

# Fluctuation bounds for interface free energies of spin glasses

Louis-Pierre Arguin  
Université de Montréal

*joint with*

C. Newman (NYU), D. Stein (NYU) & J. Wehr (Arizona)

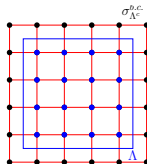
*Spin Glasses and Related Topics*  
BIRS, July 22 2014

# The Edwards-Anderson Model and the RFIM

Consider  $\Lambda \subset \mathbb{Z}^d$  a finite box and  $E(\Lambda)$  the corresponding edges.

- ▶ **Ising spin glass** or **EA model** for  $\sigma \in \{-1, +1\}^\Lambda$

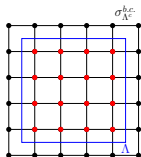
$$H_{\Lambda, \omega}(\sigma) = - \sum_{(x, y) \in E(\Lambda)} \omega_{xy} \sigma_x \sigma_y - \sum_{x \in \Lambda, y \in \Lambda^c} \omega_{xy} \sigma_x \sigma_y^{b.c.}$$



$\omega = (\omega_{xy}; (x, y) \in E(\mathbb{Z}^d))$  i.i.d. with continuous distribution  $\mathbb{P}$ .

- ▶ **Random Field Ising Model (RFIM)**

$$\tilde{H}_{\Lambda, \omega}(\sigma) = - \sum_{(x, y) \in E(\Lambda)} J \sigma_x \sigma_y - \sum_{x \in \Lambda, y \in \Lambda^c} J \sigma_x \sigma_y^{b.c.} - \sum_{x \in \Lambda} \omega_x \sigma_x$$



$(\omega_x, x \in \mathbb{Z}^d)$  are i.i.d. random variables.

## Variance bounds for the Free Energy

We want bounds on the fluctuations of the free energy:

$$F_\Lambda(\omega) = \log \sum_{\sigma \in \{-1, +1\}^\Lambda} \exp -\beta H_{\Lambda, \omega}(\sigma) \quad \beta > 0$$

Theorem (Wehr-Aizenman '90, Chatterjee '09)

*For the EA and RFIM model on  $\mathbb{Z}^d$ ,*

$$\text{Var } F_\Lambda(\omega) \geq C(\beta) |\Lambda|$$

### Our goals

1. Study the fluctuations of free energy difference between b.c.  
**Interface Free Energy**
2. Understand the impact on the structure of the Gibbs states.

## Interface Free Energy

Let  $\Gamma, \Gamma' \in \mathcal{G}_\omega(\beta)$ , two Gibbs states at disorder  $\omega$  and inv. temp.  $\beta$ .

**DLR equations**

Interface free energy:

$$F_\Lambda(\omega, \Gamma, \Gamma') = \log \Gamma \left( \exp \beta H_{\Lambda, \omega}(\sigma_\Lambda, \sigma_{\Lambda^c}) \right) - \log \Gamma' \left( \exp \beta H_{\Lambda, \omega}(\sigma_\Lambda, \sigma_{\Lambda^c}) \right)$$

- ▶ By DLR, this reduces to

$$F_\Lambda(\omega, \Gamma, \Gamma') = \log \frac{\Gamma \left( Z_{\Lambda, \omega}^{-1}(\beta, \sigma_{\Lambda^c}) \right)}{\Gamma' \left( Z_{\Lambda, \omega}^{-1}(\beta, \sigma'_{\Lambda^c}) \right)}$$

- ▶ At  $\beta = \infty$ , the analogue is the difference of energies of  $\sigma$  and  $\sigma'$  in  $\Lambda$  ground states in  $\mathbb{Z}^d$  at disorder  $\omega$ .

## RFIM: The Aizenman-Wehr theorem

Theorem (Aizenman-Wehr '90)

*In  $d = 2$ , for all  $\beta > 0$ , there is only one Gibbs state of the RFIM:*

$$\boxed{\#\mathcal{G}_\omega(\beta) = 1 \quad \omega\text{-a.s.}}$$

*Rounding of phase transition*

- ▶ The proof is partially based on the argument by Imry & Ma '85.
- ▶ In  $d \geq 3$ , RFIM exhibits a phase transition (Imbrie '85, Bricmont & Kupiainen '87).

