

On some relations between Optimal Transport and Stochastic Geometric Mechanics

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Based on joint works with M. Arnaudon, X. Chen, J.-C. Zambrini

Extensions of (some) results, with applications to fluids: with T. Ratiu,
C. Léonard

Recalling the Monge-Kantorovich (MK) optimal transport problem (flat case)

$$\inf_{\pi \in \Pi(\mu, \sigma)} \frac{1}{2} \int_{\mathbb{R}^d \times \mathbb{R}^d} \|x - y\|^2 d\pi(x, y)$$

where $\Pi(\mu, \sigma) = \{ \text{joint distributions s.t. the marginals along } x \text{ and } y \text{ coordinates are } \mu \text{ and } \sigma \text{ resp. } \}$.

Since “the cost”

$$\|x - y\|^2 = \inf_X \int_0^1 \|\dot{X}(t)\|^2 dt$$

$X \in \{C([0, 1]; \mathbb{R}^d) : X(0) = x, X(1) = y\}$, the MK problem is equivalent to

$$\inf_{\pi} \frac{1}{2} \int_{\mathbb{R}^d \times \mathbb{R}^d} \left(\int \int_0^1 \|\dot{X}(t)\|^2 dt dP_{x,y} \right) d\pi(x, y)$$

or, by desintegration, to

$$\inf_P \int \int_0^1 \frac{1}{2} \|\dot{X}(t)\|^2 dt dP$$

with P prob. measure on $C^1([0, 1]; \mathbb{R}^d)$ with marginals μ and σ at $t = 0, t = 1$ and $P_{x,y}$ the one with initial and final marginals δ_x and δ_y .

Eulerian (control) version of this problem:

$$\inf_v \frac{1}{2} \int \int_0^1 \|v(t, x^v(t))\|^2 dt dP$$

v continuous, $\dot{x}^v(t) = v(t, x^v(t))$ a.s., law of $x^v(0) = \mu$, law of $x^v(1) = \sigma$.

If $d\mu = \rho_0 dx$, $d\sigma = \rho_1 dx$ given, law of $x^v(t) = \rho_t dx$, then

$$\frac{\partial \rho}{\partial t} = -\operatorname{div}(\rho v)$$

(continuity equation)

If ψ solves the Hamilton-Jacobi equation

$$\frac{\partial \psi}{\partial t} = -\frac{1}{2} \|\nabla \psi\|^2$$

and $\frac{\partial \rho}{\partial t} = -\operatorname{div}(\rho \nabla \psi)$, with $\rho(0, x) = \rho_0(x)$, $\rho(1, x) = \rho_1(x)$,

then $v = \nabla \psi$ solves the problem. Moreover

$$\frac{\partial v}{\partial t} = -(v \cdot \nabla)v$$

On $C([0, 1]; \mathbb{R}^d)$ consider the laws (Q) of diffusions

$$dX(t) = \sqrt{\epsilon} dW(t) + Y(t)dt, \quad \text{law of } X(0) = dx$$

Entropy functional of Q w.r.t. P :

$$H(Q; P) = \int \log\left(\frac{dQ}{dP}\right) dQ$$

Choose for P the law of the Wiener process

$$H(Q; P) = H(Q_0; P_0) + \frac{1}{2} \int \int_0^1 \|Y(t)\|^2 dt dQ,$$

where Q_0, P_0 are the marginals of Q and P at time 0. This is a consequence of Girsanov's Theorem.

Schrödinger problem: minimise the entropy functional, subject to given Q_0 and Q_1 .

If $Y(t) = v(t, X(t))$, the density ρ_t of X at time t satisfies the Fokker-Planck equation

$$\frac{\partial \rho}{\partial t} = -\operatorname{div}(\rho v) + \frac{\epsilon}{2} \Delta \rho$$

If, moreover, $v = \nabla \varphi$ satisfies Hamilton-Jacobi-Bellman equation,

$$\frac{\partial \varphi}{\partial t} = -\frac{1}{2} \|\nabla \varphi\|^2 + \frac{\epsilon}{2} \Delta \varphi$$

then $v = \nabla \varphi$. So

$$\frac{\partial v}{\partial t} = -(v \cdot \nabla) v + \frac{\epsilon}{2} \Delta v$$

(Burgers equation)

Remarks:

1. With the change of variable $v = \nabla\varphi$, $\varphi = -\log \eta$, $\frac{\partial\eta}{\partial t} = \frac{\epsilon}{2}\Delta\eta$ (heat equation)
2. When $\epsilon \rightarrow 0$ formally Schrödinger problem converges to the optimal transport problem (C. Léonard, etc)
(delicate problem)

G Lie group with

$\langle \cdot, \cdot \rangle$ right (left) invariant metric

∇ Levi-Civita connection

e identity element

$\mathcal{G} \simeq T_e(G)$ Lie algebra

$\{H_i\}$ o.n. basis of \mathcal{G} , right invariant, (we suppose finite dimensional),

$\nabla_{H_i} H_i = 0$

Brownian motion on G with generator = Laplace-Beltrami operator:

$$dg^0(t) = T_e R_{g^0(t)} \left(\sum_i H_i \circ dW^i(t) \right) = T_e R_{g^0(t)} \left(\sum_i H_i \cdot dW^i(t) \right)$$

where $T_a R_{g(t)} : T_a G \rightarrow T_{ag(t)} G$ is the differential of the right translation $R_{g(t)}(x) := xg(t)$, $\forall x \in G$ at the point $x = a \in G$.

