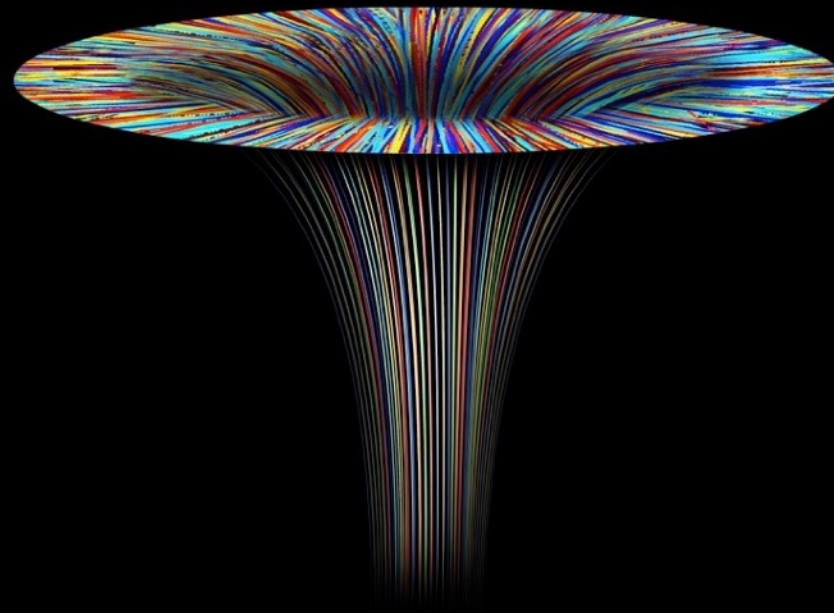
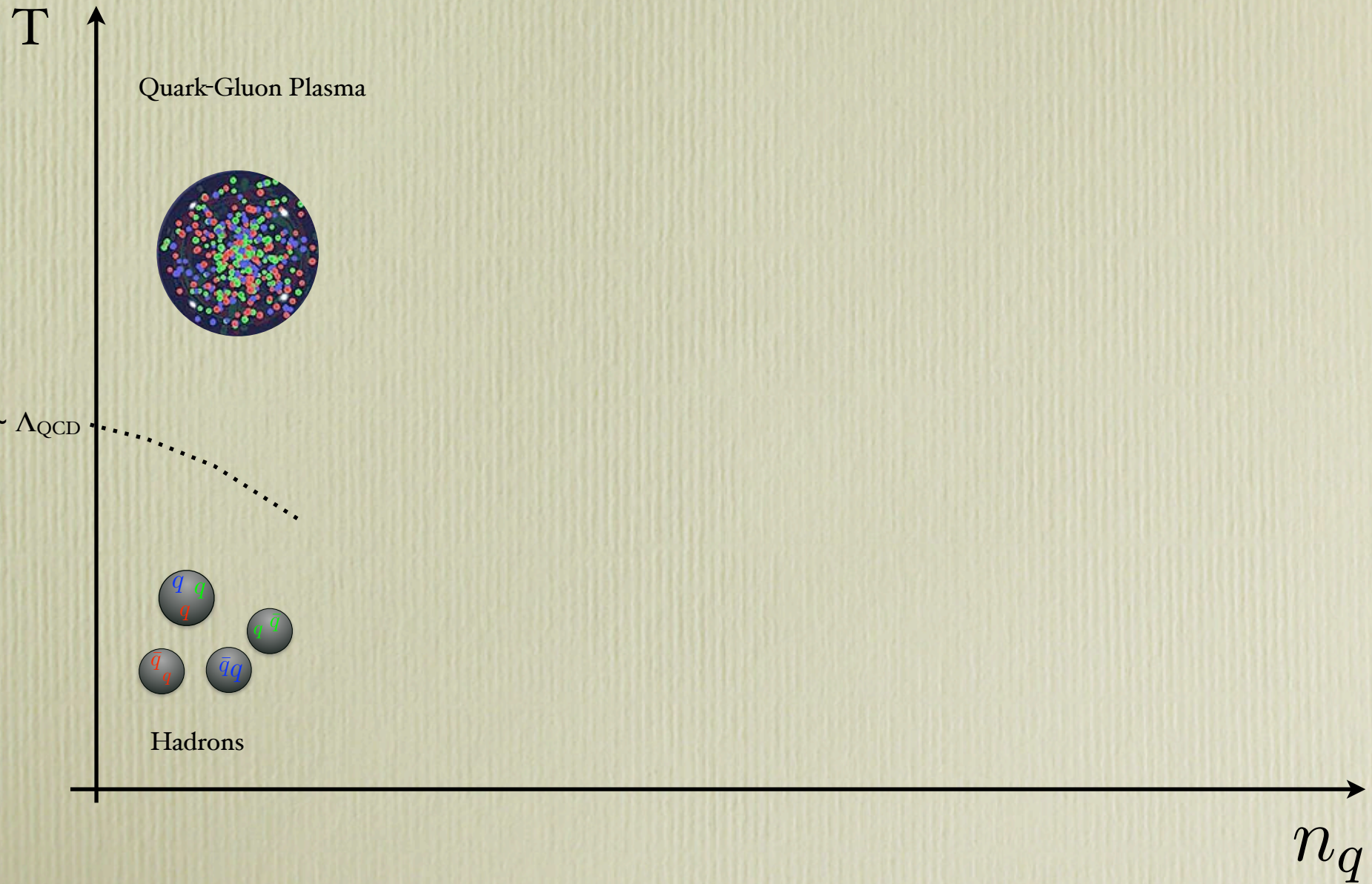


