **density** is given by

$$\rho_n(z) := \frac{1}{N} K_n(z, z) e^{-NQ(z)}.$$

QUANTIZED HELE-SHAW FLOW

Hele-Shaw time t is quantized by n/N .

So it is natural to expect

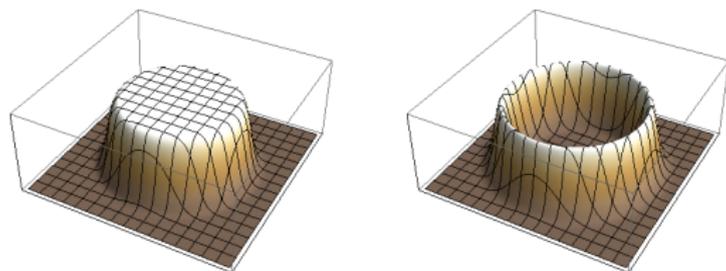
$$N(\rho_{n+1}(z) - \rho_n(z)) = \frac{|p_n(z)|^2}{h_n} e^{-NQ(z)} \sim \mathbf{1}_{\delta S}$$

where δS is the growing part of S for a small time interval.

Theorem [AHM]: $|p_n(z)|^2 e^{-NQ(z)} dA(z)$ converges to the **harmonic measure** at ∞ with respect to $\mathbb{C} \setminus S$.

THE SIMPLEST CASE: GINIBRE ENSEMBLE

When $V(z) = 0$ the orthogonal polynomials are $p_k(z) = z^k$.
The kernel is explicitly given in terms of Gamma functions.



[Left] $K_n(z, z) e^{-NQ(z)}$ for $n = N = 40$.

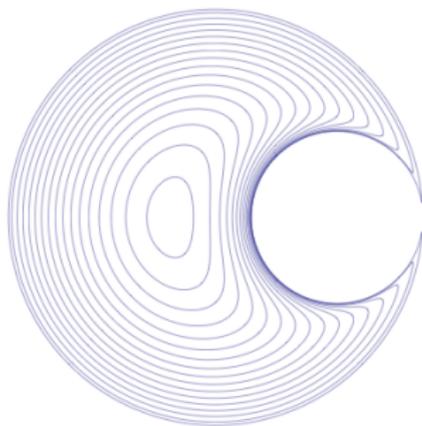
[Right] $|p_n(z)|^2 e^{-NQ(z)}$ for $n = N = 20$.

GOAL (AND SOME IMPLICATIONS)

Taking $V(z) = -c \log(z - a)$ where $c > 0$, we will obtain the **pointwise** limit of $p_n(z)$ and $K_n(z, z)$ using **Riemann-Hilbert method (DKMVZ '99)**.

- ▶ There is a topological transition.
- ▶ Similar technique may apply for any finite logarithmic singularities.
- ▶ RH method can, if necessary, give us an arbitrary order of accuracy in large N expansion.
- ▶ RH method effectively handles the singular region (for instance, where the merging transition occurs).

(Asymptotics of $p_n(z)$): Elbau-Felder)

