

# Vector Valued Modular Forms in Vertex Operator Algebras

Jean Auger  
University of Alberta

*Alberta Number Theory Days VIII, BIRS  
Banff, April 2016*

Vertex Operator Algebra = VOA

Origins in deep physics theories that aim beyond QM + GR

Philosophy : The relevance of a VOA is found in its rep theory.

# Overview : modular objects

In the VOA theory...

' $C_2$ -cofiniteness'      vs      'finite # of simple modules'  
vs      modularity of characters

Following work by Y.Zhu, M.Miyamoto proved that the linear span of trace & 'pseudo-trace' functions of such VOAs is **a representation of the modular group**.

# Overview : broad aim

An obstacle to non-ss settings : the lack of examples...

To this date, a single family of VOAs with

- $C_2$ -cofiniteness
- non-semisimple rep theory

has been known... : the  $W(p)$ -triplet VOAs.

# Overview : broad aim

An obstacle to non-ss settings : the lack of examples...

To this date, a single family of VOAs with

- $C_2$ -cofiniteness
- non-semisimple rep theory

has been known... : the  $W(p)$ -triplet VOAs.

## Broad aim

To find new examples of VOAs that are as such.

# Overview : local aim

Several people have been looking for candidate VOAs including D.Adamović, T.Creutzig, A.Milas, D.Ridout, S.Wood.

Some of the more accessible candidates with

- $C_2$ -cofiniteness
- non-ss rep theory

are constructed **out of affine VOAs**.

# Overview : local aim

Several people have been looking for candidate VOAs including D.Adamović, T.Creutzig, A.Milas, D.Ridout, S.Wood.

Some of the more accessible candidates with

- $C_2$ -cofiniteness
- non-ss rep theory

are constructed **out of affine VOAs**.

## Local aim

To expose the character modular invariance property for the most accessible candidate !!

# Overview : a candidate

The VOA  $\mathcal{D}_k$  from the following diagram :

$$L_k(\mathfrak{sl}_2) \xrightarrow{\text{Coset}} \mathcal{C}_k = \text{Com}(\mathcal{H}, L_k(\mathfrak{sl}_2)) \xrightarrow{\text{Extension}} \mathcal{D}_k$$

where

- $k < 0$  &  $k + 2 = \frac{u}{v} \in \mathbb{Q}_{>0} \setminus \{1, \frac{1}{2}, \frac{1}{3}, \dots\}$
- $\mathcal{H}$  = the Heisenberg subalgebra of  $L_k(\mathfrak{sl}_2)$



# Overview : a candidate

The VOA  $\mathcal{D}_k$  from the following diagram :

$$L_k(\mathfrak{sl}_2) \xrightarrow{\text{Coset}} \mathcal{C}_k = \text{Com}(\mathcal{H}, L_k(\mathfrak{sl}_2)) \xrightarrow{\text{Extension}} \mathcal{D}_k$$

where

- $k < 0$  &  $k + 2 = \frac{u}{v} \in \mathbb{Q}_{>0} \setminus \{1, \frac{1}{2}, \frac{1}{3}, \dots\}$
- $\mathcal{H}$  = the Heisenberg subalgebra of  $L_k(\mathfrak{sl}_2)$

Then under a suitable assumption on  $\mathcal{C}_k$ ...

'Schur-Weyl' + Extension process  $\Rightarrow \mathcal{D}_k$  is promising

# Overview : 'Schur-Weyl' duality

Assuming that the vertex tensor theory of HLZ applies for  $\mathcal{C}_k$ ...

**THEOREM** [T.Creutzig, S.Kanade, A.R.Linshaw, D.Ridout]

Then for any a simple  $L_k(\mathfrak{sl}_2)$ -module  $M$  on which  $\mathcal{H}$  acts semisimply, we have a decomposition :

$$M = \bigoplus_{y \in v^M + \text{lattice}} F_y \otimes C_y^M$$

as a  $(\mathcal{H} \otimes \mathcal{C}_k)$ -module where the  $F_y$ 's are Fock spaces and the  $C_y^M$  are simple  $\mathcal{C}_k$ -modules.