Hydrodynamics Near a Critical Point

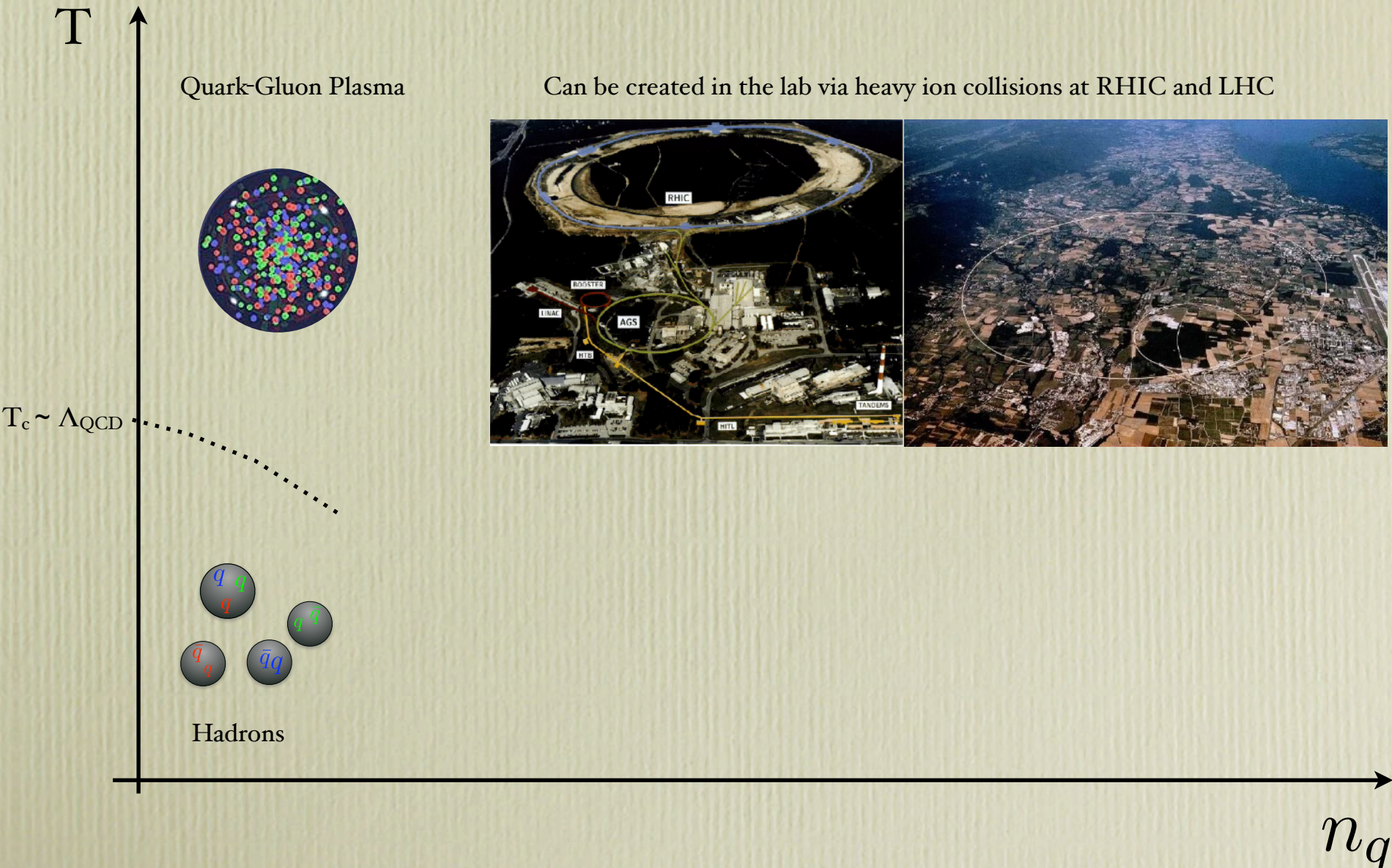


David Mateos
ICREA & University of Barcelona

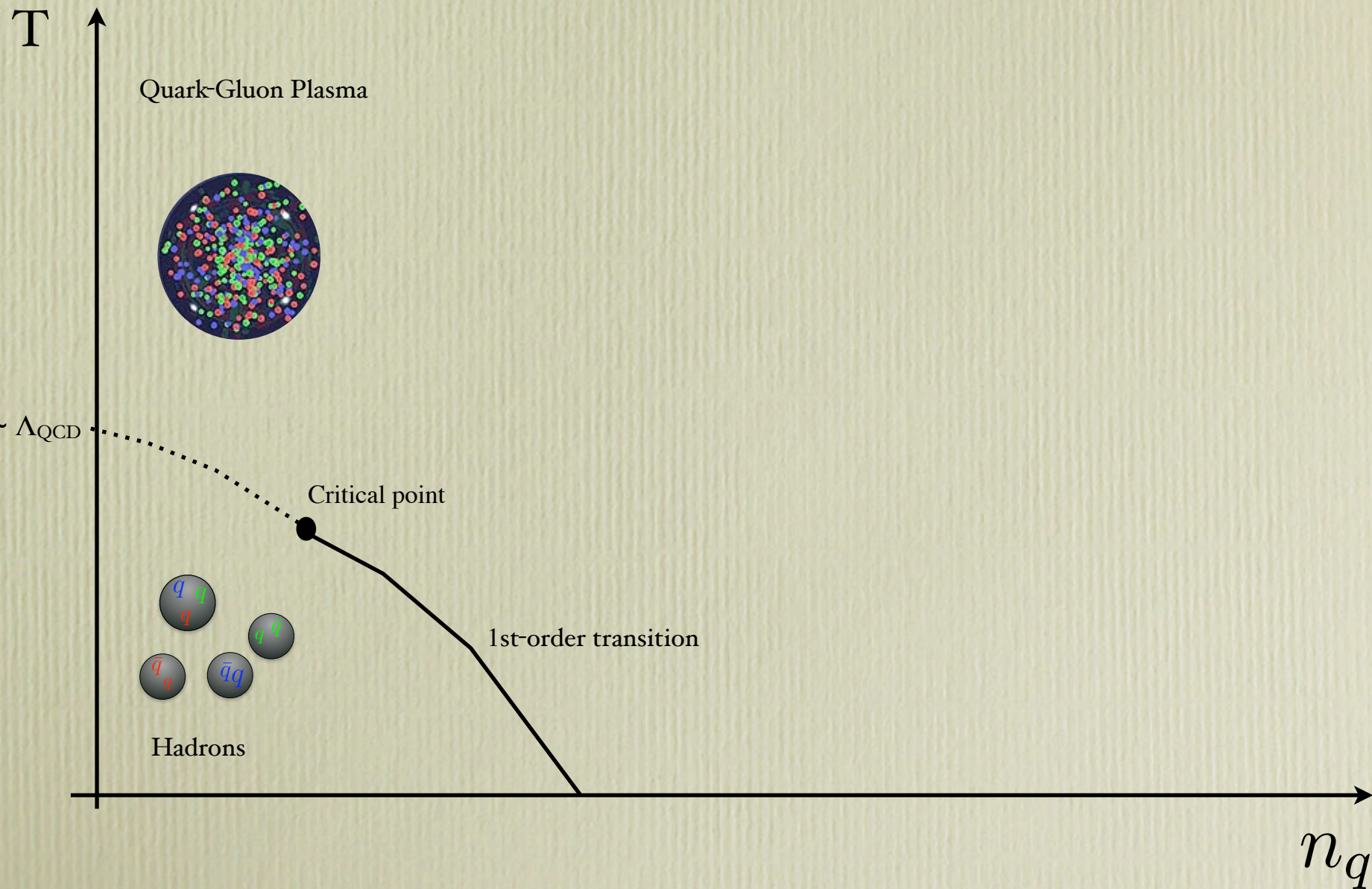
QCD phase diagram



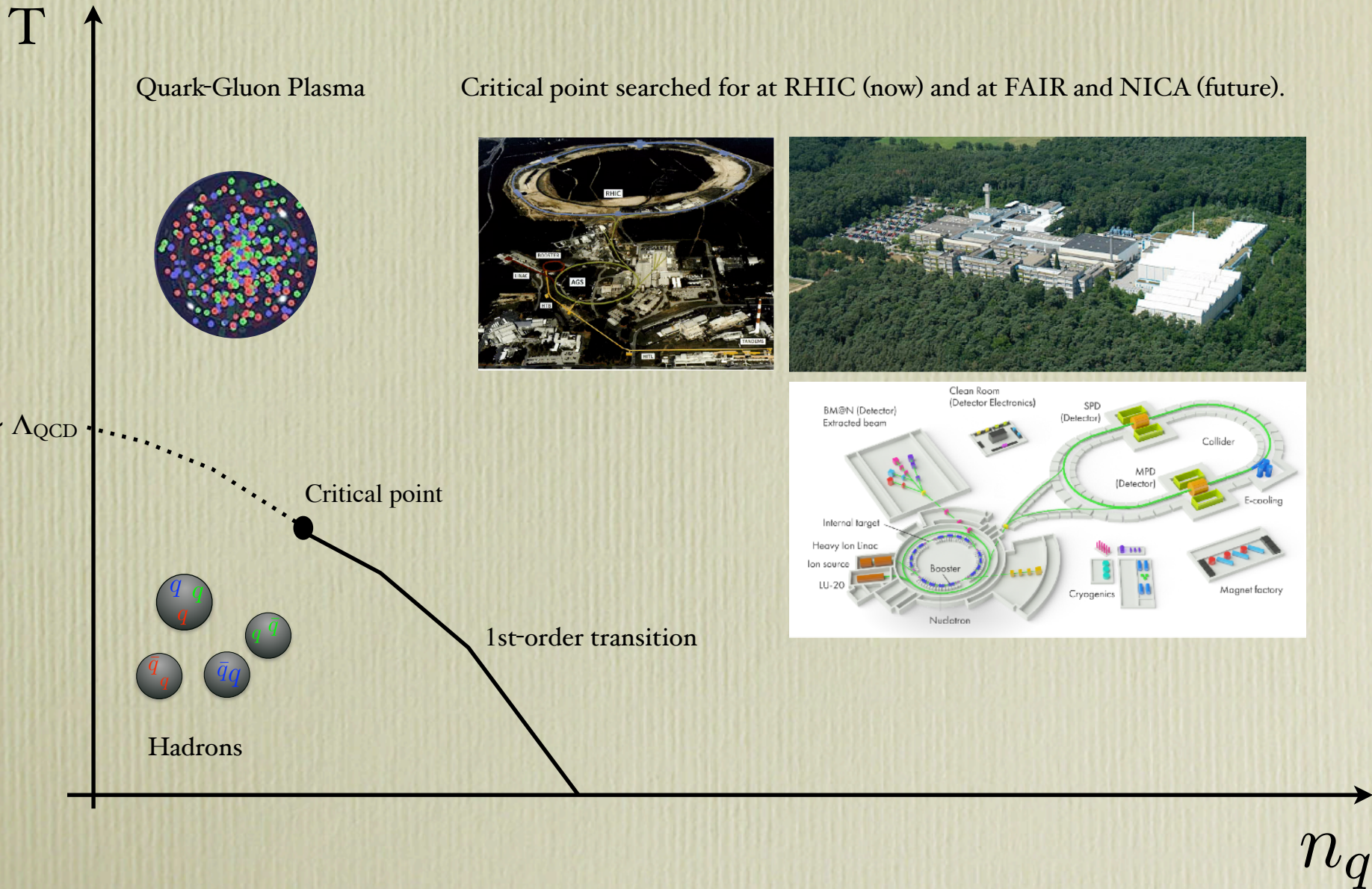
QCD phase diagram



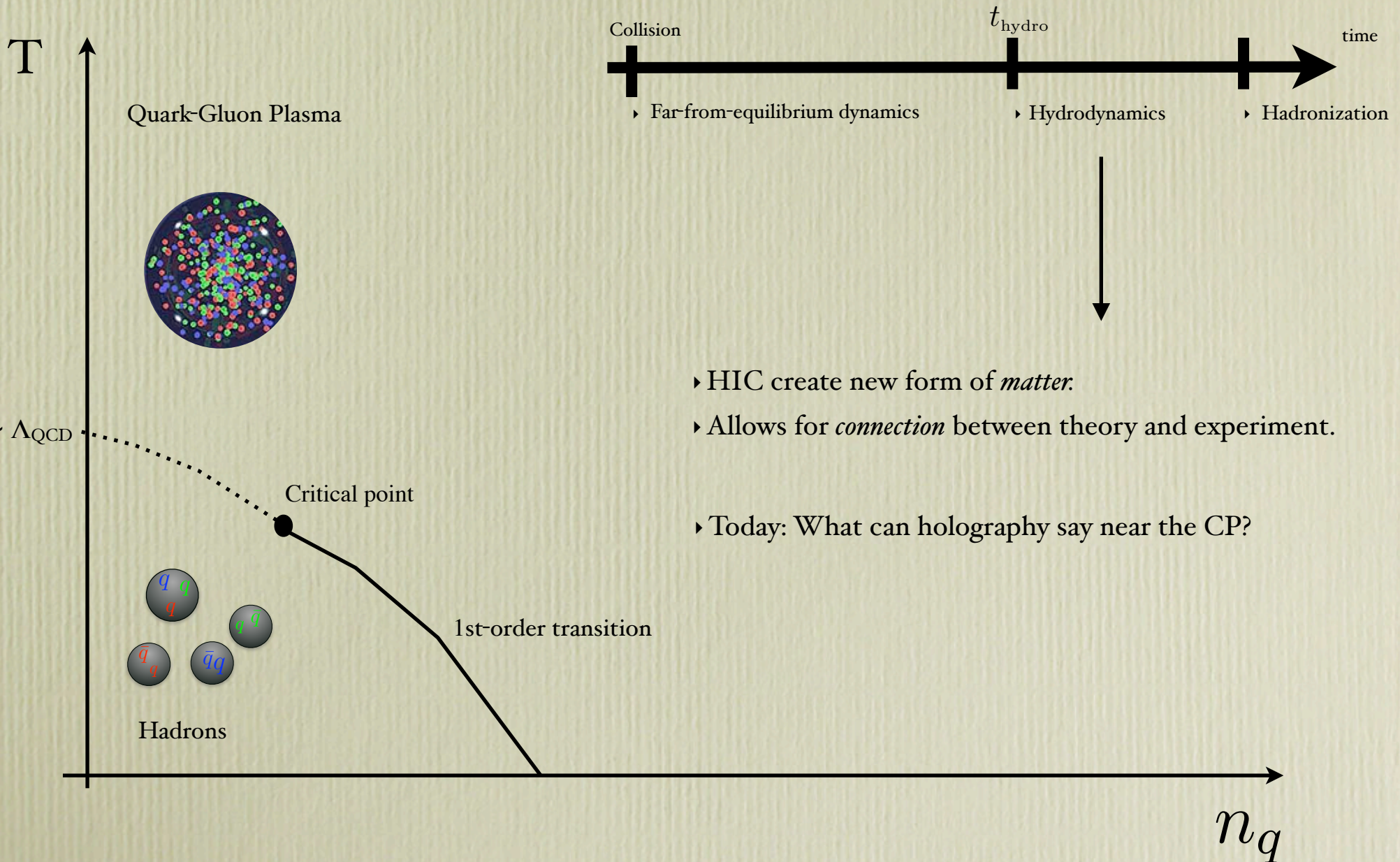
QCD phase diagram



QCD phase diagram



QCD phase diagram



- ▶ HIC create new form of *matter*.
- ▶ Allows for *connection* between theory and experiment.
- ▶ Today: What can holography say near the CP?

Plan

- Real-time dynamics of phase separation.
- Holographic collisions with phase transitions.
- What about critical fluctuations?

Dynamics of phase separation

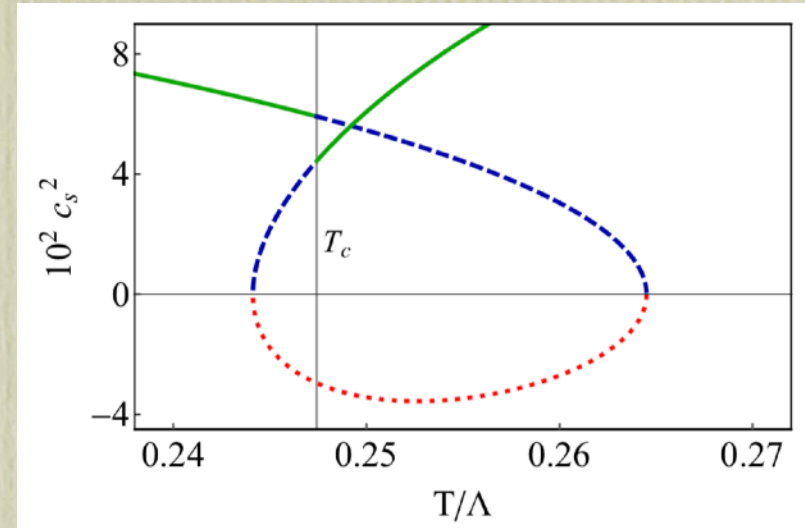
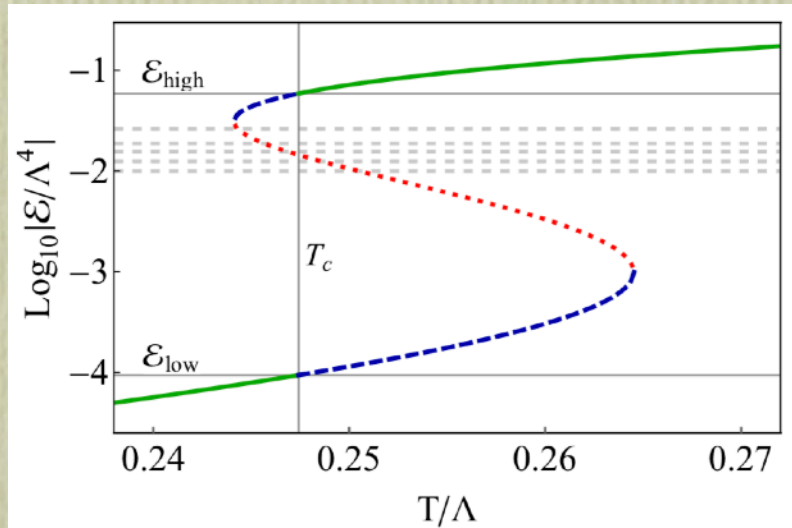
1st-order phase transition: Spinodal instability

Attems, Bea, Casalderrey, D.M., Triana & Zilhao '17

Janik, Jankowski, Soltanpanahi '17

Attems, Bea, Casalderrey, D.M. & Zilhao '19

Bellantuono, Janik, Jankowski, Soltanpanahi '19



- Thermodynamic instability implies dynamical instability:

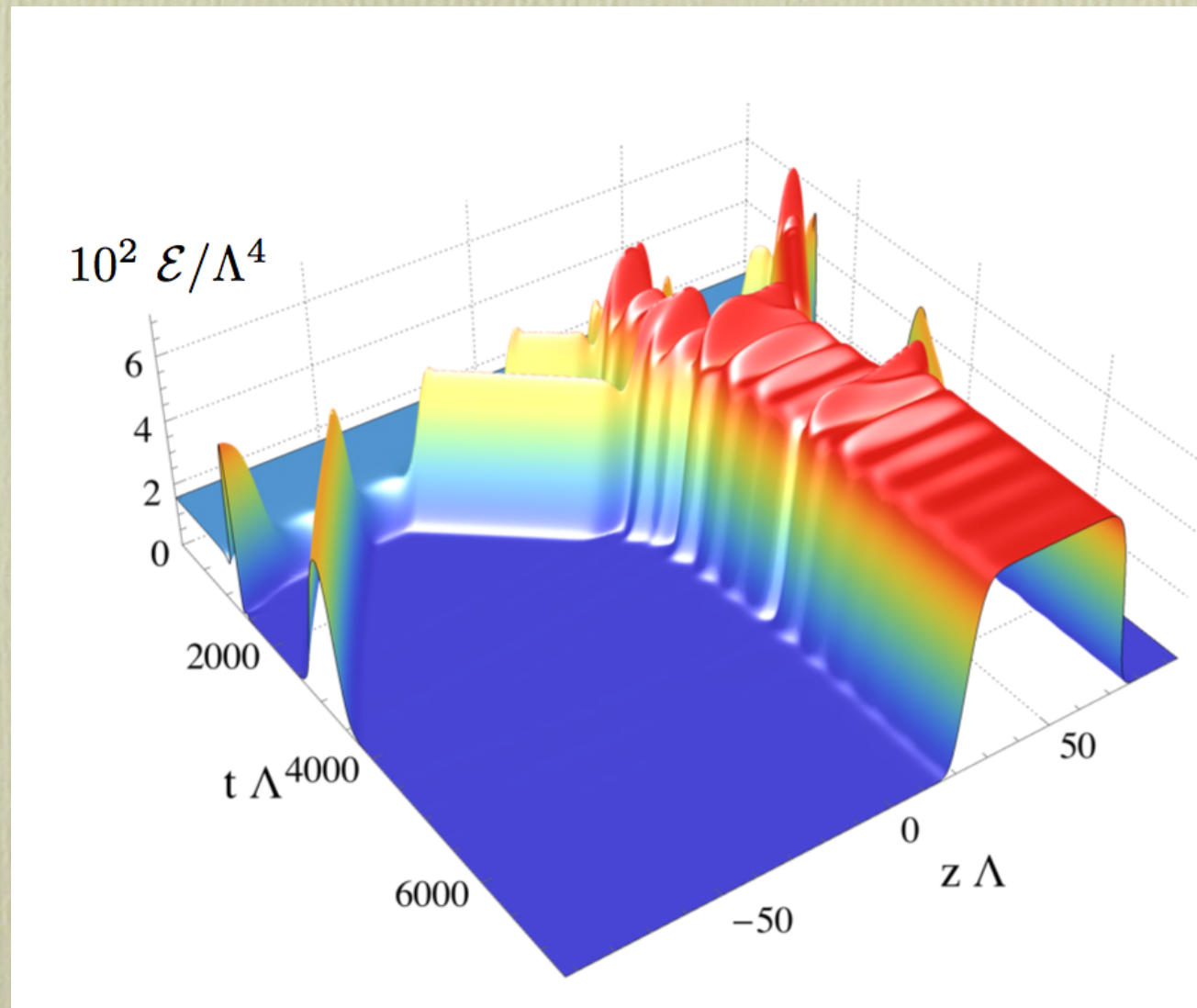
$$c_V < 0 \quad \rightarrow \quad c_s^2 = \frac{s}{c_V} < 0 \quad \rightarrow \quad c_s \text{ is imaginary}$$

$$\omega = c_s k \quad \rightarrow \quad e^{-i\omega t} = e^{+|c_s|kt}$$

1st-order phase transition: Phase separation

Attems, Bea, Casalderrey, D.M. & Zilhao '19

Perturbed homogeneous state evolves to phase-separated configuration:



1st-order phase transition: Phase separation

Attems, Bea, Casalderrey, D.M. & Zilhao '19

- Describing evolution in detail could fill an entire talk.
- Instead of that I will show you that entire evolution is well described by 2nd-order hydrodynamics.

