

# Duality for Metaplectic Ice

arXiv:1709.06500

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Also:

Brubaker, Bump, Chinta, Friedberg, Gunnells (arXiv:1009.1741)

Brubaker, Buciumas and Bump ([arXiv:1604.02206](#))

Brubaker, Buciumas, Bump and Friedberg (arXiv:1704.00701)

## Partition functions

In statistical mechanics, we consider an ensemble of many states. If  $s$  is a state let  $E(s)$  be its energy.

More energetic states become more probable if the system is at a higher temperature.

In Maxwell-Boltzmann statistics the probability of state  $s$  is proportional to  $e^{-E(s)/k_B T}$  where  $k_B$  is Boltzmann's constant.

Thus the probability of state  $s$  is

$$\frac{1}{Z} e^{-E(s)/k_B T} \quad \text{where} \quad Z = \sum_s e^{-E(s)/k_B T}.$$

This  $Z$  is the **partition function**. Its significance goes beyond this probabilistic fact. In some sense it contains complete thermodynamic information about the system.

## Solvable Lattice Models

- It was found that certain 2-dimensional lattice models can be analyzed completely.
- The first such **solvable** lattice model to be fully analyzed was the **Ising model** (Onsager, 1944).
- An identity that Onsager and Baxter called the “star-triangle relation” was used powerfully by Baxter. It was renamed the **Yang-Baxter equation** by Faddeev.
- For us, the Yang-Baxter equation characterizes solvability of the model.
- Fitting it into an algebraic framework led to the invention of quantum groups (Drinfeld, Jimbo).

## Analytic continuation of the partition function

There is utility in considering nonreal Boltzmann weights. A **spin** is  $\pm 1$ . A state of the **Ising model** consists of an assignment of spins  $s = \{s_j\}$  at vertices  $v_j$  of a lattice.

$$Z = \sum_s e^{-E(s)/k_B T}, \quad -E(s) = \sum_{i,j \text{ adjacent}} J s_i s_j + \sum_i t s_i.$$

Here  $t$  is a parameter corresponding to an external field. It may be fruitfully be regarded as a **complex parameter**.

### Theorem

*(Lee-Yang) The zeros of  $Z$  lie on the line  $\operatorname{re}(t) = 0$ .*

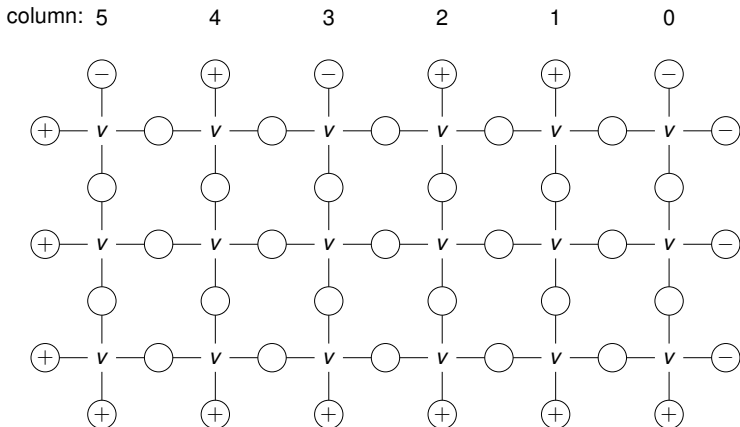
This is a kind of Riemann hypothesis. (There are other similarities between partition functions and L-functions!) We do not require the Boltzmann weights to be real.

## Duality in Lattice Models

- Sometimes a system may be described in very different ways.
- In quantum mechanics, the world looks different in position or momentum coordinates, related by the Fourier transform.
- Kramers and Wannier proved that for the Ising model, the partition functions of high and low temperature models are related, computing the location of the phase transition.
- In Type IIB string theory,  $S$  duality relates the partition functions for inverse coupling constant  $y$  with the system for  $1/y$ . The partition function is an automorphic form.
- We will describe two systems originating in number theory that have the same partition function.

