## Thermality of circular motion

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Biermann et al. PRD 102, 085006 (2020) [arXiv:2007.09523]


## Plan

1. Unruh effect

- Relativistic spacetime and analogue spacetime

2. "Quantum dot"

- Unruh-DeWitt

3. Circular motion

- Wightman function: $3+1$ and $2+1$

4. Results

- Ratio $T_{\text {circular }} / T_{\text {linear }}$

5. Summary and outlook

## 1. Unruh effect

Well established

- Uniformly linearly accelerated observer sees Minkowki vacuum as thermal, $T=\frac{a}{2 \pi}$

Unruh 1976

- Weak coupling, long time, negligible switching effects
- Thermal: Observer/detector records detailed balance:

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\frac{P_{\downarrow}}{P_{\uparrow}}=e^{E_{\mathrm{gap}} / T}
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## Uniform circular motion?

- Long time in finite size lab!
- Accelerator storage rings

Bell and Leinaas 1983,...

- Analogue spacetime: BEC, ${ }^{4} \mathrm{He}, \ldots$. Weinfurtner talk (Friday)


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## Sense of "temperature" ?

## Aims

Why now

# What today 

What not today

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- Analogue spacetime experiment proposal Gooding et al. 2020
- Finite size lab
- Time dilation $\leftrightarrow$ time-independent energy scaling

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What today


- "Quantum dot"
- Weak coupling, long time, negligible switching effects
- Sense of temperature
- Relativistic spacetime versus analogue spacetime
- $3+1$ versus $2+1$

What not today

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- Analogue spacetime experiment proposal Gooding et al. 2020
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- $3+1$ versus $2+1$

What not today

- "Quantum dot" $\rightarrow$ actual experiment?


## 2. "Quantum dot" (relativistic) Unruh(1976)-DeWitt(1979)

Quantum field
D spacetime dimension
$\phi \quad$ real scalar field
|0〉 Minkowski vacuum

Two-state detector (atom)
$\| 0\rangle$ state with energy 0
$\| 1\rangle\rangle \quad$ state with energy $E$
$\times(\tau)$ detector worldline, $\tau$ proper time

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Interaction

$$
H_{\text {int }}(\tau)=c \chi(\tau) \mu(\tau) \phi(\times(\tau))
$$

c coupling constant
$\chi$ switching function, $C_{0}^{\infty}$, real-valued
$\mu$ detector's monopole moment operator

## Probability of transition

$$
\| 0\rangle\rangle \otimes|0\rangle \longrightarrow \| 1\rangle\rangle \otimes \mid \text { anything }\rangle
$$

in first-order perturbation theory:

$$
\begin{gathered}
P(E)=c^{2} \underbrace{}_{\left.\begin{array}{c}
\text { detector internals only: } \\
\text { drop! }
\end{array} \right\rvert\,\left\langle\left.\langle 0\|\mu(0)\| 1\rangle\right|^{2}\right.} \times \underbrace{F_{\chi}(E)}_{\begin{array}{c}
\text { trajectory and }|0\rangle: \\
\text { response function }
\end{array}} \\
F_{\chi}(E)=\int \mathrm{d} \tau^{\prime} \mathrm{d} \tau^{\prime \prime} \mathrm{e}^{-i E\left(\tau^{\prime}-\tau^{\prime \prime}\right)} \chi\left(\tau^{\prime}\right) \chi\left(\tau^{\prime \prime}\right) W\left(\tau^{\prime}, \tau^{\prime \prime}\right) \\
W\left(\tau^{\prime}, \tau^{\prime \prime}\right)=\langle 0| \phi\left(\times\left(\tau^{\prime}\right)\right) \phi\left(\times\left(\tau^{\prime \prime}\right)\right)|0\rangle \quad \begin{array}{c}
\text { Wightman function } \\
\text { (distribution) }
\end{array}
\end{gathered}
$$

- Stationary motion:

$$
W\left(\tau^{\prime}, \tau^{\prime \prime}\right)=W\left(\tau^{\prime}-\tau^{\prime \prime}, 0\right)
$$

- Transition rate in the long time limit:

$$
\begin{gathered}
\frac{F_{\chi}(E)}{\Delta \tau} \underset{\Delta \tau \rightarrow \infty}{ } \\
F(E)=\int_{-\infty}^{\infty} \mathrm{d} s \mathrm{e}^{-i E s} W(s, 0) \\
\text { stationary response function }
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- Temperature via detailed balance:

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T=\frac{E}{\ln \left(\frac{F(-E)}{F(E)}\right)}
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- For uniform linear acceleration, $T_{\text {lin }}=\frac{a}{2 \pi} \quad$ Unruh 1976 (genuine KMS state)
- For other uniform motions, $T$ depends also on $E$

Letaw 1981,... . Good et al. 2020

## 3. Circular motion

- Metric: $\quad d s^{2}=-d t^{2}+\left(d x^{1}\right)^{2}+\cdots+\left(d x^{D-1}\right)^{2}$
- Trajectory: $\quad \times(\tau)=(\gamma \tau, R \cos (\gamma \Omega \tau), R \sin (\gamma \Omega \tau), \cdots)$

$$
R>0 \text { radius, } 0<R \Omega<1 \text { orbital velocity, } \gamma=1 / \sqrt{1-(R \Omega)^{2}}
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- Wightman: (massless field; $\epsilon \rightarrow 0^{+}$)

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\begin{aligned}
& D=3+1: W(s, 0)=\frac{1}{4 \pi^{2}[\times(s-i \epsilon)-\times(0)]^{2}} \\
& D=2+1: W(s, 0)=\frac{1}{4 \pi \sqrt{[x(s-i \epsilon)-\times(0)]^{2}}}
\end{aligned}
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Now examine:

- Relativistic spacetime: $\frac{T_{\text {circ }}}{T_{\text {lin }}}$ (for same proper acceleration)
- Analogue spacetime: similarly


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$D=3+1: \quad W(s, 0)=\frac{1}{4 \pi^{2}[\times(s-i \epsilon)-\times(0)]^{2}} \quad$ real for $s \neq 0$
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## 4a. Relativistic spacetime

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$3+1$ dimensions

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Not constant!

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$E_{\text {red }} \rightarrow 0$ : nonzero

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$v \rightarrow 1$ :


## 4b. Analogue spacetime $T_{\text {rlab }}:=\hat{T}_{\text {circ }} / \hat{T}_{\text {lin }}, \quad E_{\text {rab }}:=\hat{E} / \hat{a}, \quad \wedge=$ non-rel

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$$
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Suppressed at $v \rightarrow 1$

## 5. Summary and outlook

## Setting

- "Quantum dot" in circular motion; massless scalar
- Relativistic and analogue
- Unruh temperature $T_{\text {circ }}$ via detailed balance


## Outcomes

- $3+1: T_{\text {circ }} / T_{\text {lin }}$ of order unity (relativistic and analogue)
$-2+1: T_{\text {circ }} / T_{\text {lin }} \ll 1$ for:
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- near-sonic limit (analogue only)


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Fun fact: $2+1$ time-time correlations are purely imaginary (!?!)