+ a few technical properties.

Note :  $\mathcal{H} = \text{Com}(\mathcal{C}_k, L_k(\mathfrak{sl}_2))$ .

One defines characters as :  $\text{tr}_M(y^k z^{h_0} q^{L_0 - \frac{c}{24}})$ .

We should think :  $q = e^{2\pi i\tau}$ .

By some classification work, it is sufficient to consider characters of two types of  $L_k(\mathfrak{sl}_2)$ -modules...

$$\sigma^\ell \mathcal{E}_{\lambda, \Delta_{r,s}}$$

$$\sigma^\ell \mathcal{L}_{r,0}$$

where...

One defines characters as :  $\text{tr}_M(y^k z^{h_0} q^{L_0 - \frac{c}{24}})$ .

We should think :  $q = e^{2\pi i\tau}$ .

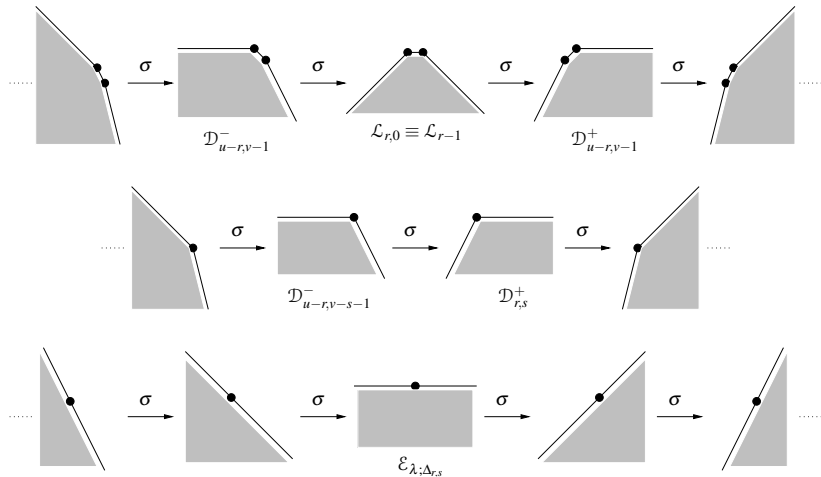
By some classification work, it is sufficient to consider characters of two types of  $L_k(\mathfrak{sl}_2)$ -modules...

$$\sigma^\ell \mathcal{E}_{\lambda, \Delta_{r,s}}$$

$$\sigma^\ell \mathcal{L}_{r,0}$$

where...

- $\ell \in \mathbb{Z}$  &  $\sigma$  is an automorphism of  $L_k(\mathfrak{sl}_2)$
- $r \in \{1, \dots, u-1\}$  &  $s \in \{0, \dots, v-1\}$
- $\lambda \in \frac{1}{v}\mathbb{Z}$



Source : T.Creutzig, D.Ridout, *Modular Data and Verlinde Formulae for Fractional Level WZW Models II*, Nucl. Phys. B 875 (2013) 423. I thank the authors for allowing me to use this picture.

Decomposing the relevant characters accordingly to the 'Schur-Weyl' result, we get :

$$\text{ch } \sigma^\ell \mathcal{E}_{\lambda, \Delta_{r,s}} = \sum_{n \in \mathbb{Z}} (\text{ch } F_{\lambda+2n+k\ell}) \cdot (\text{ch } C_{r,s, \lambda+2n}^{\mathcal{E}}(q))$$

$$\text{ch } \sigma^\ell \mathcal{L}_{r,0} = \sum_{n \in \mathbb{Z}} (\text{ch } F_{r-1+2n+k\ell}) \cdot (\text{ch } C_{r, r-1+2n}^{\mathcal{L}}(q))$$

where...