## Six Vertex Model

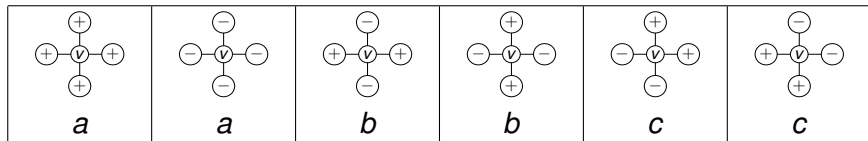
Begin with a grid, usually (but not always) rectangular:



Edges are labeled with “spins”  $\pm$ . **Boundary** spins are fixed. A **state** assigns spins to interior edges.

## Boltzmann weights: 6 Vertex Model

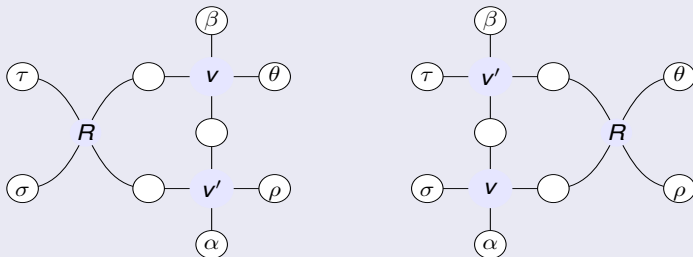
- A **state** of the system assigns spins to the interior edges.
- Each state has a **Boltzmann weight** which is the product of Boltzmann weights for each vertex.
- The Boltzmann weights for a vertex depend on the spins of the adjacent edges.
- The **field-free six vertex model** was analyzed by Lieb, Sutherland, Baxter, Korepin, Izergin. It was a primary example leading to the discovery of quantum groups.
- Let  $a = a(v)$ ,  $b = b(v)$ ,  $c = c(v)$  be complex numbers. They depend on the vertex  $v$ . Let  $\Delta(v) = \frac{a^2 + b^2 - c^2}{2ab}$ .



## Yang-Baxter equation: 6 vertex model

### Theorem

**(Baxter)** If  $\Delta(v) = \Delta(v')$  (condition on Boltzmann weights) there is a vertex  $R$  such that for all boundary spins  $\sigma, \tau, \dots$  the partition functions



are equal. (The interior edge spins are summed.)

Implies: (Baxter) **Field-free 6-vertex model is solvable.**



## Braided Categories

We may associate with every edge  $\epsilon$  a module  $V_\epsilon$  in a monoidal category. This module is a vector space with basis  $+$  and  $-$ .



**Boltzmann weights are arranged as coefficients in the R-matrix!**

is interpreted as an isomorphism  $V_\alpha \otimes V_\beta \rightarrow V_\gamma \otimes V_\delta$ . The Yang-Baxter equation means this diagram commutes:

$$\begin{array}{ccc}
 & V_\tau \otimes V_\sigma \otimes V_\beta \rightarrow V_\tau \otimes V_\beta \otimes V_\sigma & \\
 \nearrow & & \searrow \\
 V_\sigma \otimes V_\tau \otimes V_\beta & & V_\tau \otimes V_\sigma \otimes V_\beta \\
 \searrow & & \nearrow \\
 & V_\sigma \otimes V_\beta \otimes V_\tau \rightarrow V_\beta \otimes V_\sigma \otimes V_\tau &
 \end{array}$$

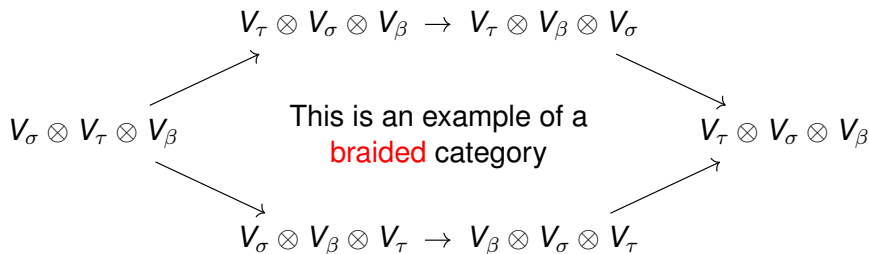
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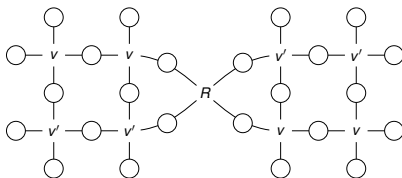