Evolution described by 2nd-order hydrodynamics

Attems, Bea, Casalderrey, D.M., Triana & Zilhao '17

Attems, Bea, Casalderrey, D.M. & Zilhao '19

bulk & shear viscosities

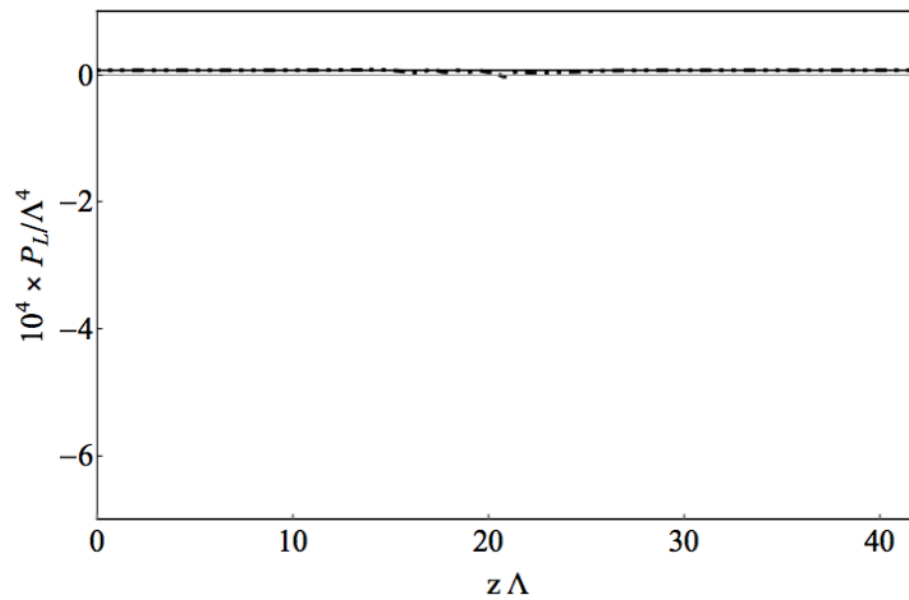
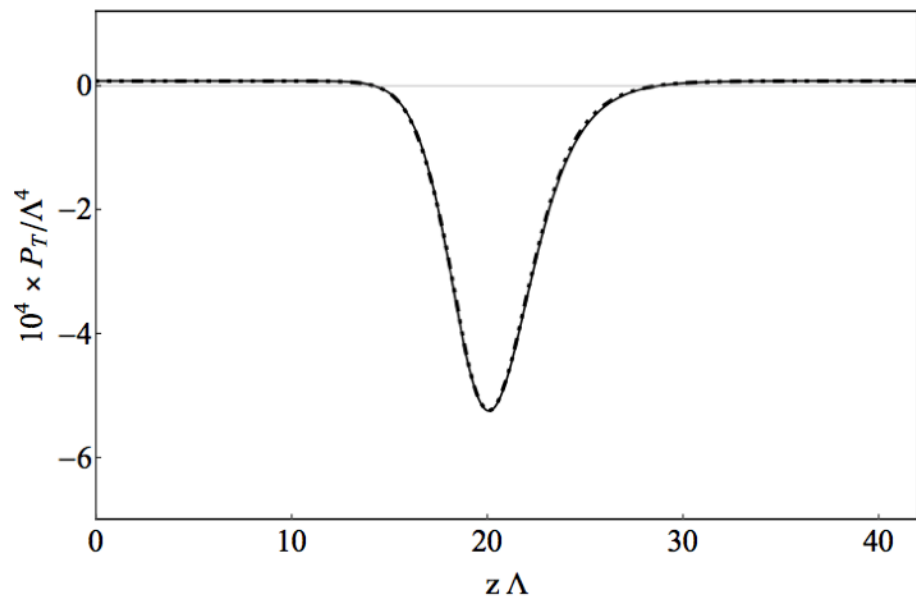


$$T_{\mu\nu} = T_{\mu\nu}^{\text{ideal}} + \partial_{\text{spatial}} + \partial_{\text{spatial}}^2$$

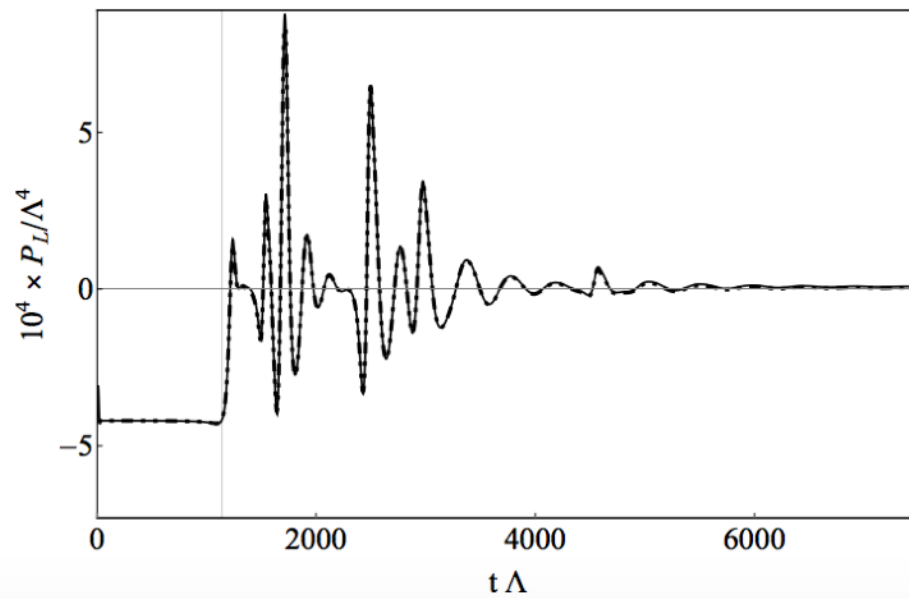
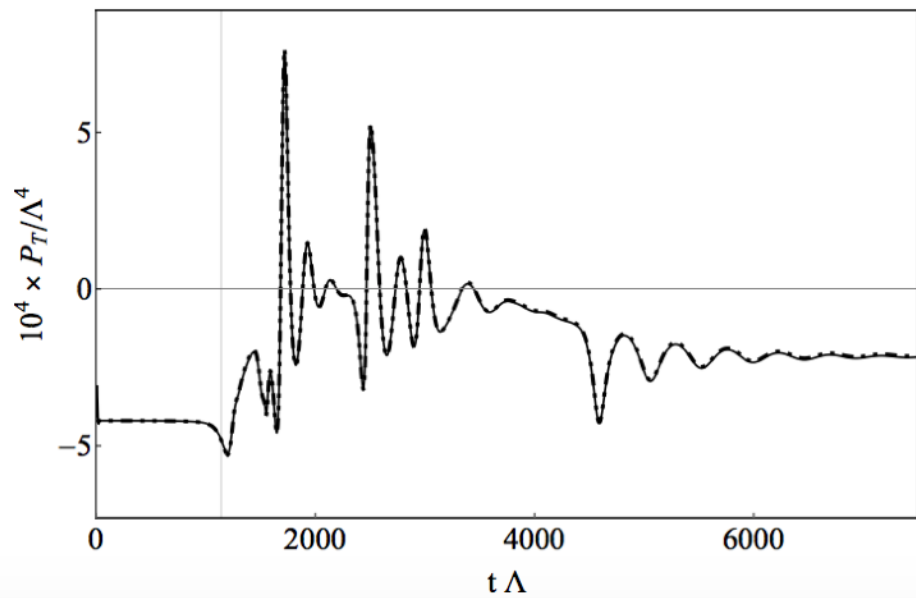
“Purely spatial formulation”

— P_T - - - P_{eq} ···· P^{hyd} ···· $P^{\text{hyd}(1)}$ - - - $P^{\text{hyd MIS}}$

Phase-separated configuration

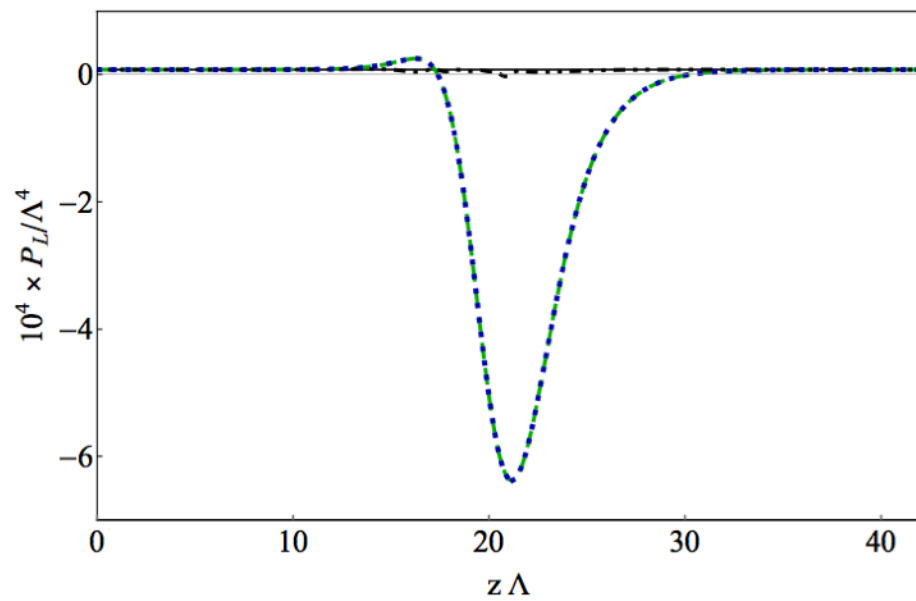
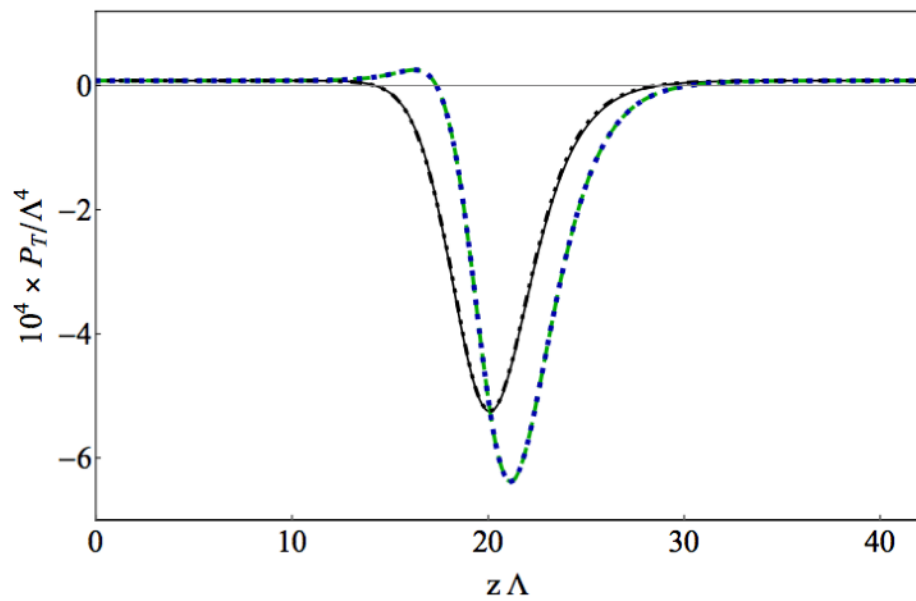


Time evolution at fixed z

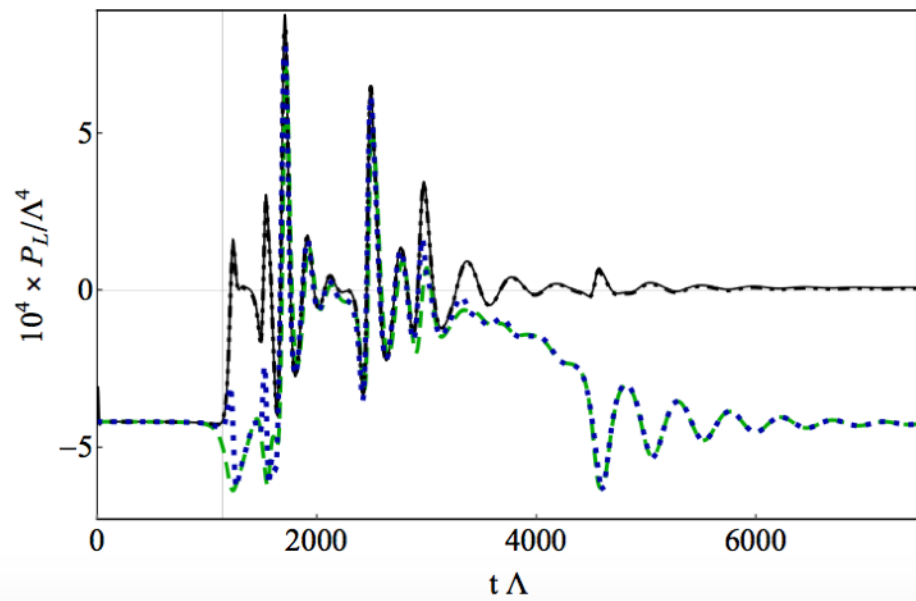
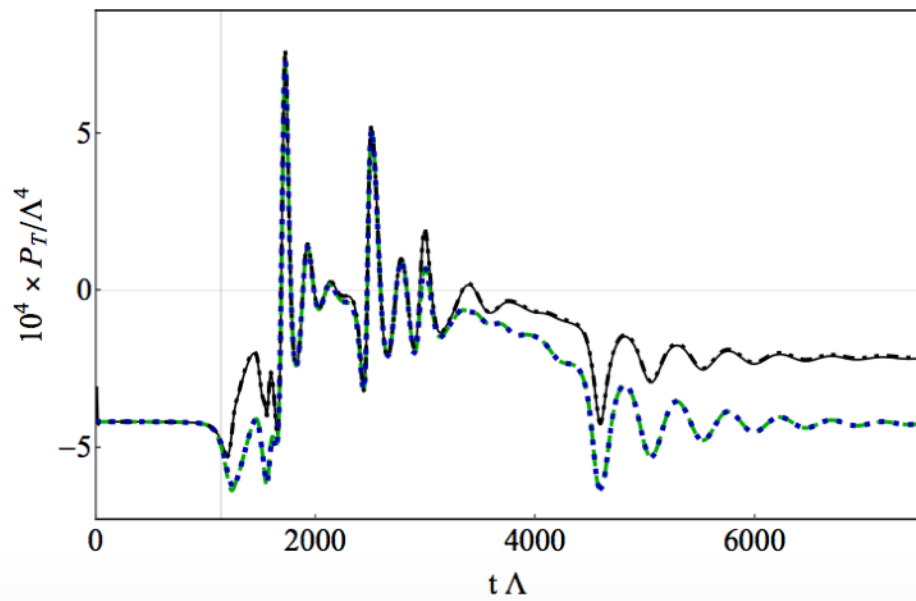


— P_T
- - - P_{eq}
... P^{hyd}
... $P^{\text{hyd}(1)}$
- - - $P^{\text{hyd MIS}}$