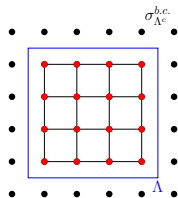
## Variance bounds and Gibbs States

In  $d = 2$ , absence of phase transition in AW is shown by **contradiction**:

### UPPER bound

$|F_\Lambda(\omega, \Gamma, \Gamma')| \leq C|\partial\Lambda|$   $\omega$ -a.s. for any  $\Gamma, \Gamma'$

The r.v.  $|F_\Lambda(\omega, \Gamma, \Gamma')|/|\partial\Lambda|$  is bounded.



### Variance LOWER bound

- ▶ The r.v.  $F_\Lambda/|\Lambda|^{1/2}$  is unbounded

▶

$$\liminf_{\Lambda \uparrow \mathbb{Z}^d} \mathbb{E}[\exp t F_\Lambda/|\Lambda|^{1/2}] \geq e^{Ct^2} \sim \text{Martingale CLT}$$

- ▶  $\text{If } \Gamma \neq \Gamma', \text{ Var } F_\Lambda(\omega, \Gamma, \Gamma') \geq C|\Lambda|$
- ▶ Domination by the (+)-b.c. and the (-)-b.c. states:

$$\Gamma_\omega^-(f(\sigma)) \leq \Gamma_\omega(f(\sigma)) \leq \Gamma_\omega^+(f(\sigma)) \quad \text{FKG inequality}$$

- ▶ It suffices to show  $\Gamma_\omega^- = \Gamma_\omega^+$  for  $\#\mathcal{G}_\omega(\beta) = 1$ .

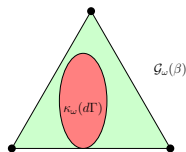
## EA: choosing a Gibbs state

For EA

- ▶ there is no domination of (+) and (-) states.
- ▶ there might be more than one limit state for given b.c.

Way out: sample a state using **metastate**:

$$\kappa_\omega(d\Gamma) \text{ prob. measure on } \mathcal{G}_\omega(\beta)$$



such that

1. **Coupling covariance**: If  $\omega_B = 0$  except on edges in a box  $B$

$$\kappa_{\omega+\omega_B}(d\Gamma) = \kappa_\omega(dL_{\omega_B}\Gamma) \quad \text{where } L_{\omega_B}\Gamma(\dots) = \frac{\Gamma(\cdot \exp -\beta H_{B,\omega_B}(\sigma))}{\Gamma(\exp -\beta H_{B,\omega_B}(\sigma))}$$

2. **Translation covariance**: for any translation  $T$ ,  $\kappa_{T\omega}(d\Gamma) = \kappa_\omega(dT\Gamma)$ .

The interface free energy  $F_\Lambda(\omega, \Gamma, \Gamma')$  is now a r.v. over

$$M = d\mathbb{P}(\omega) \kappa_\omega(d\Gamma) \times \kappa'_\omega(d\Gamma')$$

This measure is **translation-invariant**.

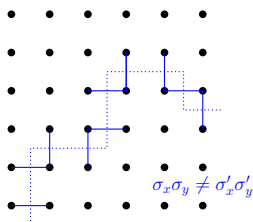
## Incongruent states

Some Gibbs states are more physically relevant.

$\Gamma, \Gamma' \in \mathcal{G}_\omega(\beta)$  are **incongruent** if

$$\liminf_{\Lambda \rightarrow \mathbb{Z}^d} \frac{1}{|\Lambda|} \sum_{(x,y) \in E(\Lambda)} \mathbf{1}_{\{|\Gamma(\sigma_x \sigma_y) - \Gamma'(\sigma_x \sigma_y)| > \delta\}} > 0$$

Positive density of edges with different edge-correlation function.



### Conjecture (EA model)

- ▶  $d = 2$ : At all  $\beta > 0$ , there is no incongruent Gibbs states.  
In fact,  $\#\mathcal{G}_\omega(\beta) = 1$  a.s.
- ▶  $d > 2$ : Are there incongruent Gibbs states ????

Our result is to prove a variance LOWER bound for the interface energy in  $\mathbb{Z}^d$  under an assumption that incongruent states exist.



## Main Result

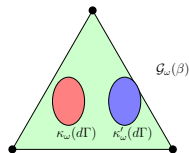
Consider  $\kappa_\omega$  and  $\kappa'_\omega$  two metastates on  $\mathcal{G}_\omega$  and

$M = d\mathbb{P}(\omega) \kappa_\omega(d\Gamma) \times \kappa'_\omega(d\Gamma')$       Probability measure on the triplet  $(\omega, \Gamma, \Gamma')$ .