## Quantum Groups

- Given a braided category Tannakian theory or other methods (Drinfeld, FRT, etc.) tries to produce a Hopf algebra whose (co)modules are the given category.
- For the field-free Yang-Baxter equation, we get  $U_q(\widehat{\mathfrak{sl}}_2)$  where

$$\frac{1}{2}(q - q^{-1}) = \Delta = \frac{a^2 + b^2 - c^2}{2ab}.$$

- This is the historical route to quantum groups.



## Metaplectic Whittaker functions as partition functions

If  $F$  is a  $p$ -adic field containing the  $n$ -th roots of unity  $\mu_n$ , there is a central extension

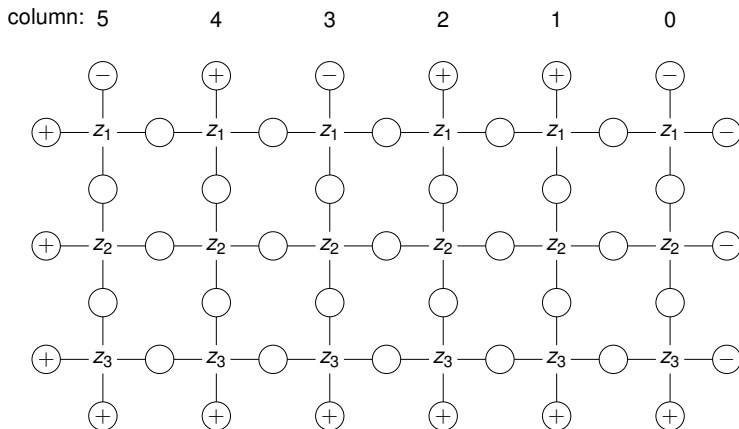
$$1 \longrightarrow \mu_n \longrightarrow \widetilde{\mathrm{GL}}(r) \longrightarrow \mathrm{GL}(r) \longrightarrow 1.$$

The group  $\widetilde{\mathrm{GL}}(r)$  is the **metaplectic group**.

- Interesting automorphic forms live on metaplectic groups.
- Brubaker, Bump, Chinta, Friedberg, Gunnells proposed ...
- ... and Brubaker, Buciumas and Bump, showed:
- **Whittaker functions on  $\widetilde{\mathrm{GL}}(r)$  may be represented as partition functions of solvable lattice models.**
- The relevant quantum group is  $U_{\sqrt{q^{-1}}}(\widehat{\mathfrak{gl}}(n|1))$ .
- This setup exhibits an interesting duality.

## Metaplectic Ice

Proceed as before but label vertices by  $z_i \in \mathbb{C}^\times$ :



## Boltzmann Weights for metaplectic $\Gamma$ Ice







Choose parameters  $g(a)$  depending on  $a \bmod n$ :

$$g(a)g(-a) = 1/v \text{ if } n \nmid a,$$

$$g(a) = -v \text{ if } n|a.$$

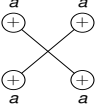
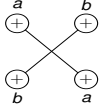
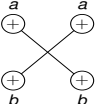
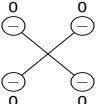
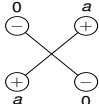
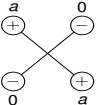
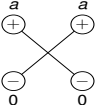
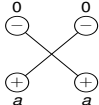
(Secretly Gauss sums).

Let  $z_i \in \mathbb{C}^\times$ . Here are the Boltzmann weights:

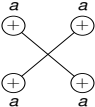
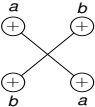
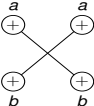
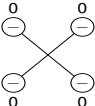
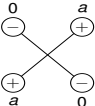
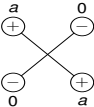
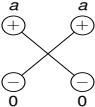
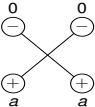
					
$1$	$z_i$	$g(a)$	$z_i$	$(1 - v)z_i$	$1$

- Horizontal edges have  $n + 1$  possible states:  
 $+a \ (a \bmod n)$  or  $-0$ .
- The module corresponding to a horizontal edge is a  $(n|1)$ -dimensional supervector space.