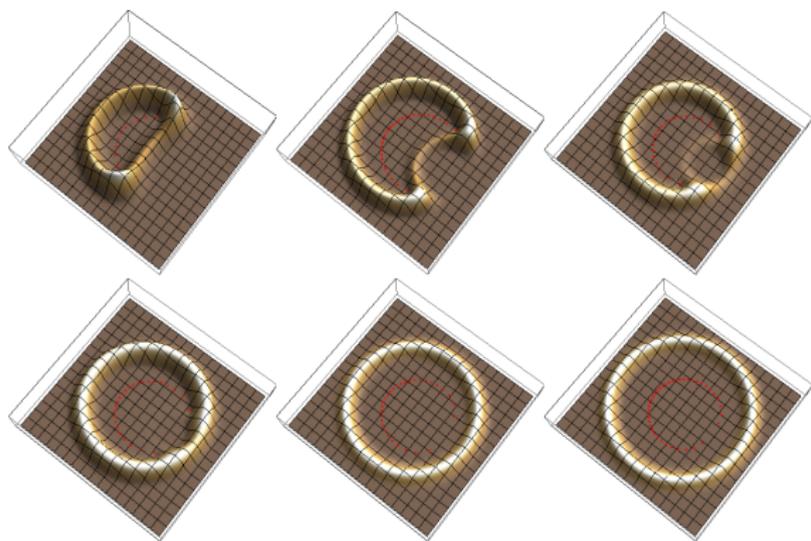


UNIVERSALITY

- ▶ The kernel in the bulk is heat kernel.
- ▶ The kernel at the boundary is written in terms of **erf**.
- ▶ Merging transition is described, in the classical limit, by scaling solution such that $y^2 \sim x^2(x^2 + t)$.
- ▶ In the merging transition, orthogonal polynomial (and quantum Hele-Shaw) is given by PII parametrix (conjectured by Bettelheim-Lee-Wiegmann, Its-Bleher in 1-matrix model)
- ▶ The spectral edges in Hermitian matrix model correspond to the cusp-singularities.

$$\boxed{(4n + 1, 2) \text{ cusp}} \Leftrightarrow \boxed{\text{density that vanishes as } \sim x^{2n+1/2}}.$$

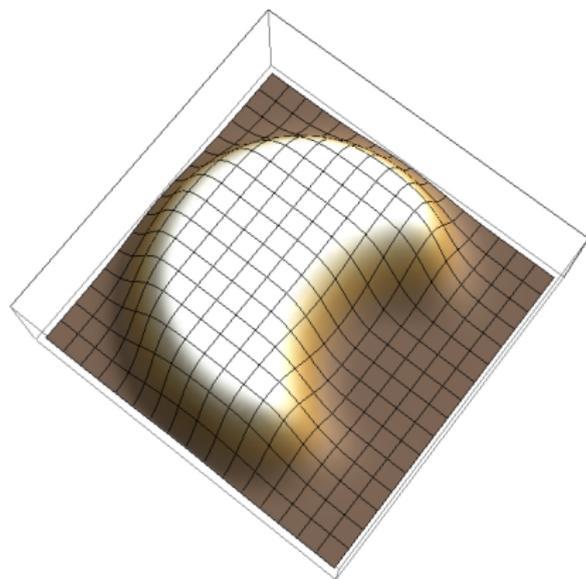
AND THE POLYNOMIALS LOOK LIKE...



The plots of $|p_n(z)|^2 e^{-NQ(z)}$ for various n .

- The roots (red dots) are on 1 dimensional curve.
- The peak is on the **growing part** of the boundary.

THE DENSITY LOOKS LIKE...



(Summing up to 12th polynomials...)

1ST STEP: AREA INTEGRAL INTO CONTOUR INTEGRAL

$$\begin{aligned}
 & \lim_{R \rightarrow \infty} \int_{|z| < R} p_j(z) (\bar{z} - a)^k |z - a|^{2Nc} e^{-Nz\bar{z}} d^2z \\
 &= \lim_{R \rightarrow \infty} \int_{|z| < R} p_j(z) (z - a)^{Nc} \frac{d}{d\bar{z}} \left(\int_a^{\bar{z}} (s - a)^{Nc+k} e^{-Nzs} ds \right) \frac{d\bar{z} \wedge dz}{2i} \\
 &= \lim_{R \rightarrow \infty} \oint_{|z|=R} p_j(z) (z - a)^{Nc} \left(\int_a^{\bar{z}} (s - a)^{Nc+k} e^{-Nzs} ds \right) \frac{dz}{2i} \\
 &\Rightarrow \lim_{R \rightarrow \infty} \frac{\Gamma(Nc + k + 1)}{2i} \oint_{|z|=R} \frac{p_j(z) (z - a)^{Nc}}{(Nz)^{Nc+k+1}} e^{-Naz} dz.
 \end{aligned}$$

In the last line, $\lim_{R \rightarrow \infty}$ can be dropped.

$$\begin{aligned}
 0 &= \oint p_j(z) \frac{(z - a)^{Nc} e^{-Naz}}{z^{Nc+j}} z^{j-k-1} dz \quad \text{for } k = 0, 1, \dots, j - 1. \\
 &= \oint p_j(z) w_n(z) z^s dz \quad \text{for } s = 0, 1, \dots, j - 1, \quad w_j(z) = \frac{(z - a)^{Nc} e^{-Naz}}{z^{Nc+j}}.
 \end{aligned}$$

$p_j(z)$ is the orthogonal polynomial with respect to $w_j(z) dz$!

