Long Transients in Ecology: Theory and Observations

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Plan of the talk

- Introduction: what are long transients?
- Basic mechanisms generating long transients (nonspatial systems)
- Relation to tipping points
- A (brief) look at spatial systems
- Conclusions

What is it all about

Transient: lasting for only a short time; temporary

(Cambridge English Dictionary)

Typically, transients are associated with the effect of the initial conditions and disappear relatively fast.

Long-term dynamics are usually associated with the system's attractors.

"Long transient" is apparently an oxymoron??

However...

Dynamics of a nonspatial, time-discrete, single-species model:



⁽from Schreiber, 2003)

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Time-continuous single-species model with time-delay:



(from Morozov et al., 2016)

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Space-time-continuous, 3-species model (plankton dynamics):



(from Petrovskii et al., 2017)

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Empirical examples are abundant too



(from Cushing et al., 1998)

Forage fishes (field) (from Frank et al., 2011)

Empirical examples are abundant too



Flour beetle data (lab) (from Cushing et al., 1998) Forage fishes (field) (from Frank et al., 2011) In all above examples, a regime shift occurs

A well-known theory of regime shifts relates it to a tipping point: a bifurcation (e.g. saddle-node) due to a slow change in some system's parameter (environmental conditions) (e.g. Scheffer et al. 2009, 2012; Kuehn 2011; Dakos et al., 2012, 2014)

Interestingly, in all above examples, parameters (environmental conditions) are constant!

How that can be possible?

Overview of the baseline mechanisms

"Crawl-by": transients induced by a saddle



Here A is the 'small' vicinity of the saddle, B the range of appropriate initial conditions

Transients induced by a saddle

Consider a generic population dynamics model:

$$\frac{du_k(t)}{dt} = f_k(\mathbf{u}), \qquad k = 1, \dots, n,$$

where $\mathbf{u} = (u_1, \dots, u_n)$ are the population densities, *t* is time.

Linearized system in the vicinity of a steady state $\bar{\boldsymbol{u}}$:

$$\frac{dx_k(t)}{dt} = a_{k1}x_1 + \ldots a_{kn}x_n, \qquad k = 1, \ldots, n,$$

where $x_k(t) = u_k(t) - \overline{u}_k$.

Solution is a linear combination of exponents $e^{\lambda_i t}$. Let λ_1 be the eigenvalue with the largest real part, $\text{Re}\lambda_1 > 0$. The time spent in the vicinity of the (unstable) steady state is estimated as

$$au \propto rac{1}{{
m Re}\lambda_1}.$$

Transients induced by a saddle

Nonlinear effects can substantially increase the range of appropriate initial conditions:



A is the 'small' vicinity, B the range of appropriate initial conditions, S is a separatrix

Example: Rosenzweig–MacArthur model



This will result in recurrent long transients:



Generalization 1

A modified prey-predator system can have a saddle point in the interior of the domain (not at the origin), so that the decay to low density is not a necessary property

Example: strong Allee effect for prey, quadratic mortality for predator



(Sen & Banerjee 2015)

Generalization 2

Saddle-induced transients in a higher-dimensional systems

A case of more complex dynamics: connected saddles:



(Ashwin & Timme, 2005)

Consider a generic two-species system:



Two-species nonlinear competition model (Hastings et al. 2018)

A change in the parameter value can bring the system beyond the saddle-node bifurcation:



However, the local bifurcation does not change the global structure of the phase flow: the system slows down in the vicinity of the pre-bifurcation steady state location

The long transient dynamics occur:



The transient's duration depends on the closeness to the bifurcation:

$$au \propto \left| oldsymbol{p} - oldsymbol{p}_{oldsymbol{c}}
ight|^{-0.5}$$
 .

A similar mechanism applies to more complicated dynamics, e.g. periodic solutions (limit cycles) and chaos.

Example: long-term chaotic transient (chaotic ghost) in a resource-consumer-predator system(Hastings and Powell 1991; McCann and Yodzis 1994)



Pre-bifurcation: chaotic attractor coexists with a stable limit cycle



Post-bifurcation: the two basins merge, chaotic attractor disappears

Chaotic transients can be particularly long: $\tau \propto \exp\left(k|p-p_c|^{-\gamma}\right)$ $(k,\gamma>0)$

(Grebogi et al. 1983, 1985)

Example of the time-series generated by a chaotic ghost:



(Petrovskii et al., 2017)

Slow-fast systems

Consider

$$\frac{du(t)}{dt} = f(u, v, \epsilon), \qquad \frac{dv(t)}{dt} = \epsilon g(u, v, \epsilon), \qquad \epsilon \ll 1.$$
(1)

Introducing a rescaled time $\tau = \epsilon t$, it turns into

$$\epsilon \frac{du(\tau)}{d\tau} = f(u, v, \epsilon), \qquad \frac{dv(\tau)}{d\tau} = g(u, v, \epsilon).$$
⁽²⁾

In the limit $\epsilon \rightarrow$ 0, system (1) turns into

$$\frac{du(t)}{dt} = f(u, v, 0), \qquad \frac{dv(t)}{dt} = 0,$$

and system (2) turns into

$$0=f(u,v,0),\qquad \frac{dv(\tau)}{d\tau}=g(u,v,0),$$

Slow-fast systems

Example 1: periodical dynamics in a prey-predator system ($\epsilon = 0.01$)



Slow-fast systems

Example 2: aperiodical dynamics in a two-species competition system



Black (dashed) curve for $\epsilon = 1$, red curve for $\epsilon = 0.002$

Relation between long transients and tipping points

Relation between long transients and tipping points



Parameter change very slow or with limited variation: regime shift after LT ghost dynamics

Long transients in higher dimensional systems

- Effect of time-delay is known to generate long transients but the scaling law is unknown
- Effect of noise broad and variable. For non-chaotic systems (saddles and ghosts), tends to decrease the transient's life-time but would not normally destroy it. Can create the transient dynamics (e.g. in bistable systems):



- For chaotic transients, noise can increase as well as decrease the transient's life-time (Grebogi et al. 1983; Do and Lai 2004, 2005).
- Spatial systems: new types of transients (e.g. related to population waves propagation).

What are the new phenomena brought in by explicit space?

- Pattern formation
- Synchronization / desynchronization & onset of spatiotemporal chaos
- Travelling waves

Consider the space-continuous, time-discrete single-species system:

$$u(x,t+1) = \int_0^L g(x-y)F(u(x,t))dx, \quad F(u) = ue^{r(1-u)}.$$

For distributed random initial conditions, the system's dynamics exhibit a chaotic saddle:



(Hastings and Higgins, 1994)

The above system exhibits long transients in terms of the spatially average values

Knowledge of the spatial population distribution can provide a different angle on long transients

Example: "wave of chaos" in a space-time-continuous prey-predator system:



⁽Petrovskii and Malchow, 2001)

Spread of the chaotic phase over the system can take a very long time, $\tau \propto \frac{l}{c}$.

For compact initial conditions, the system's dynamics usually consists of a succession of population waves

Example: space-time-continuous (diffusion-reaction) prey-predator system, invasion of predator; dynamical stabilization in the wake of the invasion front



(Petrovskii and Malchow, 2000)

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(Petrovskii and Malchow, 2000)

(Petrovskii et al., 2016)

2.5

1.5

0.5

Conclusions

- Long transients do occur
- The life-time of long transients can be arbitrary long (cf. scaling laws)
- We have identified a few basic mechanisms for the long transients to occur
- Long transients provide an alternative scenario of regime shifts

References

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Thanks for listening