General diffusion processes:

$$dg(t) = T_e R_{g(t)} \left(\sum_i H_i \circ dW^i(t) + u(t)dt \right)$$

$$T_e R_{g(t)} u(t) = \frac{D^\nabla g(t)}{dt}$$

where, by definition, D^∇ is the (mean) derivative defined as: for $\xi(t) = \int_0^t T_{0 \leftarrow s} \circ dg(s)$, where $T_{\cdot \leftarrow 0} : T_{g(0)}G \rightarrow T_{g(\cdot)}G$ is the stochastic parallel transport associated to ∇ , define

$$\frac{D\xi(t)}{dt} = \lim_{\epsilon \rightarrow 0} \frac{1}{\epsilon} E \left[\xi(t + \epsilon) - \xi(t) | \mathcal{P}_t \right]$$

and

$$\frac{D^\nabla g(t)}{dt} := T_{t \leftarrow 0} \frac{D\xi(t)}{dt}$$

From Girsanov's Theorem the law Q of g on $C([0, 1]; G)$ is absolutely continuous w.r.t. the law P of g^0 (same initial distributions) with density given by

$$\frac{dQ}{dP} = \exp\left\{ \int_0^1 \sum_i \left\langle T_g R_{g^{-1}} \frac{D^\nabla g(t)}{dt}, H_i dW^i(t) \right\rangle - \frac{1}{2} \int_0^1 \left\| T_g R_{g^{-1}} \frac{D^\nabla g(t)}{dt} \right\|^2 dt \right\}$$

Therefore the entropy is

$$H(Q; P) = \frac{1}{2} \int \int_0^1 \left\| T_g R_{g^{-1}} \frac{D^\nabla g(t)}{dt} \right\|^2 dt dQ$$

Denote by α the invariant measure on G , P_x law of the Brownian motion g^0 on G starting at x , $P_t(x, \alpha(dy))$ its transition semigroup.

$$\tilde{P}_\mu = \int_G P_x \mu(dx)$$

Assume μ and σ prob. measures, to be abs. cont. w.r.t. α .

$\mathcal{M}(\mu, \sigma) := \{\pi \text{ prob. measures on } G \times G : \pi(g^0(0) \in \cdot) = \mu(\cdot), \pi(g^0(1) \in \cdot) = \sigma(\cdot), \pi \text{ abs. cont. w.r.t. } \mu \otimes P_1\}$

For $\pi \in \mathcal{M}(\mu, \sigma)$, define

$$P_\pi(d\omega) = \int_{G \times G} \pi(dx, dy) P(d\omega | 0, x; 1, y)$$

prob. measure on $C([0, 1]; G)$.

By Csiszár's Theorem, if $\exists \pi_0 \in \mathcal{M}(\mu, \sigma) : H(\pi_0, \mu \otimes P_1) < \infty$,
 then $\exists^1 Q$ on $C([0, 1]; G)$ attaining the

$$\inf_{Q: Q((g(0), g(1)) \in \cdot) \in \mathcal{M}(\mu, \sigma)} H(Q; \tilde{P}_\mu)$$

with $Q = P_\pi$, π attaining

$$\inf_{\pi} H(\pi; \mu \otimes P_1)$$

Moreover Q is the law of a Markov process.

Combining with Girsanov's result, we have

$$\begin{aligned} H(Q; \tilde{P}_\mu) &= \frac{1}{2} \int \int_0^1 \left\| T_g R_{g^{-1}} \frac{D^\nabla g(t)}{dt} \right\|^2 dt dQ \\ &:= A[g] \end{aligned}$$

Stochastic Euler-Poincaré reduction theorem

Theorem. The G -valued semi-martingale

$$dg(t) = T_e R_{g(t)} \left(\sum_i H_i \circ dW^i(t) + u(t)dt \right)$$

is a critical point of A if and only if the time-dependent vector field $u(\cdot)$ satisfies the equation,

$$\frac{d}{dt} u(t) = -ad_{u(t)}^* u(t) - K(u(t)),$$

where

for $u \in \mathcal{G}$, $ad_u^* : \mathcal{G} \rightarrow \mathcal{G}$ is the adjoint of $ad_u : \mathcal{G} \rightarrow \mathcal{G}$ with respect to the metric $\langle \cdot, \cdot \rangle$,

and the operator $K : \mathcal{G} \rightarrow \mathcal{G}$ is defined as

$$K(u) = -\frac{1}{2} \sum_i (\nabla_{H_i} \nabla_{H_i} u + R(u, H_i) H_i) = -\frac{1}{2} \square(u)$$