RIEMANN-HILBERT PROBLEM

$$Y(z) = \begin{pmatrix} p_n(z) & \frac{1}{2\pi i} \oint \frac{p_n(z')}{z' - z} w_n(z') dz' \\ -2\pi i Q_{n-1}(z) & - \oint \frac{Q_{n-1}(z')}{z' - z} w_n(z') dz' \end{pmatrix} .$$

Here $Q_{n-1}(z)$ is the orthogonal polynomial of degree $n - 1$ with respect to the measure $w_n(z)dz$.

Then Y satisfies the **Riemann-Hilbert problem**:

$Y(z)$ is holomorphic in $\mathbb{C} \setminus \Gamma$.

$$Y_+(z) = Y_-(z) \begin{pmatrix} 1 & w_n(z) \\ 0 & 1 \end{pmatrix}, \quad z \in \Gamma .$$

$$Y(z) = \left(I + \mathcal{O} \left(\frac{1}{z} \right) \right) z^{n\sigma_3} \quad z \rightarrow \infty .$$

Above, Γ is the contour that goes around the origin.

EIGENVALUE SUPPORT: S

... CAN BE DEFINED BY THE FOLLOWING STEP.

Write a general rational function that maps ∞ and α to ∞ .

$$f(v) := \rho v + \frac{\kappa}{v - \alpha} + z_0.$$

$$S(v) := f(1/v).$$

Require that

- ▶ $S(v) \sim c/(f(v) - a)$ as $v \rightarrow 1/\alpha$,
- ▶ $S(v) \rightarrow (c + t)/f(v)$ as $v \rightarrow \infty$.

This determines $\rho, \kappa, \alpha, z_0$.

Now S is given by $f(\partial\mathbb{D})$.

The “inverse” function f^{-1} maps $\mathbb{C} \setminus S$ to the outside of **unit disk**.
 $\Rightarrow \log |f^{-1}(z)| = G(\infty, z)$. (f^{-1} can be extended inside S .)

RH METHOD GIVES...

The first column of Y : (in terms of geometric quantity)

$$p_n(z) = \sqrt{\rho(f^{-1}(z))'} e^{ng(z)} \left(1 + \mathcal{O}\left(\frac{1}{n}\right) \right),$$

$$Q_{n-1}(z) = c_n \frac{\sqrt{\rho(f^{-1}(z))'}}{\rho(f^{-1}(z) - \alpha)} e^{ng(z)} \left(1 + \mathcal{O}\left(\frac{1}{n}\right) \right).$$

- ▶ $\operatorname{Re} g(z)$ is the logarithmic potential of $\mathbf{1}_S$, i.e.

$$g(z) := \frac{1}{\pi t} \int_{\mathbf{1}_S} \log(z - \zeta) dA(\zeta).$$

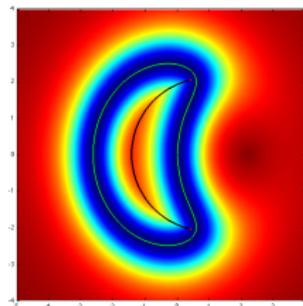
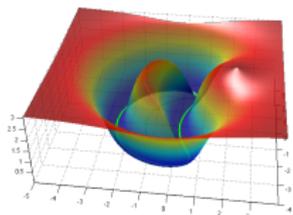
(g can be analytically extended inside S .)

- ▶ $|(f^{-1}(z))'|$ gives the harmonic measure.

$\mathbf{1}_S$ was defined by a potential problem such that

$$Q(z) - 2t \operatorname{Re} g(z) = |z|^2 + 2 \operatorname{Re} \left(-c \log(z - a) - t g(z) \right)$$

