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- Exotic Hodge theories
- The RRG theorem: two trivial cases
 - The proof when S is a point
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A Riemann-Roch theorem in Bott-Chern cohomology

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Banff, October 28th - November 1st 2019

BRIDGING THE GAP BETWEEN KÄHLER AND
NON-KÄHLER COMPLEX GEOMETRY

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Bott-Chern cohomology

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$$H_{\text{BC}}^{(p,q)}(S, \mathbf{C}) = \frac{\ker d^S \cap \Omega^{(p,q)}(S, \mathbf{C})}{\bar{\partial}^S \partial^S \Omega^{(p-1,q-1)}(S, \mathbf{C})}.$$

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- In general $H_{\text{BC}}^i(X, \mathbf{C})$ strictly finer than $H_{\text{DR}}^i(X, \mathbf{C})$.

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- In general $H_{\text{BC}}(X, \mathbf{C})$ strictly finer than $H_{\text{DR}}(X, \mathbf{C})$.

- $H_{\text{BC}}^{(=)}(S, \mathbf{R}) = \bigoplus_{0 \leq p \leq n} H_{\text{BC}}^{(p,p)}(S, \mathbf{R})$.

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Characteristic classes in $H_{\text{BC}}(S, \mathbf{R})$

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- E holomorphic vector bundle, g^E Hermitian metric.

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- Bott-Chern class $[\text{ch}(E, g^E)] \in H_{\text{BC}}^{(=)}(S, \mathbf{R})$ does not depend on g^E .
- It will be denoted $\text{ch}_{\text{BC}}(E) \in H_{\text{BC}}^{(=)}(S, \mathbf{R})$.

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A theorem of RRG

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- $p : M \rightarrow S$ proper submersion of complex manifolds, with fibre $X_s = p^{-1}(s)$.

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A theorem of RRG

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Theorem

- If Rp_*F locally free,

$$\mathrm{ch}_{\mathrm{BC}}(Rp_*F) = p_* [\mathrm{Td}_{\mathrm{BC}}(TX) \mathrm{ch}_{\mathrm{BC}}(F)] \text{ in } H_{\mathrm{BC}}^{(=)}(S, \mathbf{R}).$$

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$$\mathrm{ch}_{\mathrm{BC}}(Rp_*F) = p_* [\mathrm{Td}_{\mathrm{BC}}(TX) \mathrm{ch}_{\mathrm{BC}}(F)] \text{ in } H_{\mathrm{BC}}^{(=)}(S, \mathbf{R}).$$

- For $c_{1,\mathrm{BC}}(Rp_*F)$, the result is valid in full generality.

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Remarks

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- Families index theorem of Atiyah-Singer implies de Rham version of this result, valid even if Rp_*F not locally free.

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Remarks

- Families index theorem of Atiyah-Singer implies de Rham version of this result, valid even if Rp_*F not locally free.
- If M, S projective, the result follows from versions of Riemann-Roch-Grothendieck, even if Rp_*F not locally free.
- In general, if \mathcal{F} coherent sheaf, $\text{ch}_{\text{BC}}(\mathcal{F})$ was in principle defined by Schweitzer.

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Hodge theory without a metric

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Hodge theory without a metric

- X smooth compact oriented manifold.
- One can scale the intersection product $\int_X \alpha \wedge \beta \dots$
- ...so as to obtain a nondegenerate Hermitian form of signature (∞, ∞) .

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An elementary example

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- If $\dim_{\mathbf{R}} M = 2$, $\eta(\alpha^{(0)}, \beta^{(2)}) = -i \int_M \alpha \wedge \bar{\beta} \dots$

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An elementary example

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- $\eta(\alpha^{(1)}, \beta^{(1)}) = i \int_M \alpha \wedge \bar{\beta} \dots$
- $\eta(\alpha^{(2)}, \beta^{(0)}) = i \int_M \alpha \wedge \bar{\beta}$.
- η is a Hermitian form of signature (∞, ∞) .

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The adjoint of d

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- d^* adjoint of d with respect to η .

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- If M complex, $\bar{\partial}^* = \partial, \partial^* = \bar{\partial}$.
- $[\bar{\partial}, \bar{\partial}^*] = 0, [\partial, \partial^*] = 0$.

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The case of a Hermitian vector bundle

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The case of a Hermitian vector bundle

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- Then $\nabla^{E//*} = \nabla^{E'}$.
- Curvature R^E is the Hodge Laplacian $[\nabla^{E//}, \nabla^{E'}] \dots$
- ... which is nilpotent.

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A modified Hermitian form on $\Omega^{\cdot}(M)$

A modified Hermitian form on $\Omega^*(M)$

- Assume M complex and ω a real $(1, 1)$ form.

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A modified Hermitian form on $\Omega^1(M)$

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- $d^* = d - id\omega \wedge, \bar{d}^* = \bar{\partial} - id\omega.$
- $[d, d^*] = 0, [\bar{\partial}, \bar{d}^*] = -i\bar{\partial}d\omega.$
- Holomorphic Laplacian vanishes if and only if $\bar{\partial}d\omega = 0.$

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Description of two trivial cases

Description of two trivial cases

- First case: $M = S$, fibre is a point.