Assumption (Sufficient for existence of incongruent states)

*There exists an edge  $(x, y)$  such that with positive  $\mathbb{P}$ -probability*

$$\kappa_\omega(\Gamma(\sigma_x \sigma_y)) \neq \kappa'_\omega(\Gamma'(\sigma_x \sigma_y)) .$$



Theorem (A-Newman-Stein-Wehr '14)

*Under the above assumption, there exists  $C > 0$  such that*

$$\text{Var}_M(F_\Lambda(\omega, \Gamma, \Gamma')) \geq C|\Lambda|$$

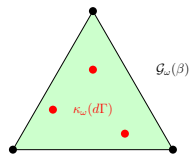
Lower bound for the variance of the interface free energies in  $\mathbb{Z}^d$

## Corollary in $d = 2$

### Corollary (A-Newman-Stein-Wehr '14)

In  $d = 2$ , a metastate  $\kappa_\omega$  *cannot* be supported on a countable infinite number of incongruent states

$$\kappa_\omega(d\Gamma) = \sum_{i=1}^{\infty} p_\omega^i \delta_{\Gamma_\omega^i} .$$



Of course, still far from  $\#\mathcal{G}_\omega(\beta) = 1$ .

Proof.

- ▶ Use  $p_\omega^i$  as tags (like + and - in RFIM):  $\{p_\omega^i\}$  is translation-invariant.
- ▶  $\kappa_\omega = \delta_{\Gamma_\omega^i}$  and  $\kappa'_\omega = \delta_{\Gamma_\omega^{i'}}$  where  $p_\omega^i \neq p_\omega^{i'}$  are **metastates**.
- ▶ On one hand:  $\text{Var}_M(F_\Lambda(J, \Gamma, \Gamma')) \geq C|\Lambda|$ .  
 $\liminf_{\Lambda \uparrow \mathbb{Z}^d} \mathbb{E}[\exp tF_\Lambda / |\Lambda|^{1/2}] \geq e^{Ct^2}$  Martingale CLT
- ▶ On the other hand: upper bound

$$\frac{|F_\Lambda(\omega, \Gamma, \Gamma')|}{|\Lambda|^{1/2}} \leq C \frac{1}{|\partial\Lambda|} \sum_{(x,y) \in E(\partial\Lambda)} |J_{xy}| \leq C'$$



## Possible directions

The interest of the lower bound is to rule out more assumptions

1.  $d = 2$  Prove that there is no incongruent states in  $\mathbb{Z}^2$ .

$$\text{Var}_M(F_\Lambda) = \underbrace{M\left(\text{Var}_M(F_\Lambda|\omega_\Lambda)\right)}_{\text{Fluct. of b.c.}} + \underbrace{\text{Var}_M\left(M(F_\Lambda|\omega_\Lambda)\right)}_{\text{Fluct. of couplings in } \Lambda}$$

2.  $d > 2$  Find a variance UPPER bound to get a contradiction

$$\boxed{\text{Var}\left(\log\frac{Z_{\Lambda,\omega}^{per.}(\beta)}{Z_{\Lambda,\omega}^{anti}(\beta)}\right) \leq C|\partial\Lambda|}$$

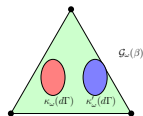
(Aizenman & Fisher, Newman & Stein, Contucci & Giardinà)

Other than gauge-related ?

## Picture of the proof

We want to show  $\text{Var}_M(F_\Lambda(\omega, \Gamma, \Gamma')) \geq C|\Lambda|$  for  $M = d\mathbb{P}(\omega) \kappa_\omega(d\Gamma) \times \kappa'_\omega(d\Gamma')$  under

$$\kappa_\omega(\Gamma(\sigma_x \sigma_y)) \neq \kappa'_\omega(\Gamma'(\sigma_x \sigma_y)) \quad \text{w.p.p.}$$



We have

$$\text{Var}_M(F_\Lambda) \geq \underbrace{\text{Var}_M(M(F_\Lambda | \omega_\Lambda))}_{\text{Fluct. of couplings in } \Lambda}$$

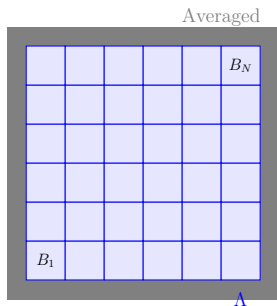
Divide  $\Lambda$  into equally sized blocks

$B_1, \dots, B_k, \dots, B_N$  where  $N = c|\Lambda|$ .

$$\text{Var}_M(F_\Lambda) \geq \sum_{k=1}^N \text{Var}_M(M(F_\Lambda | \omega_{B_k}))$$

To show

- $\text{Var}_M(M(F_\Lambda | \omega_{B_k})) = \text{Var}_M(M(F_\Lambda | \omega_{B_1}))$
- $\text{Var}_M(M(F_\Lambda | \omega_{B_1})) > c'$  for  $c' > 0$  independent of  $\Lambda$ .