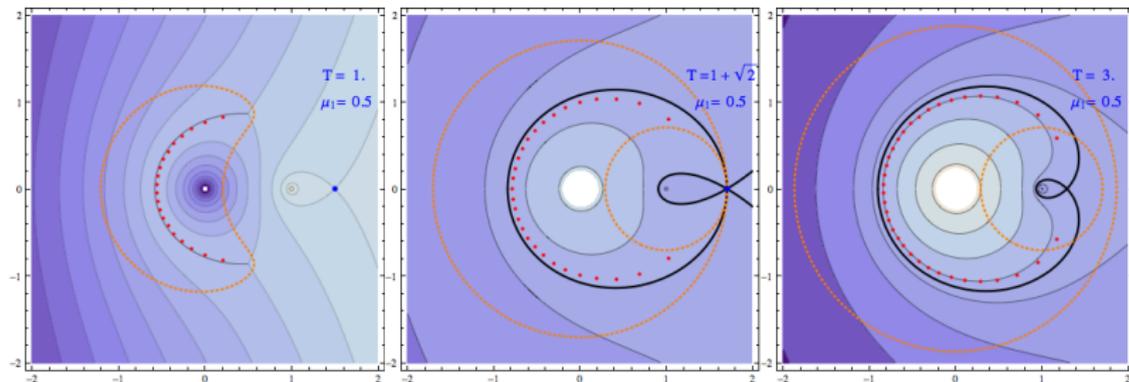
gets minimized on S (or on ∂S).



$$\Rightarrow |p_n(z)|^2 e^{-nQ(z)} = \rho |(f^{-1}(z))'| e^{-n(Q(z) - 2t \operatorname{Re} g(z))} (1 + \mathcal{O}(n^{-1}))$$

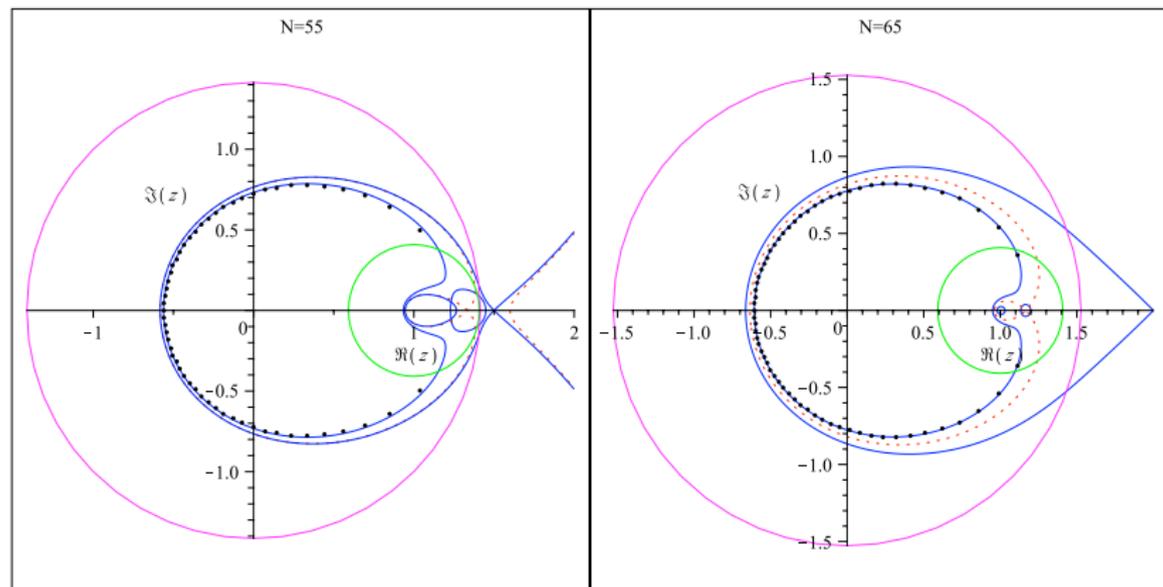
is peaked on ∂S and the line density proportional to the harmonic measure.

ZERO LOCUS AND S (ORANGE DOTTED LINE).



⇒ At the edge of the zero locus, special functions are used to describe the strong asymptotics (Airy function, Painlevé II function, Parabolic cylinder function).

BETTER APPROXIMATION



— Parabolic cylinder function.

CHRISTOFFEL-DARBOUX IDENTITY

When applying “Riemann-Hilbert technique” to 1 matrix model, one obtains the asymptotic expansion of $p_j(z)$'s and obtains the kernel by using the identity

$$\sum_{j=0}^{n-1} \frac{p_j(z)p_j(w)}{h_j} \propto \frac{p_n(z)p_{n-1}(w) - p_n(w)p_{n-1}(z)}{z - w}.$$

This relates the large sum into a few highest degree polynomials. In normal matrix model (\in 2 matrix model) CD identity exists but does not relate to $K_n(z, w)$ directly.