Decomposing the relevant characters accordingly to the 'Schur-Weyl' result, we get :

$$\text{ch } \sigma^\ell \mathcal{E}_{\lambda, \Delta_{r,s}} = \sum_{n \in \mathbb{Z}} (\text{ch } F_{\lambda+2n+k\ell}) \cdot (\text{ch } C_{r,s,\lambda+2n}^{\mathcal{E}}(q))$$

$$\text{ch } \sigma^\ell \mathcal{L}_{r,0} = \sum_{n \in \mathbb{Z}} (\text{ch } F_{r-1+2n+k\ell}) \cdot (\text{ch } C_{r,r-1+2n}^{\mathcal{L}}(q))$$

where...

$$\text{ch } C_{r,s,x}^{\mathcal{E}}(q) = \frac{\chi_{r,s}^{\text{Vir}}(q)}{\eta(q)} q^{-\frac{1}{4k}x^2}$$

$$\begin{aligned} \text{ch } C_{r,x}^{\mathcal{L}}(q) = & \sum_{d=1}^{v-1} (-1)^{d-1} \frac{\chi_{r,d}^{\text{Vir}}(q)}{\eta(q)} \cdot \sum_{a=0}^{\infty} q^{-\frac{1}{4k}(x-k(2av+d))^2} \\ & - q^{-\frac{1}{4k}(x-k(2(a+1)v-d))^2} \end{aligned}$$

Set  $p = -kv^2$  and  $\Gamma = \sqrt{2p}\mathbb{Z}$ .

Lifting the  $\mathcal{C}_k$ -modules  $\mathcal{C}_{r,s,x}^{\mathcal{E}}$  and  $\mathcal{C}_{r,x}^{\mathcal{L}}$  results in the apparition of lattice  $\Theta$ -functions and derivatives :

$$\underbrace{D_{r,s,\omega}^{\mathcal{E},0}(q)}_{\Theta} + \underbrace{0}_{\Theta'}$$

$$\underbrace{D_{r,t}^{\mathcal{L},0}(q)}_{\Theta} + \underbrace{D_{r,t}^{\mathcal{L},1}(q)}_{\Theta'}$$

where...



Set  $p = -kv^2$  and  $\Gamma = \sqrt{2p}\mathbb{Z}$ .

Lifting the  $\mathcal{C}_k$ -modules  $\mathcal{C}_{r,s,x}^{\mathcal{E}}$  and  $\mathcal{C}_{r,x}^{\mathcal{L}}$  results in the apparition of lattice  $\Theta$ -functions and derivatives :

$$\underbrace{D_{r,s,\omega}^{\mathcal{E},0}(q)}_{\Theta} + \underbrace{0}_{\Theta'}$$

$$\underbrace{D_{r,t}^{\mathcal{L},0}(q)}_{\Theta} + \underbrace{D_{r,t}^{\mathcal{L},1}(q)}_{\Theta'}$$

where...

$$D_{r,s,\omega}^{\mathcal{E},0}(q) = \frac{\chi_{r,s}^{Vir}(q)}{\eta(q)} \Theta_{\frac{\omega}{\sqrt{2p}} + \Gamma}(1, q)$$

$D_{r,t}^{\mathcal{L},0}(q)$  = a linear combination of expressions of the form  $D_{r,s,\omega}^{\mathcal{E},0}(q)$

$$D_{r,t}^{\mathcal{L},1}(q) = \sum_{d=1}^{v-1} (-1)^{d-1} \frac{\chi_{r,d}^{Vir}(q)}{\eta(q)} \left( \Theta'_{\frac{(r-1+2t)v+kv d}{\sqrt{2p}} + \Gamma}(1, q) - \Theta'_{\frac{(r-1+2t)v-kv d}{\sqrt{2p}} + \Gamma}(1, q) \right)$$

$$\frac{\chi_{r,s}^{Vir}(q)}{\eta(q)} \Theta_{\frac{\omega}{\sqrt{2p}} + \Gamma}(1, q)$$