(de Rham-Hodge Laplacian)

(Left) variations:

$$\partial_L \mathbf{A}[\xi(\cdot)] = \left. \frac{d}{d\varepsilon} \right|_{\varepsilon=0} \mathbf{A}[\xi_{\varepsilon, v}(\cdot)]$$

with $\xi_{\varepsilon, v}(t) = \mathbf{e}_{\varepsilon, v}(t)\xi(t)$

$$\begin{cases} \frac{d}{dt} \mathbf{e}_{\varepsilon, v}(t) = \varepsilon T_e R_{\mathbf{e}_{\varepsilon, v}(t)} \dot{v}(t), \\ \mathbf{e}_{\varepsilon, v}(0) = \mathbf{e} \end{cases} \quad (1)$$

for $v(\cdot) \in C^1([0, T]; \mathcal{G})$, $v(0) = v(T) = 0$

We differentiate in the direction of v ($\left. \frac{d}{d\varepsilon} \right|_{\varepsilon=0} \mathbf{e}_{\varepsilon, v} = v$)

"Proof" of the Theorem:

For $g_\varepsilon(t) = e_{\varepsilon, v}(t)g(t)$, by Itô's formula ,

$$\begin{aligned}
 dg_\varepsilon(t) &= \sum_i T_e R_{g_\varepsilon(t)} \text{Ad}_{e_\varepsilon^{-1}(t)} H_i \circ dW_t^i \\
 &+ T_e R_{g_\varepsilon(t)} \left(\text{Ad}_{e_\varepsilon^{-1}(t)} (u(t)) + T_{e_\varepsilon(t)} R_{e_\varepsilon^{-1}(t)} \dot{e}_\varepsilon(t) \right) dt,
 \end{aligned} \tag{2}$$

(no contraction term)

We have $T_{e_\varepsilon(t)} R_{e_\varepsilon^{-1}(t)} \dot{e}_\varepsilon(t) = \varepsilon \dot{v}(t)$, $\frac{d}{d\varepsilon} \Big|_{\varepsilon=0} \text{Ad}_{e_\varepsilon^{-1}(t)} u(t) = -\text{ad}_{v(t)} u(t)$
 and $\frac{d}{d\varepsilon} \Big|_{\varepsilon=0} \text{Ad}_{e_\varepsilon^{-1}(t)} H_i = \text{ad}_{v(t)} H_i$

$$\begin{aligned}
\frac{dA[g_\varepsilon(\cdot)]}{d\varepsilon} \Big|_{\varepsilon=0} &= \mathbb{E} \int_0^1 \left\langle \frac{d}{d\varepsilon} (T_{g_\varepsilon(t)} R_{g_\varepsilon^{-1}(t)} \frac{D^\nabla g_\varepsilon(t)}{dt}) \Big|_{\varepsilon=0}, u(t) \right\rangle dt \\
&= \int_0^1 \left\langle u(t), \dot{v}(t) - \text{ad}_{(u(t))} v(t) \right. \\
&\quad \left. + \frac{1}{2} \sum_i (\nabla_{\text{ad}_{v(t)} H_i} H_i + \nabla_{H_i} (\text{ad}_{v(t)} H_i)) \right\rangle dt \\
&= \int_0^1 \left\langle -\dot{u}(t) - \text{ad}_{u(t)}^* u(t) - K(u(t)), v(t) \right\rangle dt
\end{aligned} \tag{3}$$

Remark : When $H_i = 0$ we recover the deterministic Euler-Poincaré reduction theorem.

The change of variables $u(t) = \nabla\varphi(t)$ gives the Hamilton-Jacobi-Bellman equation

$$\frac{\partial\varphi}{\partial t} = -\frac{1}{2}\|\nabla\varphi\|^2 + \frac{1}{2}\Delta_{LB}(\varphi)$$

Extra remarks:

In the assumptions of Csiszär's theorem

($\exists\pi_0 \in \mathcal{M}(\mu, \sigma) : H(\pi_0, \mu \otimes P_1) < \infty$), if π_0 is abs. cont. w.r.t. $\mu \otimes P_1$, then the minimiser is of the form

$$\pi(dx, dy) = \psi(x)\phi(y)\alpha(dx)P_1(x, \alpha(dy))$$

and

$$\frac{d\mu}{d\alpha} = \psi P_1 \phi, \quad \frac{d\nu}{d\alpha} = \phi P_1^* \psi \quad a.e.$$

which allows to transform initial and final conditions μ, σ in ϕ, ψ , initial and final conditions for the equations satisfied by the drift and the one of its time reversed.

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