CD IDENTITY FOR THE BI-ORTHOGONAL POLYNOMIAL

Let us consider the following (normalized) BOP with polynomial potentials V and W . (Our case: $W(w) \rightarrow \bar{V}(\bar{w}); w \rightarrow \bar{z}$)

$$\delta_{nm} = \iint dz dw p_n(z) q_m(w) e^{-N(zw + V(z) + W(w))}.$$

Define $\tilde{p}_m(z) := p_m(z) e^{-NV(z)}$; $\tilde{q}_m(w) := q_m(w) e^{-NW(w)}$.

Differentiating inside the integral,

$$\begin{aligned} 0 &= \iint dz dw \frac{d}{Ndw} (\tilde{p}_n(z) \tilde{q}_m(w) e^{-Nzw}) \\ &= \iint \tilde{p}_n(z) \frac{d\tilde{q}_m(w)}{Ndw} e^{-Nzw} - \iint \tilde{p}_n(z) z \tilde{q}_m(w) e^{-Nzw}. \end{aligned}$$

- ▶ Obviously, $z \tilde{\mathbf{p}}_n(z)$ is spanned by $\{\tilde{p}_0(z), \dots, \tilde{p}_{n+1}(z)\}$.
- ▶ $\frac{d}{dz} \tilde{q}_m(z)$ is a linear combination of $\{\tilde{q}_{m+d}(z), \dots, \tilde{q}_0(z)\}$ where d is the degree of $W'(z)$.

(This is not true if $W'(w)$ is **not polynomial**.)

(See, “Biorthogonal polynomials for two-matrix models with semiclassical potentials” by Bertola)

We consider the truncated vectors:

$$\begin{aligned}\tilde{\mathbf{p}}_n(z) &:= (\tilde{p}_{n-1}(z), \dots, \tilde{p}_0(z))^T, \\ \tilde{\mathbf{q}}_n(w) &:= (\tilde{q}_{n-1}(w), \dots, \tilde{q}_0(w))^T.\end{aligned}$$

The kernel is given by $\tilde{\mathbf{q}}_n(w)^T \tilde{\mathbf{p}}_n(z)$.

We have,

$$z \tilde{\mathbf{p}}_n(z) = \left(\iint d\zeta dw e^{-N\zeta w} \tilde{\mathbf{p}}_n(\zeta) \zeta \tilde{\mathbf{q}}_n(w)^T \right) \tilde{\mathbf{p}}_n(z) + a_n(\tilde{p}_n(z), 0, \dots, 0)^T$$

where **the red part** is the projection operation into $\{\tilde{p}_0, \dots, \tilde{p}_{n-1}\}$
(The big bracket is an $n \times n$ matrix).

Similarly,

$$\frac{d}{Nd w} \tilde{\mathbf{q}}_n(w)^T = \tilde{\mathbf{q}}_n(w)^T \left(\iint dz d\zeta e^{-Nz\zeta} \tilde{\mathbf{p}}_n(z) \frac{d\tilde{\mathbf{q}}_n(\zeta)^T}{Nd\zeta} \right) + (\text{mostly zero vector}).$$

Note that the **same $n \times n$ matrix** appears.

$\tilde{\mathbf{q}}_n(w)^T \times$ (1st eq.) $-$ (2nd eq.) $\times \tilde{\mathbf{p}}_n(z)$ gives

$$\begin{aligned} \left(z - \frac{d}{Nd w} \right) \tilde{\mathbf{q}}_n(w)^T \tilde{\mathbf{p}}_n(z) &= -\frac{e^{Nzw}}{N} \frac{d}{dw} \left(e^{-Nzw} \tilde{\mathbf{q}}_n(w)^T \tilde{\mathbf{p}}_n(z) \right) \\ &= (\text{involving a few highest degree polynomials}). \end{aligned}$$

LET'S WORK... REMIND $w_n(z) = (z - a)^{Nc} e^{-Naz} / z^{Nc+n}$

We have defined $Q_{n-1}(z) = c_n z^{n-1} + \dots$ such that

$$\oint Q_{n-1}(z) w_n(z) z^k dz = \delta_{k,n-1} \quad \text{for } k = 0, \dots, n-1.$$