## Picture of the proof

$$\begin{aligned}\text{Var } M(F_\Lambda|\omega_B) &= \frac{1}{2} \int d\mathbb{P}(\omega_B) \int d\mathbb{P}(\omega'_B) \left\{ M(F_\Lambda|\omega_B) - M(F_\Lambda|\omega'_B) \right\}^2 \\ &= \frac{1}{2} \int d\mathbb{P}(\omega_B) \int d\mathbb{P}(\omega'_B) \left\{ \int_{\omega'_B \rightarrow \omega_B} \nabla_B M(F_\Lambda|z_B) \cdot dz_B \right\}^2\end{aligned}$$

### Lemma

For any  $(x, y) \in E(B)$

$$\frac{\partial}{\partial \omega_{xy}} M(F_\Lambda|\omega_B) = \beta M(\Gamma(\sigma_x \sigma_y) - \Gamma'(\sigma_x \sigma_y)|\omega_B) \text{ a.s.}$$

*Do not depend on  $\Lambda$  AND Translation invariant*

Let  $\omega_{B^c}$  be the couplings  $\omega$  where the couplings in  $B$  are set to 0:

$$\kappa_\omega(d\Gamma) = \kappa_{\omega_{B^c}}(dL_{\omega_B} \Gamma) \text{ Coupling covariance}$$

$$L_{\omega_B} \Gamma \left( \exp \beta H_{\Lambda, \omega}(\sigma) \right) = \Gamma \left( \exp \beta H_{\Lambda, \omega}(\sigma) \exp -\beta H_{B, \omega}(\sigma) \right) / \Gamma \left( \exp -\beta H_{\omega, B}(\sigma) \right)$$

## Picture of the proof

$$\begin{aligned} & \text{Var } M(F_\Lambda | \omega_B) \\ &= \frac{\beta^2}{2} \int d\mathbb{P}(\omega_B) d\mathbb{P}(\omega'_B) \left\{ \int_{\omega'_B \rightarrow \omega_B} \underbrace{\sum_{(x,y) \in E(B)} M(\Gamma(\sigma_x \sigma_y) - \Gamma'(\sigma_x \sigma_y) | z_B)}_{\nabla_B M(F_\Lambda | z_B) \cdot dz_B} dz_{xy} \right\}^2 \end{aligned}$$

The assumption implies that for  $B$  large enough

$$M\left(\Gamma(\sigma_x \sigma_y) - \Gamma'(\sigma_x \sigma_y) | \omega_B\right) \neq 0 \text{ w.p.p.}$$

since as  $B \rightarrow \mathbb{Z}^d$

$$M\left(\Gamma(\sigma_x \sigma_y) - \Gamma'(\sigma_x \sigma_y) | \omega_B\right) \rightarrow \kappa_\omega\left(\Gamma(\sigma_x \sigma_y)\right) - \kappa'_\omega\left(\Gamma(\sigma_x \sigma_y)\right)$$

## Possible directions

The interest of the lower bound is to rule out more assumptions

1.  $d = 2$  Prove that there is no incongruent Gibbs states in  $\mathbb{Z}^2$ .

$$\text{Var}_M(F_\Lambda) = \underbrace{M\left(\text{Var}_M(F_\Lambda|\omega_\Lambda)\right)}_{\text{Fluct. of b.c.}} + \underbrace{\text{Var}_M\left(M(F_\Lambda|\omega_\Lambda)\right)}_{\text{Fluct. of couplings in } \Lambda}$$

2.  $d > 2$  Find a variance UPPER bound to get a contradiction

$$\text{Var} \left( \log \frac{Z_{\Lambda,\omega}^{per.}(\beta)}{Z_{\Lambda,\omega}^{anti}(\beta)} \right) \leq C|\partial\Lambda|$$

(Aizenman & Fisher, Newman & Stein, Contucci & Giardinà)

Other than gauge-related ?

Thank you!