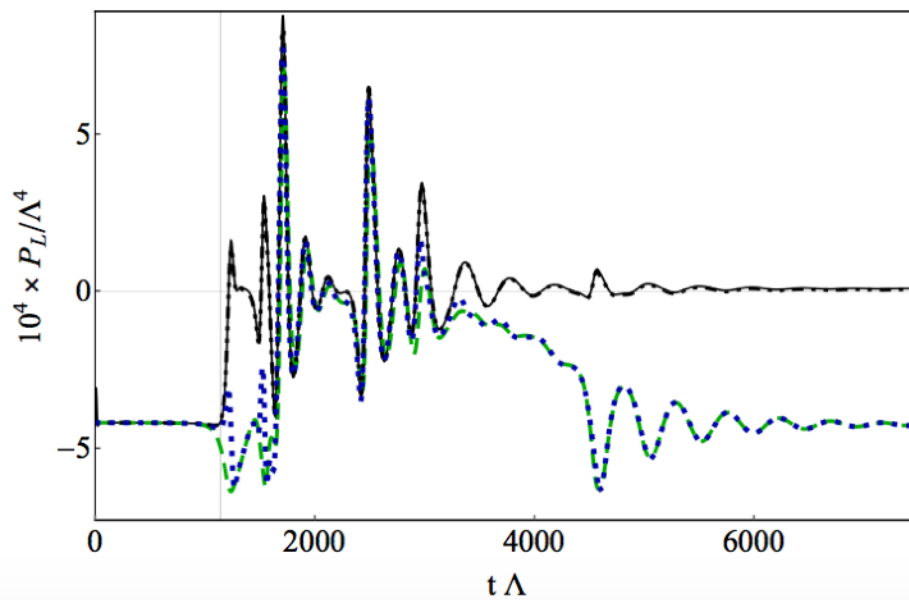
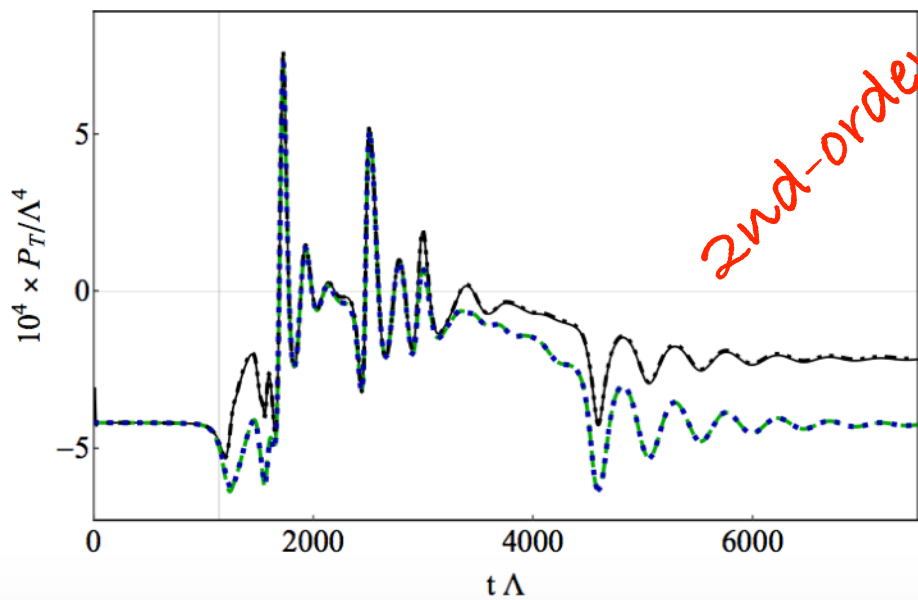
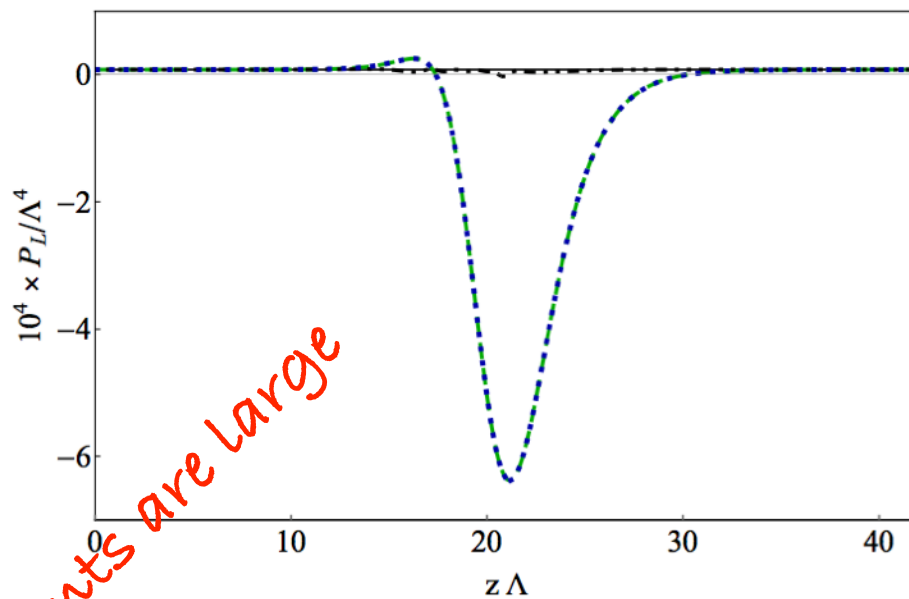
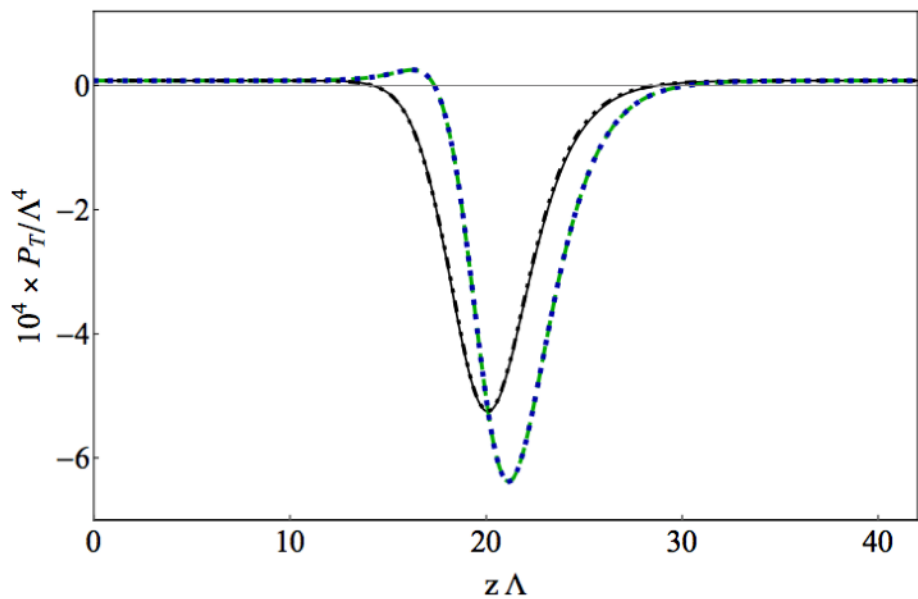
Phase-separated configuration



Time evolution at fixed z



— P_T
- - - P_{eq}
... p^{hyd}
· · · $p^{hyd(1)}$
- - - $p^{hyd MIS}$



2nd-order gradients are large

Evolution described by 2nd-order hydrodynamics

Attems, Bea, Casalderrey, D.M., Triana & Zilhao '17
Attems, Bea, Casalderrey, D.M. & Zilhao '19

- We are not doing time evolution, just checking constitutive relations.
- Problem for time evolution: Hydrodynamics is acausal.

$$T_{\mu\nu} = T_{\mu\nu}^{\text{ideal}} + \partial_{\text{spatial}} + \partial_{\text{spatial}}^2$$

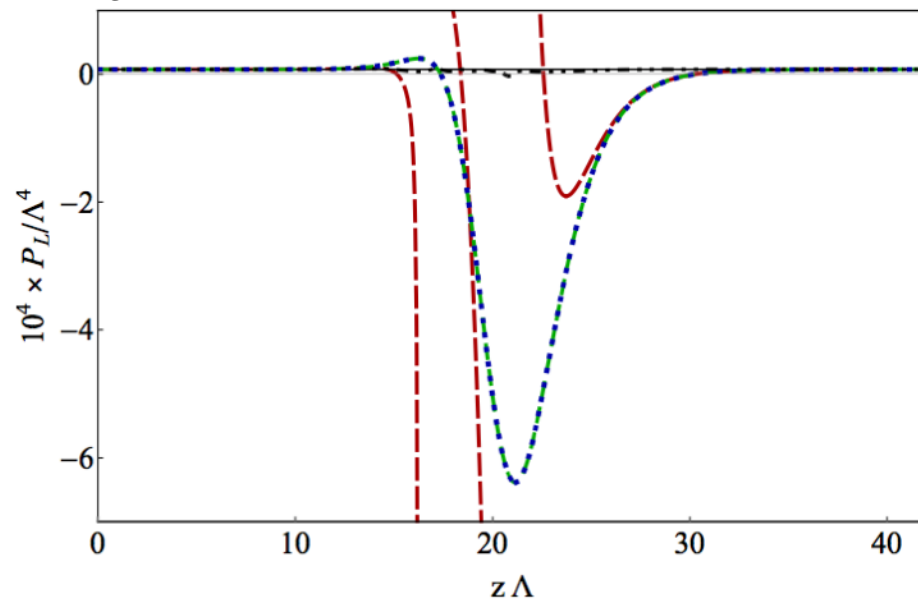
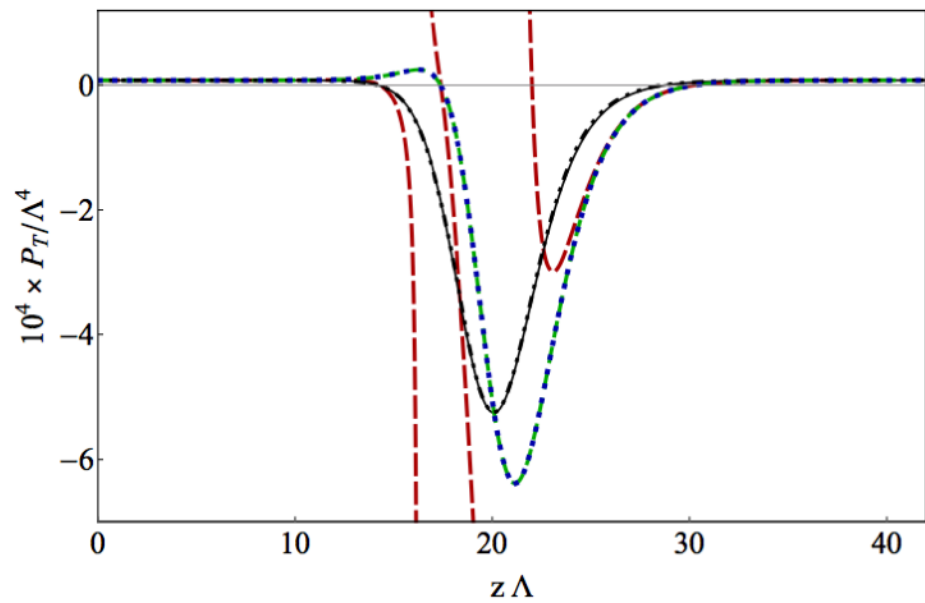
- One fix (Muller-Israel-Stewart): Use lower order equations to get:

$$T_{\mu\nu}^{\text{MIS}} = T_{\mu\nu}^{\text{ideal}} + \partial_{\text{spatial}} + \partial_{\text{spatial}} \partial_{\text{time}}$$

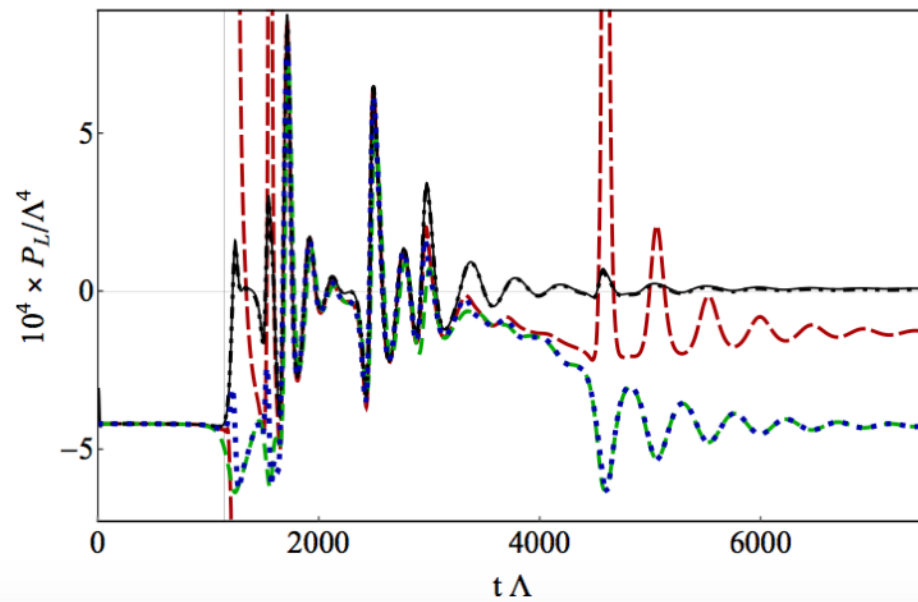
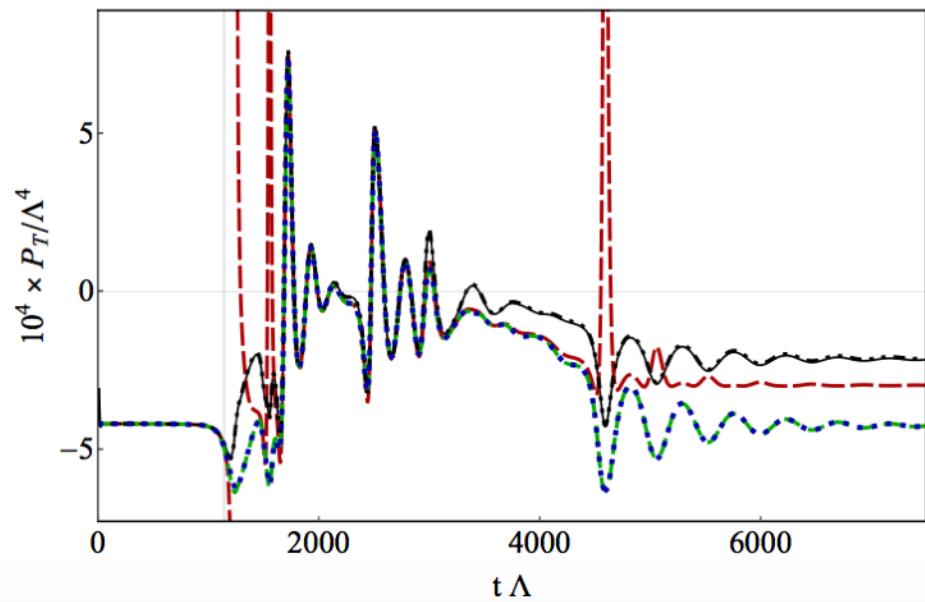
- Produces equivalent descriptions if gradients are small, but not in our case.