Consider the generating modular transformations

$$S : \tau \mapsto -\frac{1}{\tau}$$

$$T : \tau \mapsto \tau + 1$$

$\text{Span}_{\mathbb{C}} \{D_{r,s,\omega}^{\mathcal{E},0}(q)\}$  is then **automatically** a representation of  $PSL(2, \mathbb{Z})$  !

$$\sum_{d=1}^{v-1} (-1)^{d-1} \frac{\chi_{r,d}^{Vir}(q)}{\eta(q)} \left( \Theta'_{\frac{(r-1+2t)v+kvd}{\sqrt{2p}}+\Gamma}(1, q) - \Theta'_{\frac{(r-1+2t)v-kvd}{\sqrt{2p}}+\Gamma}(1, q) \right)$$

Fix parameters  $r, t$  and write

$$D_{r,t}^{\mathcal{L},1}\left(-\frac{1}{\tau}\right) = \sum \text{Coeff}_{(r',s'),\omega} \cdot \left( \frac{\chi_{r',s'}^{Vir}(\tau)}{\eta(\tau)} \Theta'_{\frac{\omega}{\sqrt{2p}}+\Gamma}(\tau) \right)$$

$$\sum_{d=1}^{v-1} (-1)^{d-1} \frac{\chi_{r,d}^{Vir}(q)}{\eta(q)} \left( \Theta'_{\frac{(r-1+2t)v+kvd}{\sqrt{2p}}+\Gamma}(1, q) - \Theta'_{\frac{(r-1+2t)v-kvd}{\sqrt{2p}}+\Gamma}(1, q) \right)$$

Fix parameters  $r, t$  and write

$$D_{r,t}^{\mathcal{L},1}\left(-\frac{1}{\tau}\right) = \sum \text{Coeff}_{(r',s'),\omega} \cdot \left( \frac{\chi_{r',s'}^{Vir}(\tau)}{\eta(\tau)} \Theta'_{\frac{\omega}{\sqrt{2p}}+\Gamma}(\tau) \right)$$

Fix  $d$ . Then for any  $r', t'$ , one can find that

$$(-1)^{d-1} \text{Coeff}_{(r',d), (r'-1+2t')v \pm kvd} = \pm [\#(r, t, r', t')]$$

... and that the irrelevant 'Coeffs' vanish !

## RESULT

The vector space

$$V = \text{Span}_{\mathbb{C}} \{ D_{r,s,\omega}^{\mathcal{E},0}(q) + 0, D_{r,t}^{\mathcal{L},0}(q) + D_{r,t}^{\mathcal{L},1}(q) \}$$

is a representation of  $PSL(2, \mathbb{Z})$  !

## RESULT

The vector space

$$V = \text{Span}_{\mathbb{C}} \{ D_{r,s,\omega}^{\mathcal{E},0}(q) + 0, D_{r,t}^{\mathcal{L},0}(q) + D_{r,t}^{\mathcal{L},1}(q) \}$$

is a representation of  $PSL(2, \mathbb{Z})$  !

More interestingly  $\text{Span}_{\mathbb{C}} \{ D_{r,t}^{\mathcal{L},1}(q) \}$  also is ;

$$D_{r,t}^{\mathcal{L},1} \left( -\frac{1}{\tau} \right) = \sum S_{(r,t),(r',t')}^{\mathcal{L},1} \cdot D_{r',t'}^{\mathcal{L},1}(\tau)$$

where

$$S_{(r,t),(r',t')}^{\mathcal{L},1} = \underbrace{X_{(r',t')}}_{1 \text{ or } 1/2} \cdot \frac{4i\tau}{\sqrt{u}\sqrt{2v-u}} \sin\left(\pi \frac{v}{u} rr'\right) \cos\left(\pi \frac{(r-1+2t)(r'-1+2t')v}{2v-u}\right)$$