The  $\Gamma\Gamma$  R-matrix

<p>Another type of vertex, the <b>R-matrix</b> <math>R_{z_1, z_2}^{\Gamma\Gamma}</math>.</p>	 $z_2^n - v z_1^n$	 $g(a-b)(z_1^n - z_2^n)$
 $(1-v)z_1^{a-b}z_2^{n-a+b}$	 $z_1^n - v z_2^n$	 $v(z_1^n - z_2^n)$
 $z_1^n - z_2^n$	 $(1-v)z_1^a z_2^{n-a}$	 $(1-v)z_1^{n-a} z_2^a$

The  $\Gamma\Gamma$  R-matrix

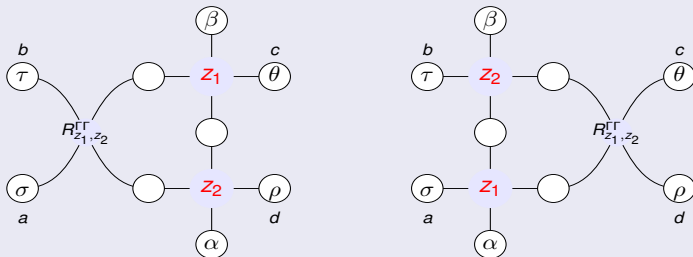
<p>Another type of vertex, the R-matrix <math>R_{Z_1, Z_2}^{\Gamma\Gamma}</math>. Drinfeld twist of <math>\widehat{\mathfrak{gl}}(n 1)</math></p>	 $z_2^n - v z_1^n$	 $g(a-b)(z_1^n - z_2^n)$
 $(1-v)z_1^{a-b}z_2^{n-a+b}$	 $z_1^n - v z_2^n$	 $v(z_1^n - z_2^n)$
 $z_1^n - z_2^n$	 $(1-v)z_1^a z_2^{n-a}$	 $(1-v)z_1^{n-a} z_2^a$



## Gamma-Gamma Yang-Baxter Equation

### Theorem

The following partition functions are equal.

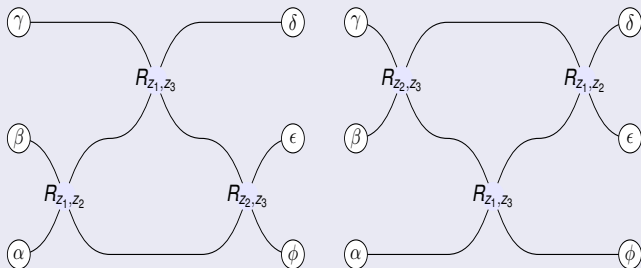


(Here  $a, b, c, d$  are the charges.)

## Another Yang-Baxter equation

### Theorem

Fix  $z_1, z_2$  and  $z_3$  and (decorated) boundary spins. The following partition functions are equal:

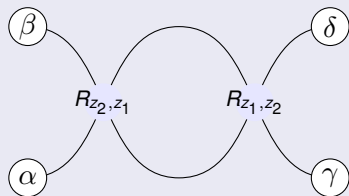


Brubaker, Buciumas, Bump and Friedberg use this to construct representations of the affine Hecke algebra.

## Complementary Equation

## Theorem

Let  $\alpha, \beta, \gamma, \delta$  be decorated spins. Then the partition function



$$\text{equals } \begin{cases} (z_1^n - vz_2^n)(z_2^n - vz_1^n) & \text{if } \alpha = \gamma, \beta = \delta \\ 0 & \text{otherwise.} \end{cases}$$

The horizontal edges correspond to evaluation modules for a Drinfeld twist of  $U_{\sqrt{v}}(\widehat{\mathfrak{gl}}(n|1))$