— P_T - - - P_{eq} ···· p^{hyd} ···· $p^{\text{hyd}(1)}$ - - - $p^{\text{hyd MIS}}$

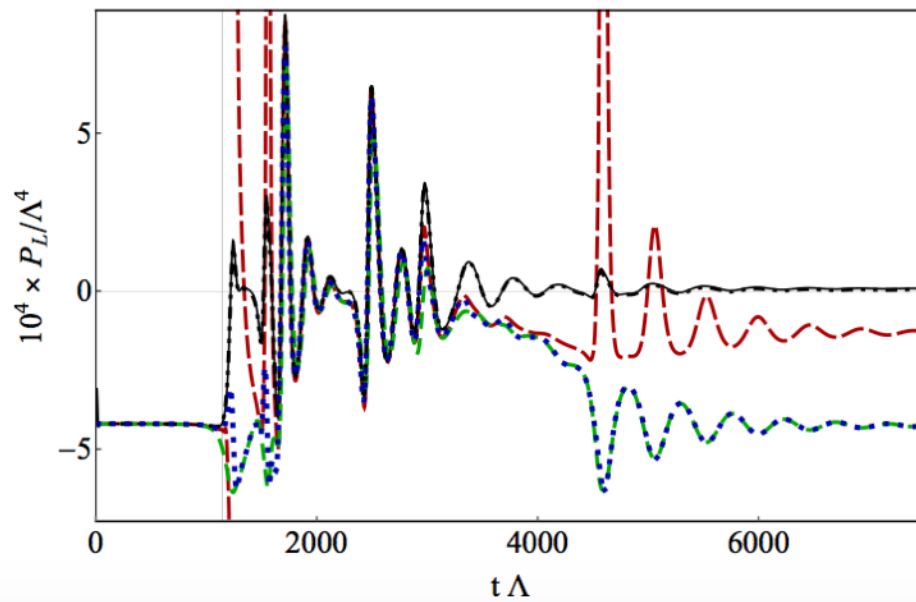
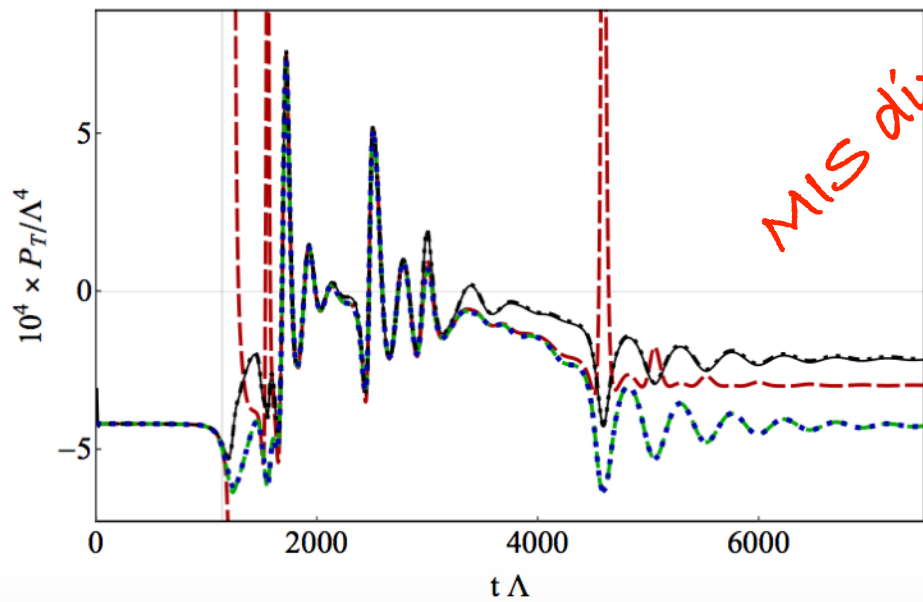
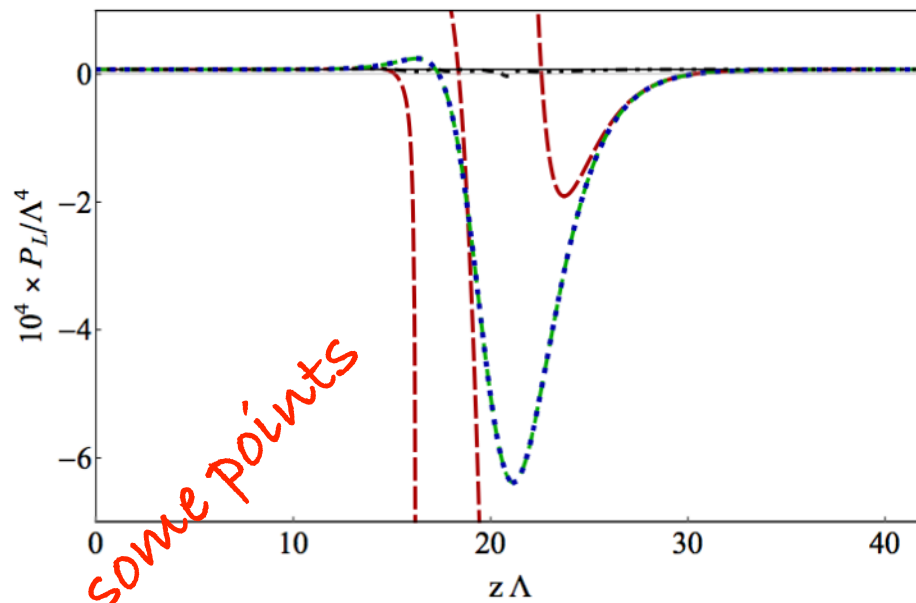
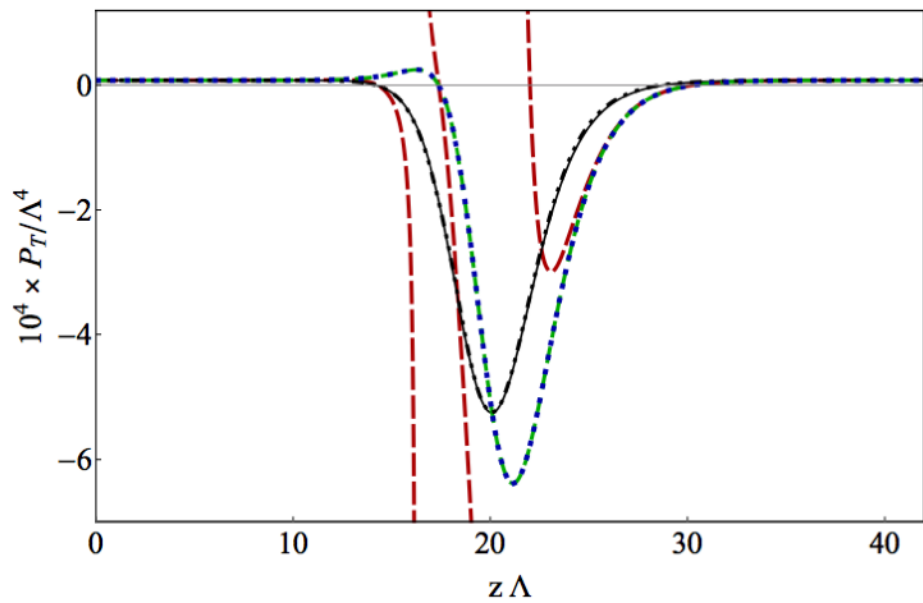
Phase-separated configuration



Time evolution at fixed z



— P_T - - - P_{eq} . . . p^{hyd} . . . $p^{hyd(1)}$ - - - $p^{hyd MIS}$

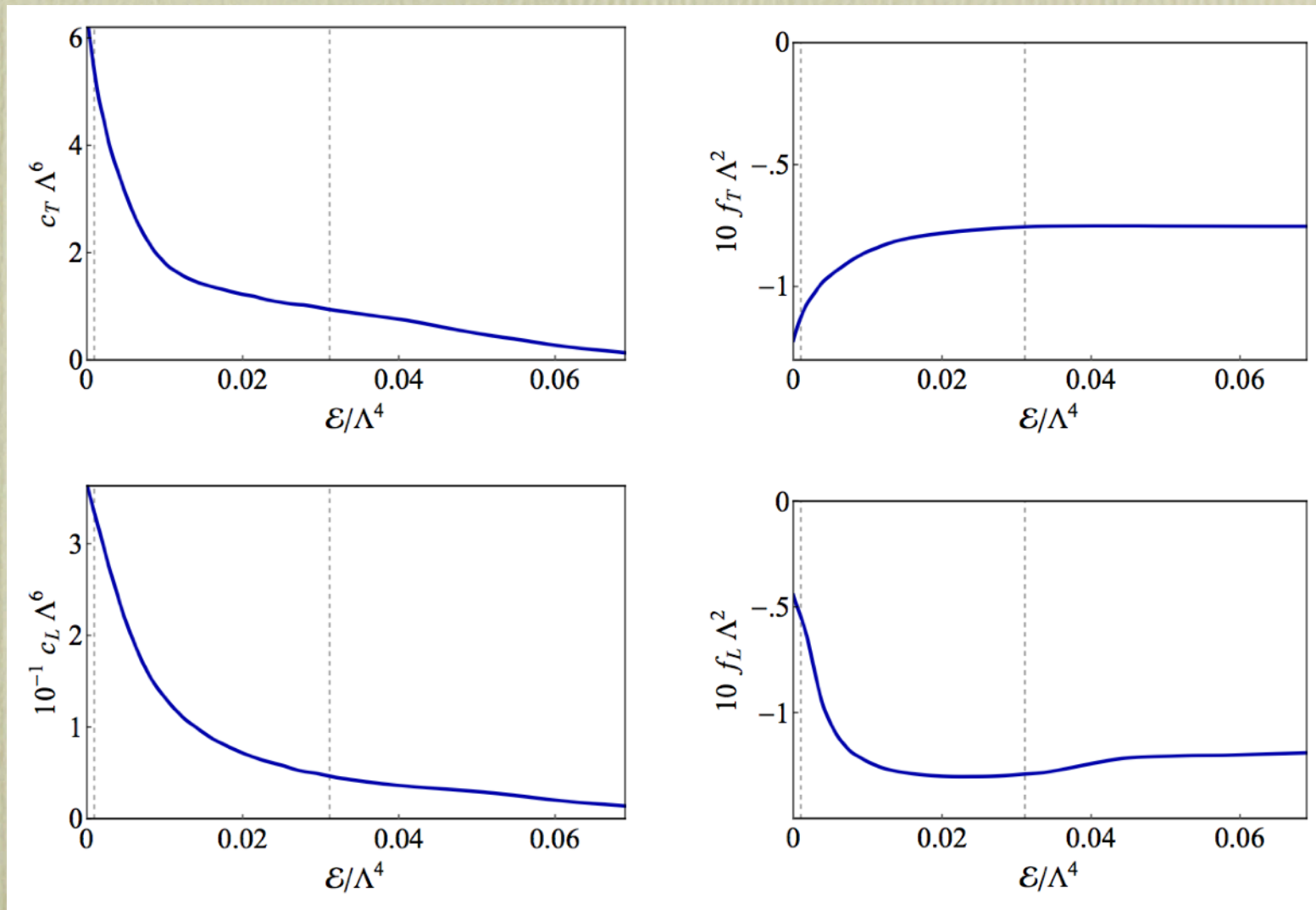


MIS diverges at some points

Purely spatial coefficients are smooth and finite

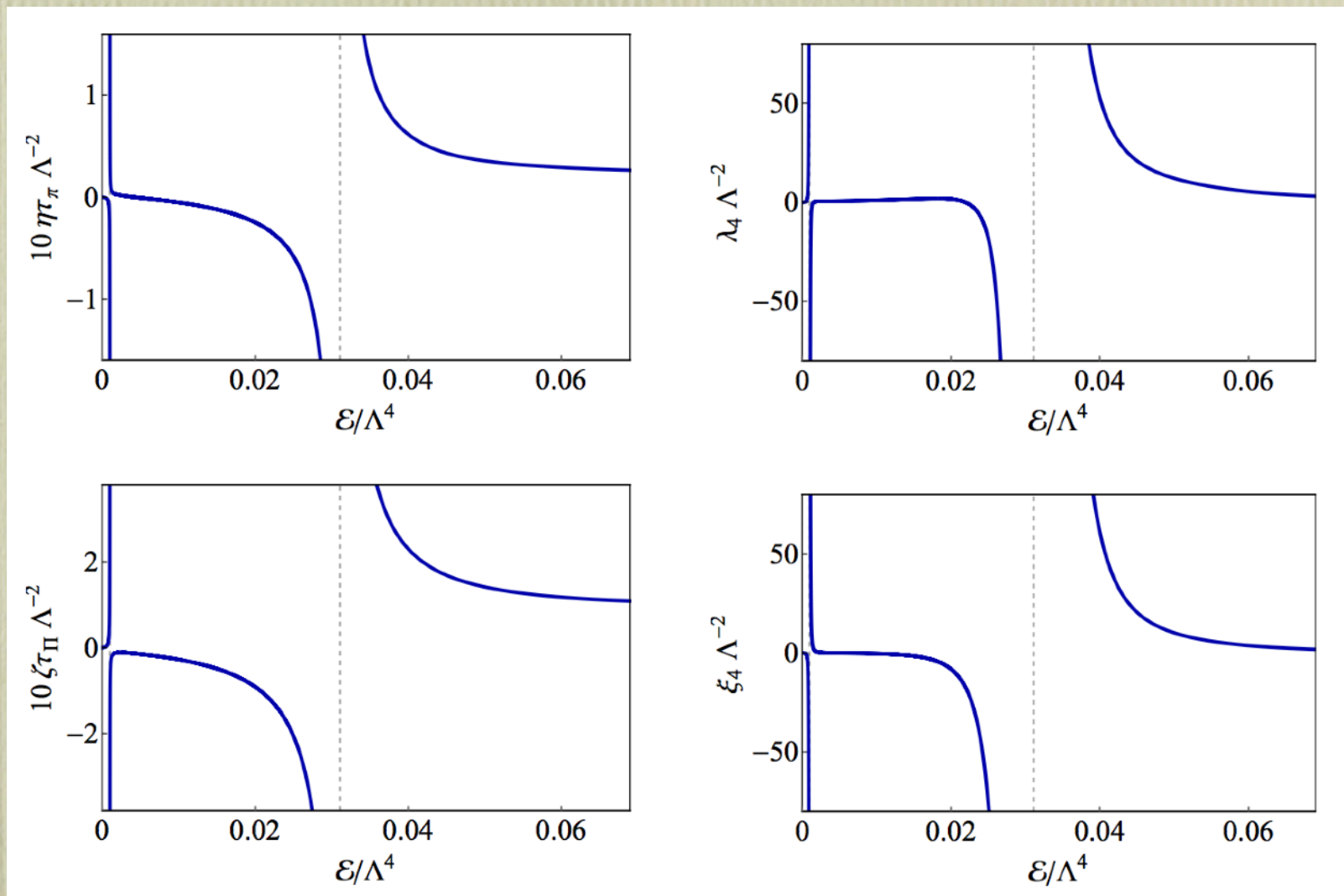
Attems, Bea, Casalderrey, D.M., Triana & Zilhao '17

Attems, Bea, Casalderrey, D.M. & Zilhao '19



MIS coefficients diverge at points where $c_s=0$

Attems, Bea, Casalderrey, D.M., Triana & Zilhao '17
Attems, Bea, Casalderrey, D.M. & Zilhao '19