Description of two trivial cases

- First case: $M = S$, fibre is a point.
- Second case: base S is a point.

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References

The case where the fibre is a point

The case where the fibre is a point

- Take $M = S$, $F = \mathbf{C}$.

The case where the fibre is a point

- Take $M = S$, $F = \mathbf{C}$.
- The theorem to be proved is the known fact $1 = 1$.

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- X arbitrary, deformation of $D^X = \bar{\partial}^X + \bar{\partial}^{X*}$ in smooth category to classical Dirac operator (Atiyah-Singer).
- In families, smooth deformation destroys the holomorphic structure: Bott-Chern information is lost!

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A simple idea of the proof

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- How to prove RRH by heat equation while preserving $\bar{\partial}^X$ in the non-Kähler case ?

A simple idea of the proof

- How to prove RRH by heat equation while preserving $\bar{\partial}^X$ in the non-Kähler case ?
- By enlarging the set of permissible metrics.

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The local index theorem

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- For $t > 0$, $p_t(x, x')$ smooth kernel for $\exp(-tD^{X,2})$.

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- The local index theorem says that in certain cases, as $t \rightarrow 0$, $\text{Tr}_s [p_t(x, x)]$ has a geometrically computable limit...
- ... which proves the index theorem.
- This holds in particular when X is Kähler.

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- Exotic Laplacian $\bar{\partial}^X \partial^X \omega^X$ obstruction to local index theorem.

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A Lichnerowicz formula for the Bochner Laplacian

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$$\begin{aligned} \left(\bar{\partial}^X + \bar{\partial}^{X*} \right)^2 &= -\frac{1}{2} \nabla_{e_i}^{\Lambda \cdot (T^* X)} \otimes_{F,2} + \frac{K^X}{8} + \left(R^F + \frac{1}{2} \text{Tr} [R^{TX}] \right)^c \\ &\quad - \left(\bar{\partial}^X \partial^X i\omega^X \right)^c - \frac{1}{16} \left\| \left(\bar{\partial}^X - \partial^X \right) \omega^X \right\|_{\Lambda \cdot (T_{\mathbf{R}}^* X)}^2. \end{aligned}$$

A Lichnerowicz formula for the Bochner Laplacian

$$\begin{aligned} \left(\bar{\partial}^X + \bar{\partial}^{X*}\right)^2 &= -\frac{1}{2} \nabla_{e_i}^{\Lambda \cdot (T^*X)} \otimes F, 2 + \frac{K^X}{8} + \left(R^F + \frac{1}{2} \text{Tr} [R^{TX}]\right)^c \\ &\quad - \left(\bar{\partial}^X \partial^X i\omega^X\right)^c - \frac{1}{16} \left\| \left(\bar{\partial}^X - \partial^X\right) \omega^X \right\|_{\Lambda \cdot (T_{\mathbf{R}}^* X)}^2. \end{aligned}$$

The term $\left(\bar{\partial}^X \partial^X i\omega^X\right)^c$ is of length 4 in the Clifford algebra.

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The term $\left(\bar{\partial}^X \partial^X i\omega^X \right)^c$ is of length 4 in the Clifford algebra. Local index theory accepts only terms of length ≤ 2 .

If $\bar{\partial}^X \partial^X \omega^X = 0$, there is a local index theorem, compatible with RRH.

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The space \mathcal{X}

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- $\pi : \mathcal{X} \rightarrow X$ total space of TX , with fibre \widehat{TX} , $\hat{y} \in \widehat{TX}$ tautological section, $y \in TX$ corresponding section of TX .

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- The cohomology of this new complex is still equal to $H^{(0,\cdot)}(X, F)$.

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Exotic Hodge theory

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- ϵ Hermitian form of signature (∞, ∞) .

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Evaluation of the adjoint

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- It is potentially good in families: it has been obtained by replacing L_2 metric by nonpositive metric.

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- It is still true that $\chi(X, F) = \text{Tr}_s [\exp(-tL_b)]$.

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- Except when $\bar{\partial}^X \partial^X \omega^X = 0$, no local index theorem for the heat kernel for L_b .
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- $|Y|_{g^{\widehat{TX}}}^2$ critical power.

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- Analytic torsion forms $\frac{\bar{\partial}^S \partial^S}{2i\pi} T = \alpha_\infty - \alpha_0 = \text{ch} (R p_* F, g^{R p_* F}) - p_* \left[\text{Td} (TX, g^{TX}) \text{ch} (F, g^F) \right]$.

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- These results extend to $\bar{\partial}^M \partial^M \omega^M = 0$.

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- This is possible because we deform nondegenerate Hermitian forms.

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- Make $t \rightarrow 0 \dots$ finally!

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- ...and explain the given proof in the general case.

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- The term $\bar{\partial}^S \partial^S \omega^S$ appears ‘because’ it is a Laplacian in the exotic Hodge theory of S .

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- The family of fiberwise Hodge Laplacians is 0 acting on \mathbf{C} !



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