## Whittaker coinvariants as $U_{\sqrt{q^{-1}}}(\widehat{\mathfrak{gl}}(n))$ -modules

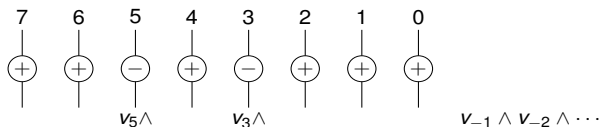
- $V$ : A representation of  $GL(r, F)$ ,  $V_{u, \psi}$  its module of **Whittaker coinvariants**, the largest quotient through which every Whittaker functional factors.
- $\mathbf{z} = (z_1, \dots, z_r) \in \mathbb{C}^r \cong \widehat{T}(\mathbb{C})$ , the Langlands dual torus. It parametrizes a principal series representation  $V_{\mathbf{z}}$ .
- $\mathcal{A}_W : V_{\mathbf{z}} \rightarrow V_{W\mathbf{z}}$  intertwining integral.

The scattering matrix of  $\mathcal{A}_{S_i}$  on  $(V_{\mathbf{z}})_{U, \psi}$  was computed by Kazhdan and Patterson. It agrees with the R-matrix of (Drinfeld twisted)  $U_{\sqrt{q^{-1}}}(\widehat{\mathfrak{gl}}(n))$ !

$$\begin{array}{ccc}
 (V_{\mathbf{z}})_{U, \psi} & \xrightarrow{\cong} & V_{z_1} \otimes \cdots \otimes V_{z_i} \otimes V_{z_{i+1}} \otimes \cdots \otimes V_{z_r} \\
 \downarrow \mathcal{A}_{S_i} & & \downarrow \text{R-matrix} \\
 (V_{S_i \mathbf{z}})_{U, \psi} & \xrightarrow{\cong} & V_{z_1} \otimes \cdots \otimes V_{z_{i+1}} \otimes V_{z_i} \otimes \cdots \otimes V_{z_r}
 \end{array}$$

## Fermionic interpretation of vertical edges

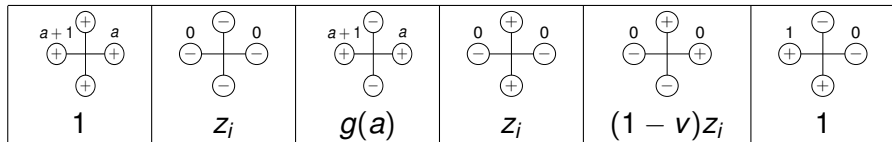
- The vertical edges which take values  $\pm$  have no interpretation as for  $U_{\sqrt{v}}(\widehat{\mathfrak{gl}}(n))$  modules, but the aggregation of them may be interpreted as an element of the **fermionic Fock space**.



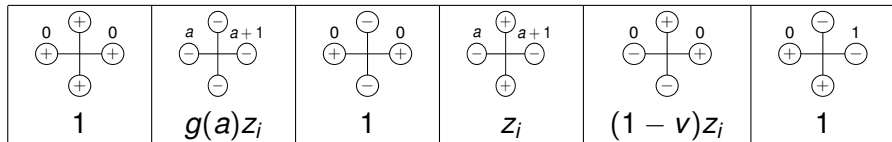
- Kashiwara, Miwa and Stern defined a  $U_q(\widehat{\mathfrak{gl}}(n))$  module structure on the degree  $k$  Fermionic Fock space  $\mathfrak{F}_{(k)}$  and a **vertex operator**  $V_z \otimes \mathfrak{F}_{k-1} \longrightarrow \mathfrak{F}_k$ . We think the row transfer matrix for Gamma ice computes this.

## Gamma ice versus Delta Ice

Gamma Ice:



Delta Ice:



Same boundary conditions for Delta ice but reverse the rows.

**Theorem (Duality Theorem)**

*The partition functions are equal!*

## Why this was discovered

- An equivalent statement was proved combinatorially by Brubaker, Bump and Friedberg (2011).
- The proof was difficult.
- Motivation: this is the  $p$ -part of a multiple Dirichlet series. This combinatorial fact implies the analytic continuation and functional equation.
- Brubaker, Bump, Chinta, Friedberg and Gunnells proposed that this could be proved using the Yang-Baxter equation (2012).
- A proof using the Yang-Baxter equation was given by Brubaker, Buciumas, Bump and Gray (2017)

### Theorem (**Duality Theorem**)

*The partition functions are equal!*