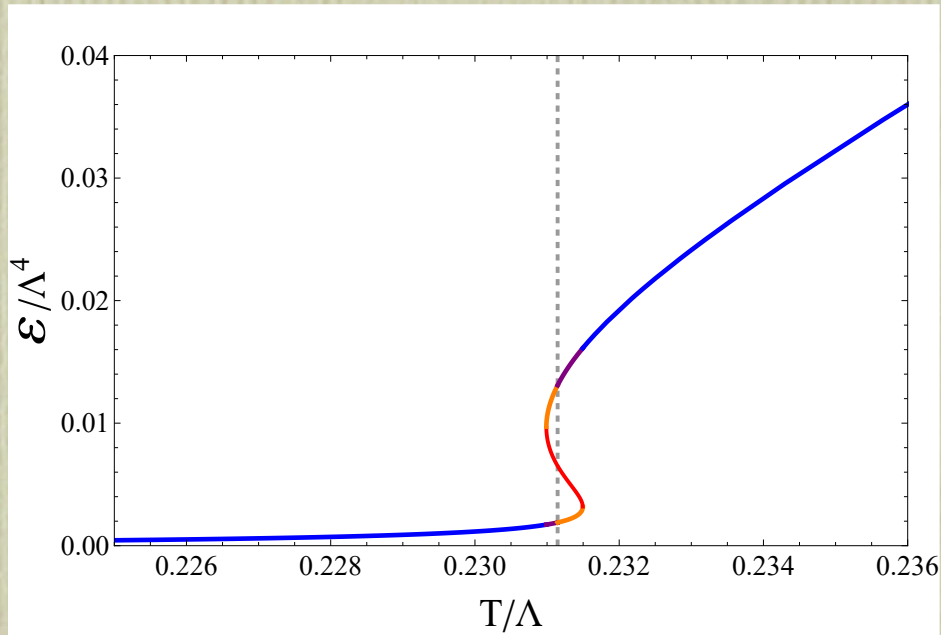


Change of basis involves powers of $1/c_s$

Collisions across a phase transition

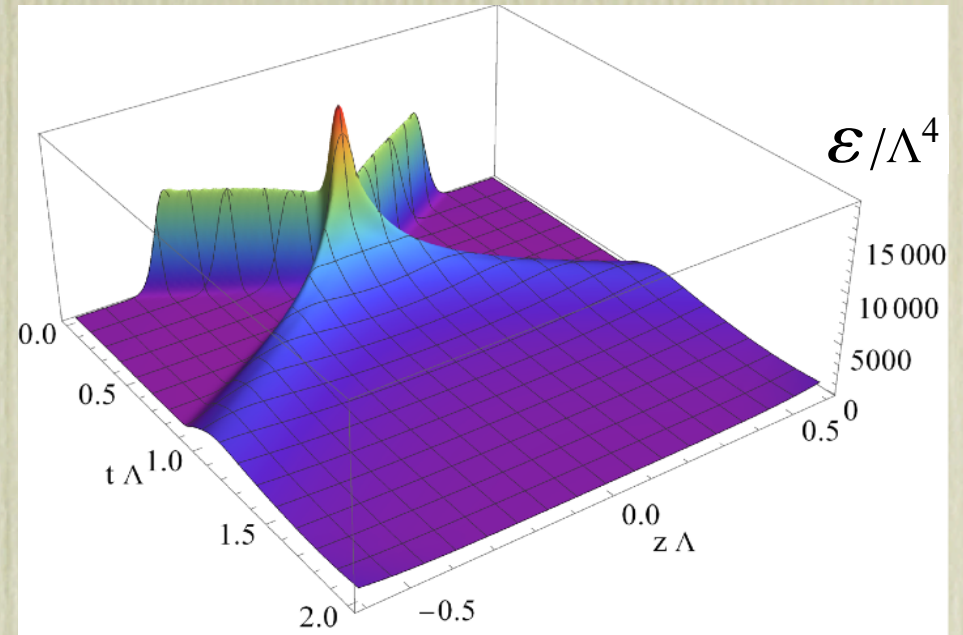
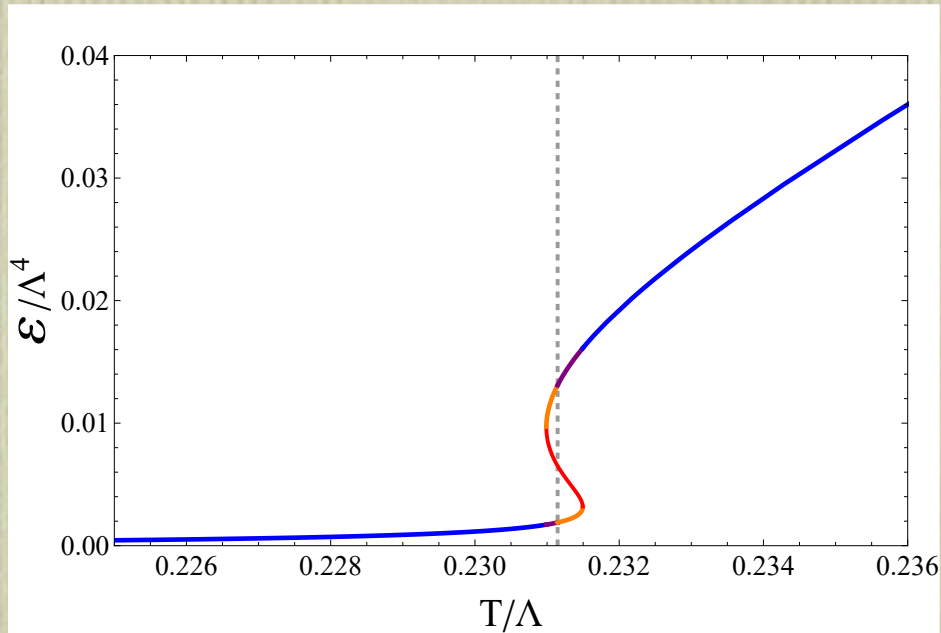
Collisions across a 1st-order phase transition

Attems, Bea, Casalderrey, D.M., Triana & Zilhao '18



Collisions across a 1st-order phase transition

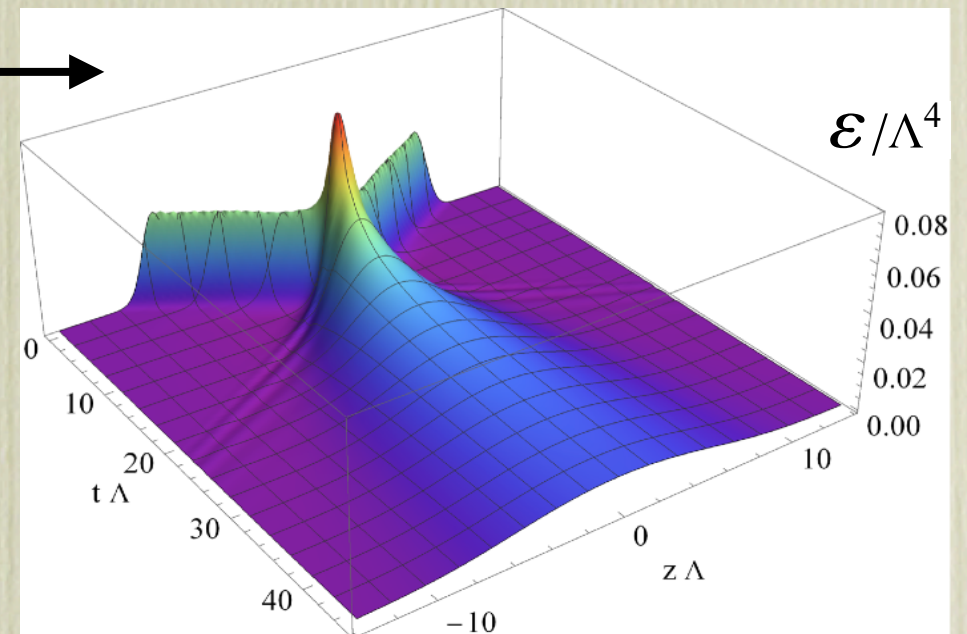
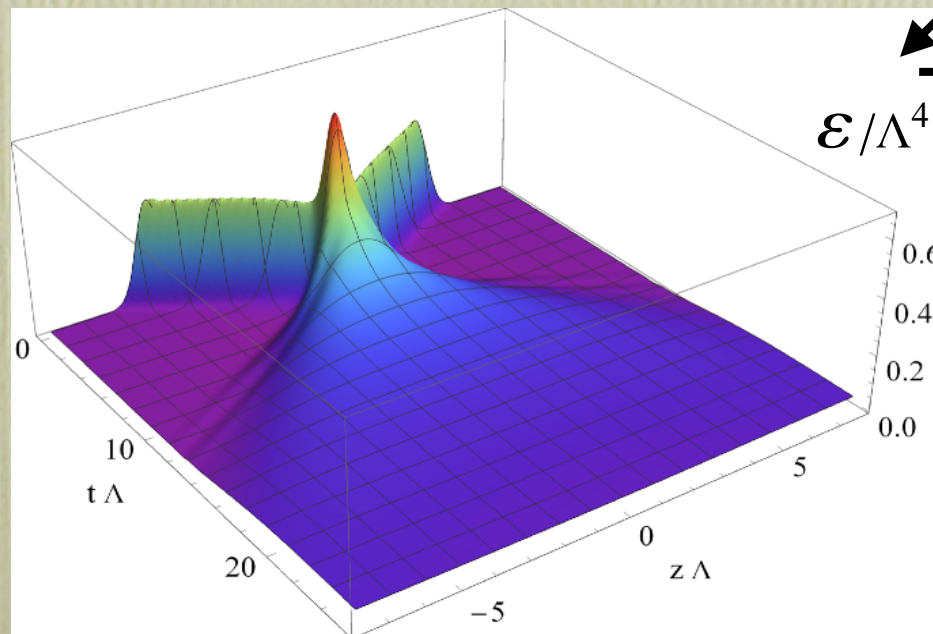
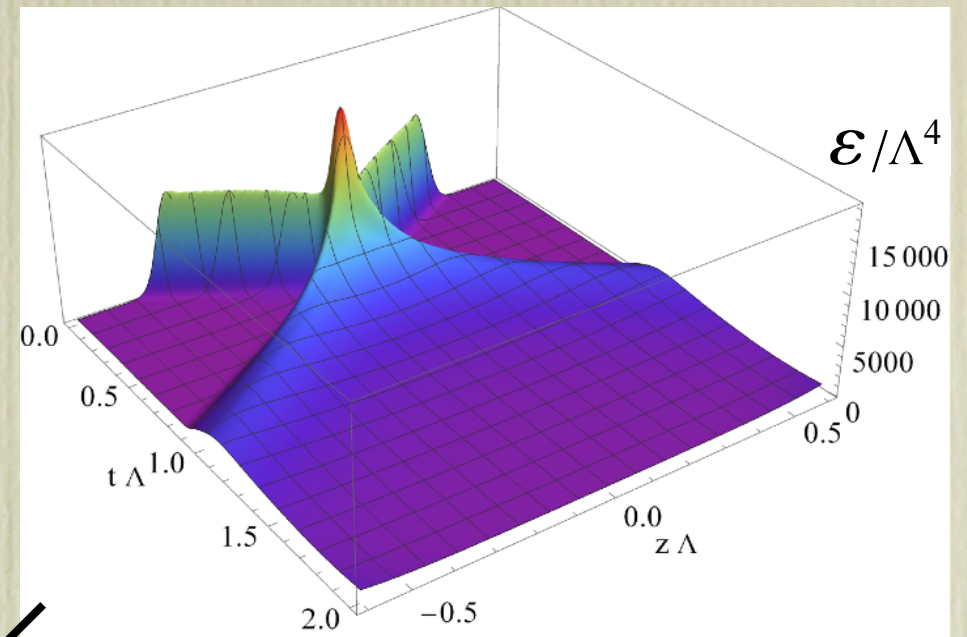
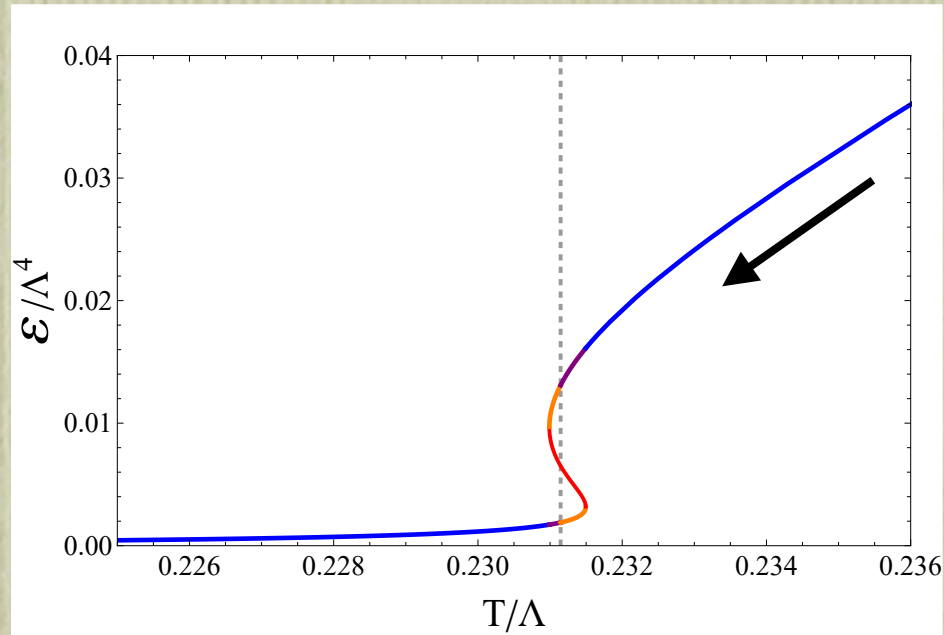
Attems, Bea, Casalderey, D.M., Triana & Zilhao '18



Extremely high energy:
Recover CFT result

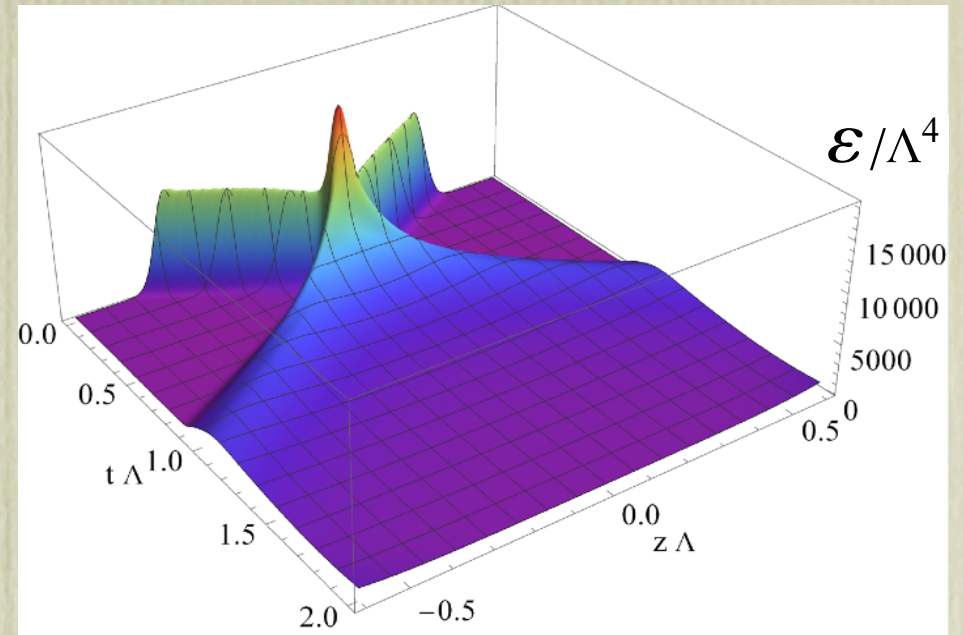
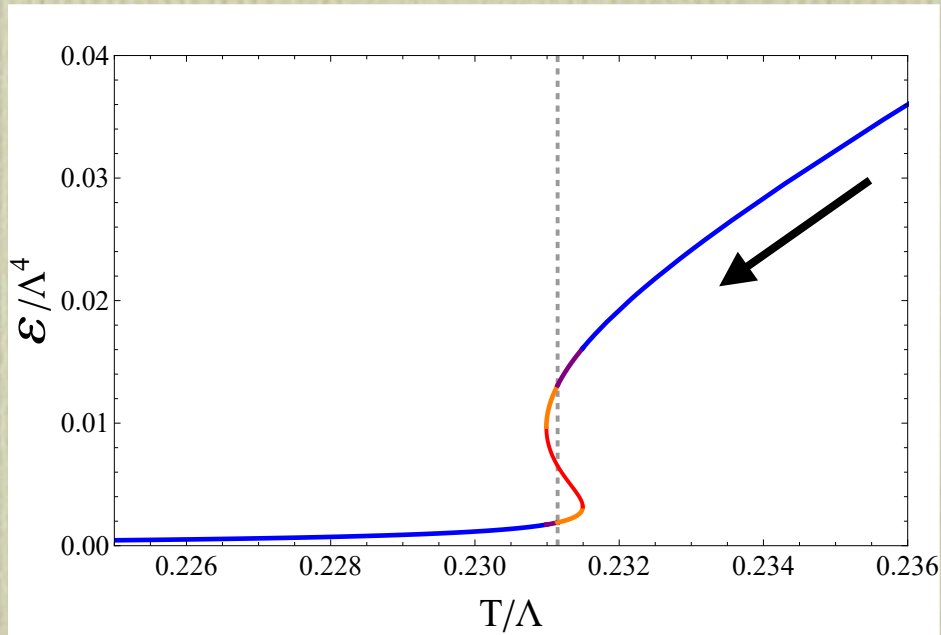
Collisions across a 1st-order phase transition