We define $(p_n(z) = z^n + b_n z^{n-1} + \dots$ is monic polynomial)

$$\tilde{h}_j := \oint p_n(z) p_n(z) w_n(z) dz.$$

For $k \leq n-1$, we have

$$\begin{aligned} & \int (p_n(z) - z p_{n-1}(z)) w_n(z) z^k dz \\ &= \int p_n(z) w_n(z) z^k dz - \int p_{n-1}(z) w_{n-1}(z) z^k dz = -\tilde{h}_{n-1} \delta_{k,n-1} \end{aligned}$$

Therefore, we have $c_n = (b_{n-1} - b_n) / \tilde{h}_{n-1}$ and

$$Q_{n-1}(z) = \frac{z p_{n-1}(z) - p_n(z)}{\tilde{h}_{n-1}}.$$

Similarly, consider the following polynomial of degree $n - 1$ for $k \leq n - 1$:

$$\begin{aligned} & \int w_n(z) z^k (p_n(z) - c_{n+1}^{-1} Q_n(z)) dz \\ &= \int w_n(z) z^k p_n(z) dz - c_{n+1}^{-1} \int w_n(z) z^k Q_n(z) dz \\ &= \frac{-1}{c_{n+1} \tilde{h}_n} \int w_n(z) z^k (z p_n(z) - p_{n+1}(z)) dz \\ &= \frac{-1}{c_{n+1} \tilde{h}_n} \left(\int w_n(z) z^{k+1} p_n(z) dz - \int w_{n+1}(z) z^{k+1} p_{n+1}(z) dz \right) \\ &= -c_{n+1}^{-1} \delta_{k, n-1} \end{aligned}$$

Therefore,

$$p_n(z) - c_{n+1}^{-1} Q_n(z) = -c_{n+1}^{-1} Q_{n-1}(z).$$

Combined with the previous equation \Rightarrow **Three-term recurrence.**

DIFFERENTIAL RELATION

Note that $\frac{d}{dz}\tilde{p}_n(z)$ is not a linear combination of $\{\tilde{p}_j\}$, but

$$\frac{1}{N} \frac{d}{dz} \left(\tilde{p}_n(z) - \frac{p_n(a)}{p_{n+1}(a)} \tilde{p}_{n+1}(z) \right) \text{ is.}$$

Even more, this is orthogonal to $\{\tilde{p}_0, \dots, \tilde{p}_{n-2}\}$ with respect to the area integral, hence

$$\frac{1}{N} \frac{d}{dz} \left(\tilde{p}_n(z) - \frac{p_n(a)}{p_{n+1}(a)} \tilde{p}_{n+1}(z) \right) = \star \tilde{p}_{n-1}(z) - \frac{p_n(a)}{p_{n+1}(a)} \left(c + \frac{n+1}{N} \right) \tilde{p}_n(z).$$

\star is obtained after some algebra to be

$$\star = \frac{P_n(a)}{P_{n-1}(a)} \frac{c_n}{c_{n+1}}.$$

KERNEL

$$\begin{aligned} e^{-Nzw} \left(z - \frac{1}{N} \frac{d}{dw} \right) \sum_{j=0}^{n-1} \frac{\tilde{p}_j(w) \tilde{p}_j(z)}{h_j} &= -\frac{1}{N} \frac{d}{dw} \left(e^{-Nzw} \sum_{j=0}^{n-1} \frac{\tilde{p}_j(w) \tilde{p}_j(z)}{h_j} \right) \\ &= \frac{e^{-Nzw}}{c_{n+1} h_n} \left(-\frac{1}{N} \frac{d}{dw} \tilde{p}_n(w) \tilde{Q}_{n-1}(z) + \star \tilde{p}_{n-1}(w) \tilde{Q}_n(z) \right). \end{aligned}$$

The kernel is obtained by the antiderivative. Especially the density $\rho_n(z) = e^{-N|z|^2} K_n(z, z)$ can be obtained by integrating

$$\partial_z \rho_n(z) = \frac{e^{-N|z|^2}}{c_{n+1} h_n} \left(-\frac{1}{N} \partial_z \tilde{p}_n(z) \tilde{Q}_{n-1}(\bar{z}) + \star \tilde{p}_{n-1}(z) \tilde{Q}_n(\bar{z}) \right).$$

This quantity converges to “ $\partial_z \mathbf{1}_S$ ”.



THANKS FOR THE ATTENTION.