Attems, Bea, Casalderrey, D.M., Triana & Zilhao '18

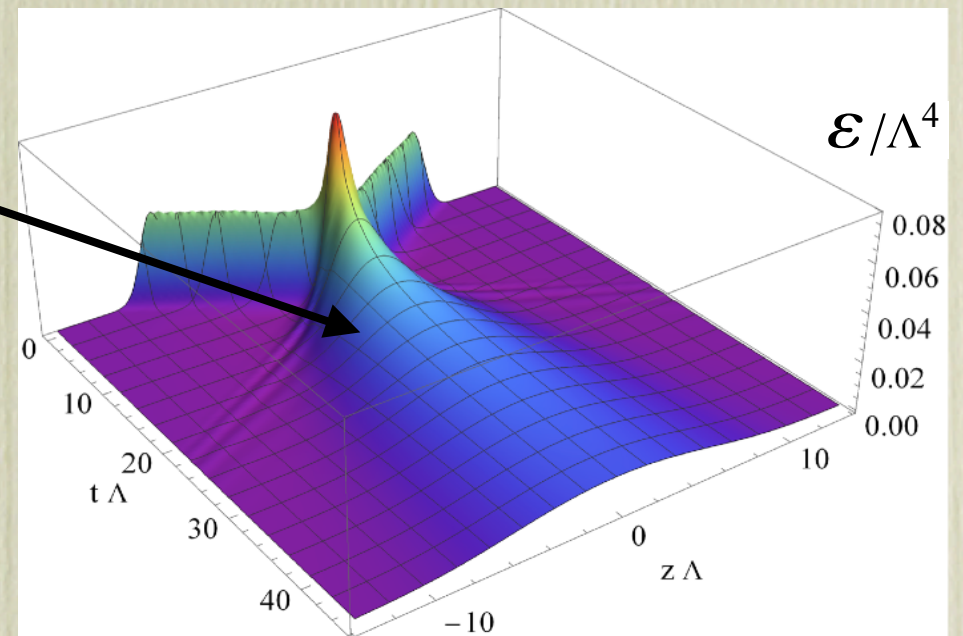


Collisions across a 1st-order phase transition

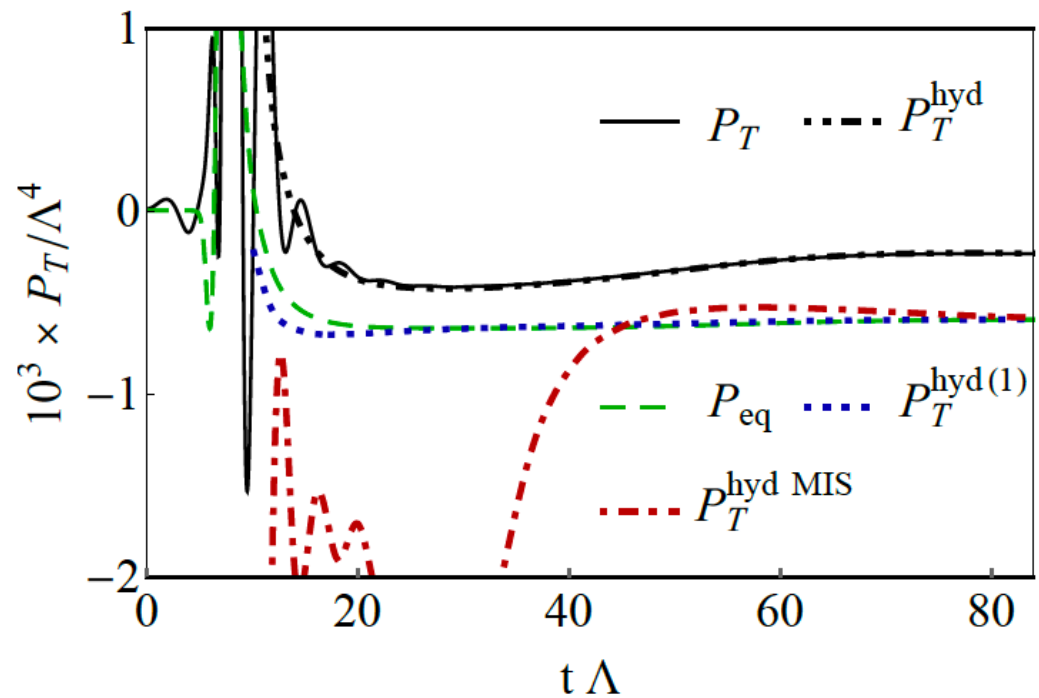
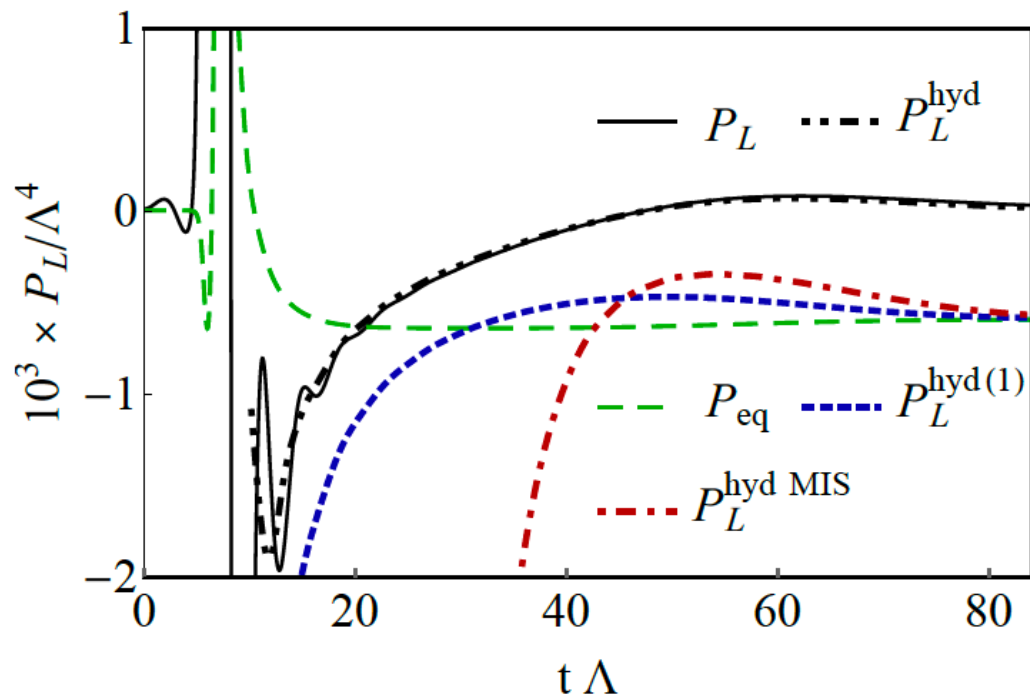
Attems, Bea, Casalderrey, D.M., Triana & Zilhao '18



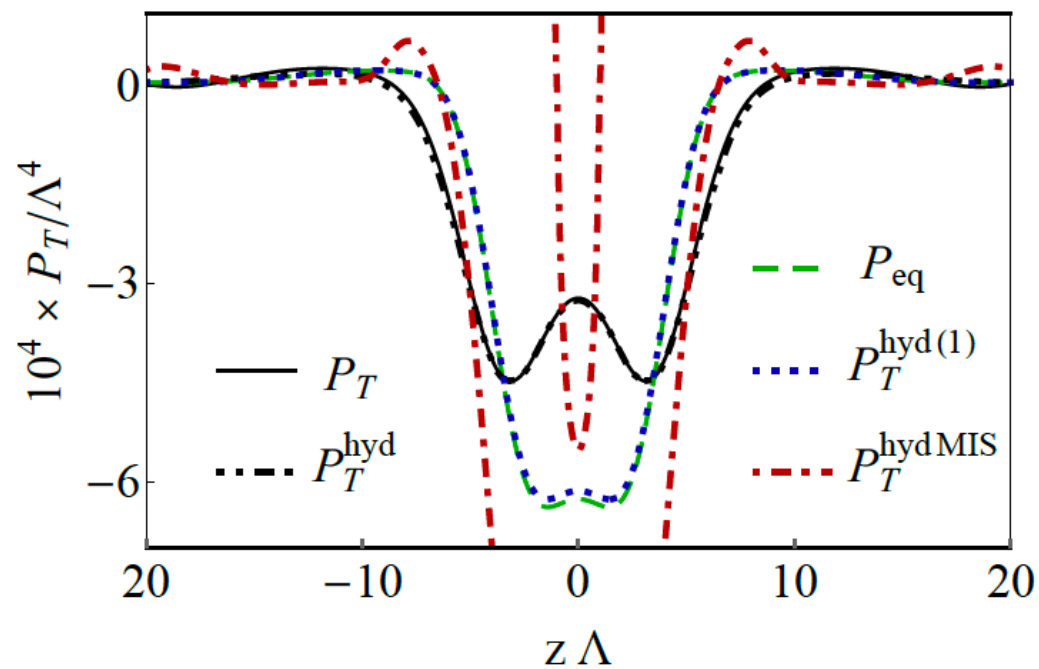
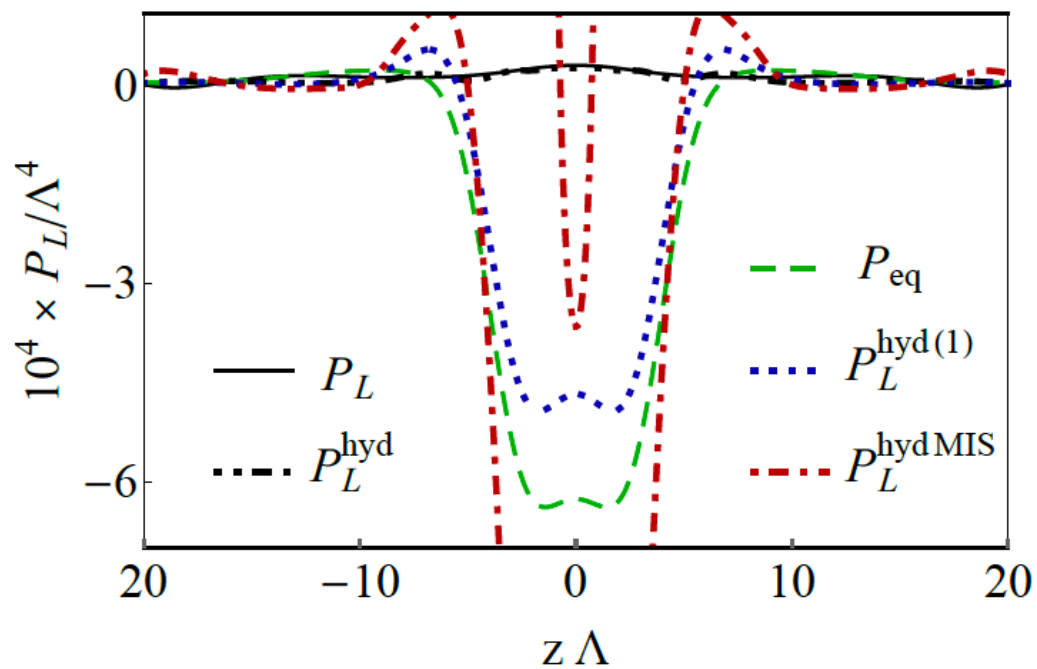
- Long-lived, quasi-static blob
- Well described by 2nd order purely spatial hydro but not MIS



Time evolution at mid-rapidity



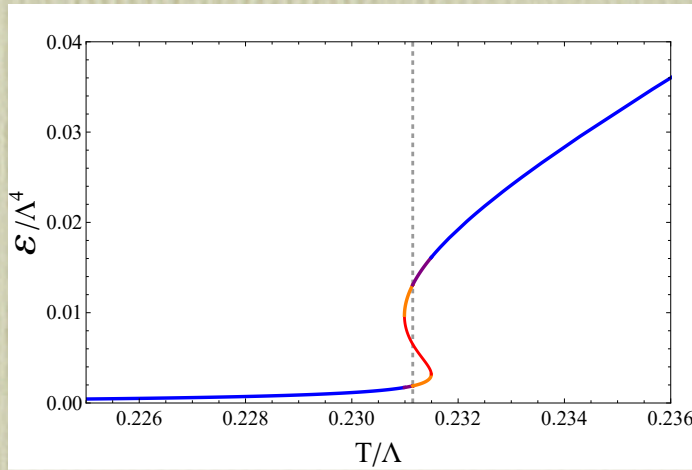
Snapshots of spatial profile after hydrodynamization



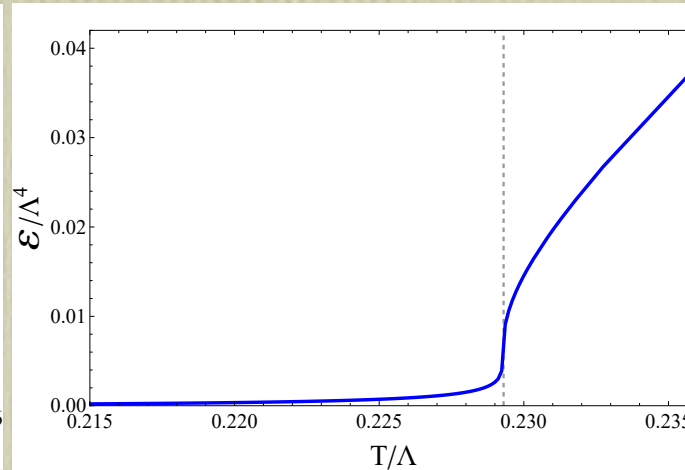
From 1st-order to 2nd-order to crossover

Attems, Bea, Casalderrey, D.M., Triana & Zilhao '18

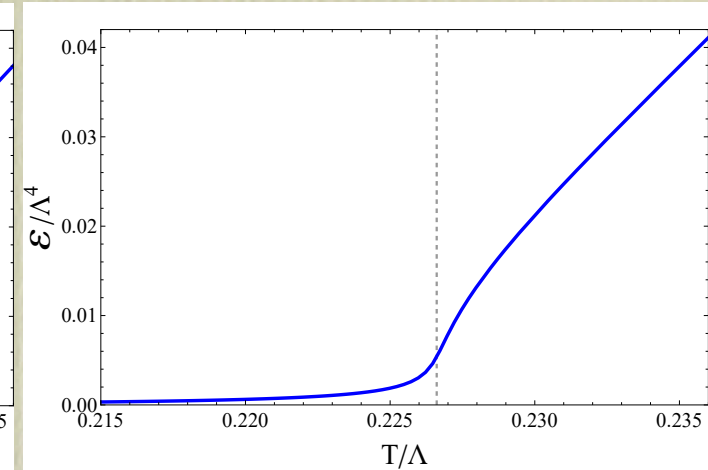
Continuous parameter \longrightarrow



1st order



2nd order

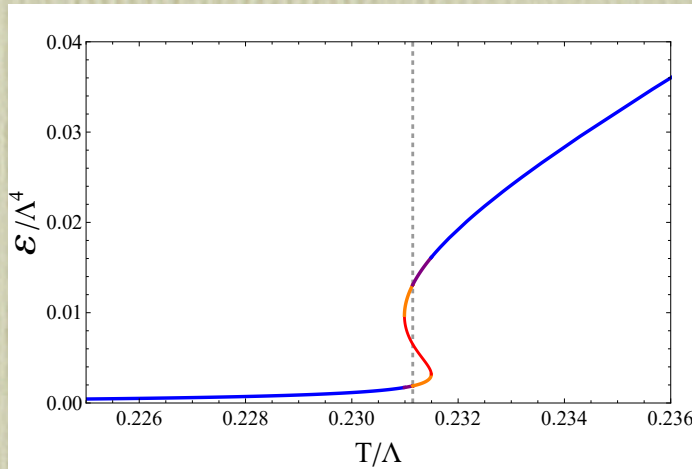


Crossover

From 1st-order to 2nd-order to crossover

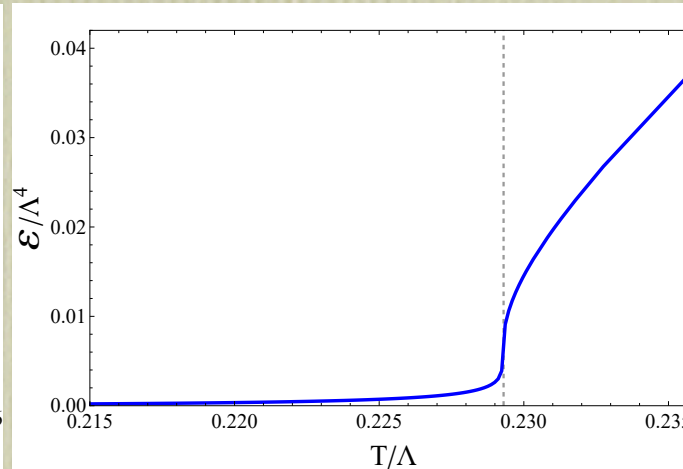
Attems, Bea, Casalderrey, D.M., Triana & Zilhao '18

Continuous parameter \longrightarrow



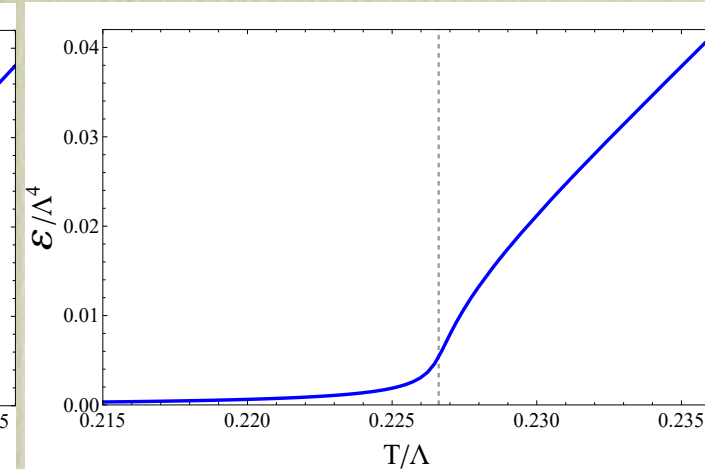
1st order

Non-zero latent heat



2nd order

Infinite correlation length



Crossover

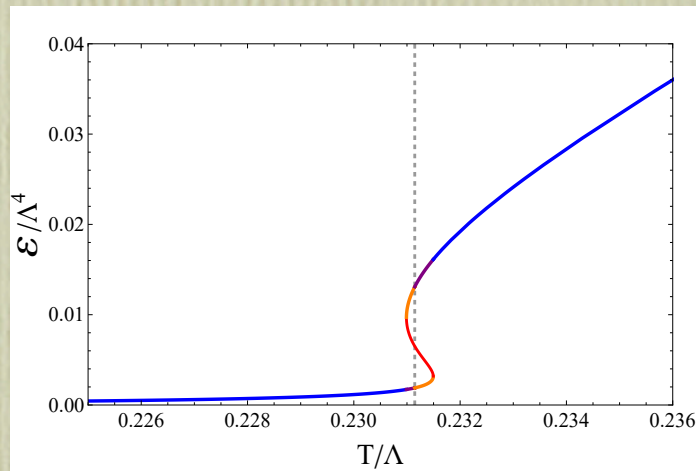
Neither of the above

Equilibrium physics is qualitatively very different

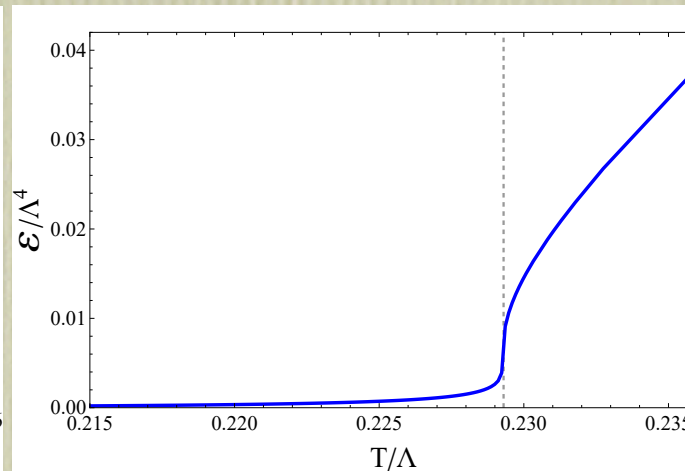
From 1st-order to 2nd-order to crossover

Attems, Bea, Casalderrey, D.M., Triana & Zilhao '18

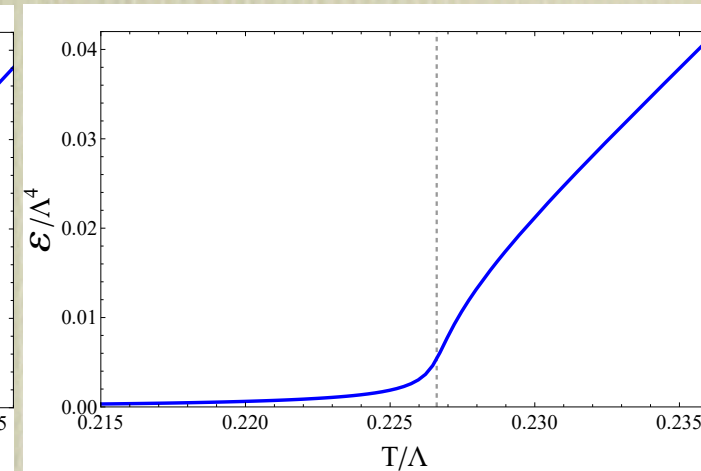
Continuous parameter \longrightarrow



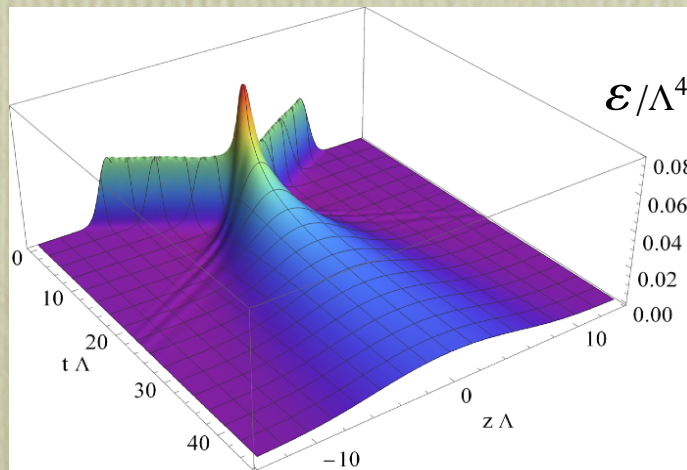
1st order



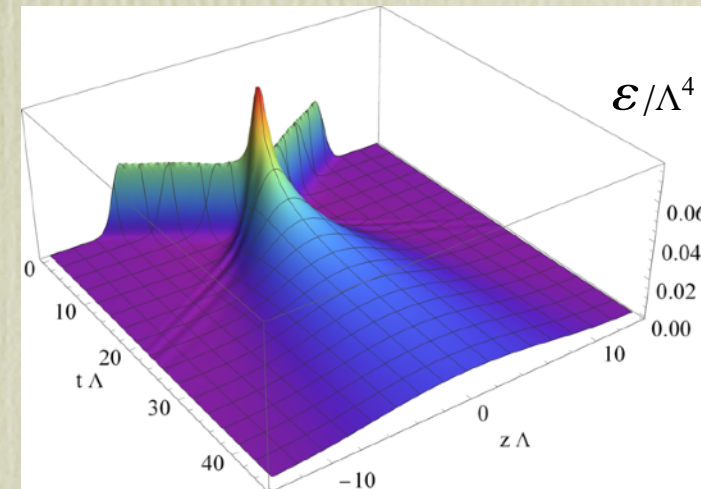
2nd order



Crossover



\longrightarrow

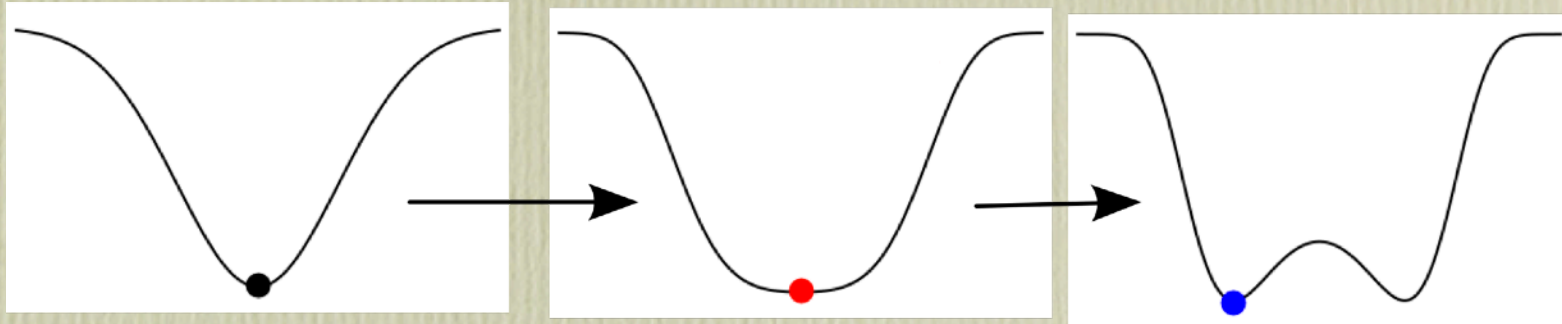


But off-equilibrium physics is qualitatively very similar

Critical fluctuations

Critical fluctuations

- Potential for order parameter flattens out at critical point:



$T > T_c$

$T = T_c$

$T < T_c$

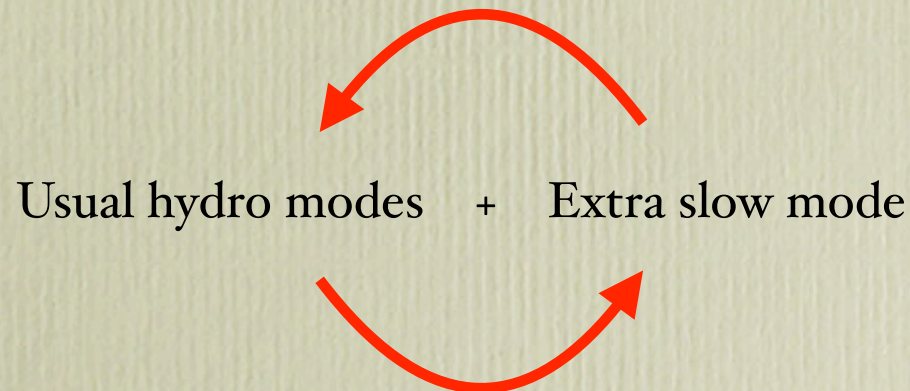
- At the critical point:

$$m \rightarrow 0, \quad \xi = m^{-1} \rightarrow \infty \quad \tau_\xi \rightarrow \infty$$

Critical fluctuations

- This leads to divergences also in transport coefficients (e.g. viscosities) because of mode-mode coupling.
- Near the critical point fluctuations of the order parameter are light and must be added to usual hydro: HYDRO+

Stephanov & Yin '17



Critical fluctuations

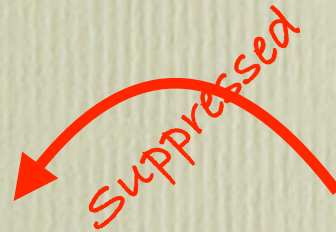
- However, at large-N:
- This leads to divergences also in transport coefficients (e.g. viscosities) because of mode-mode coupling.

Suppressed

- Near the critical point fluctuations of the order parameter are light and must be added to usual hydro: HYDRO+

Stephanov & Yin '17

Usual hydro modes + Extra slow mode



Critical fluctuations

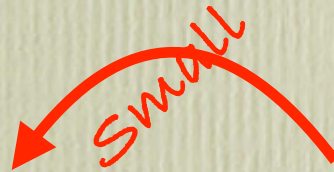
- Moreover, even at finite-N:

- Near the critical point fluctuations of the order parameter are light and must be added to usual hydro: HYDRO+

Stephanov & Yin '17

Rajagopal, Ridgway, Weller & Yin '17

Usual hydro modes + Extra slow mode



Thank you