Explicit asymptotic expansions of supercuspidal characters II

A good day for asymptotic expansions; or, 2 Asymptotic 2 Expansion

Loren Spice

2021-11-16

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The machinery that is developed in S 2021 to handle the harmonic analysis described here does not assume that **G** is connected * (so "**G** reductive" really means "**G**° reductive"), but so far we need connectedness for *applications*

^{*} Blame Jeff.

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Maybe characters are *no more* non-elementary than Fourier transforms of orbital integrals

Harish-Chandra-Howe local character expansion:

$$\Phi_{\pi}(\gamma \cdot \exp(Y)) = \sum_{\mathcal{O} \in \mathcal{O}^{H}(0)} c_{\mathcal{O}}(\pi, \gamma) \hat{\mu}_{\mathcal{O}}^{H}(Y),$$

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How elementary are the coefficients $c_{\mathcal{O}}(\pi, \gamma)$? Long-term project of Gordon–Hales–S

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For example, $\Gamma_{\pi}=0$ generalises the Adler–Korman results on the local character expansion

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So we've turned the question from "compute $\Phi_{\pi}(\gamma \cdot \exp(\gamma))$ " to "compute $c_{\mathcal{O}}(\pi, \gamma)$ "

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- a description or parameterisation of elements of G, and then the formula will be written in terms of the parameters for γ

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Adler and I generalised this notion to general p-adic groups: in the connected case, replace a commuting sum of monomials by a commuting product of good elements

Given $\gamma = \prod_{i \ge 0} \gamma_i$, we put $\gamma_{< r} = \prod_{0 \le i < r} \gamma_i$, $\gamma_{\ge r}$, $Y_{\ge r} = \log(\gamma_{\ge r})$, ...

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$$\mathsf{Cent}^{(< r)}_{\mathbf{G}}(\gamma)^\circ := \mathsf{Cent}_{\mathbf{G}}(\gamma_i : 0 \leq i < r)^\circ$$

is well defined for every r

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This decomposition is needed for fine control over the terms appearing in Frobenius-type formulæ for characters of compactly induced representations, such as supercuspidals

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Other results, such as Moy's description of supercuspidal representations of $GSp_4(k)$ and U(2,1)(k), suggested a general phenomenon

Systematically investigated by Adler, and then further generalised by Yu

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Yu takes the factorisation to be the basic object: a Yu datum is a triple $(\vec{\mathbf{G}}, \pi_{-1}, \vec{\phi})$

Hakim–Murnaghan and Kaletha bring us back to the original perspective by showing that a Yu datum can be replaced by just $(\mathbf{G}^0, \pi_{-1} \prod_{i=0}^{\ell} \phi_i)$

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Theorem (S 2018)

Let r be the depth of π . Then the coefficients $c_{\mathcal{O}}(\pi, \gamma_{< r})$ of a depth-r asymptotic expansion for Φ_{π} near $\gamma_{< r}$ of G' can be computed explicitly in terms of $c_{\mathcal{O}'}(\pi'^g, \gamma_{< r})$.

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- "Explicit" is allowed to involve some quite elaborate ingredients, a full understanding of which requires a different linearisation of the Weil representation from the one used by Yu (see Jessica's talk)
- This recipe is not suitable for induction, because the depth of π' is usually strictly less than the depth of π

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- Expand each Fourier transform of an orbital integral on $C_{G'}(\gamma_0)^\circ$ near $Y_{< r}$ in terms of Fourier transforms of orbital integrals on $C_{G'}(\gamma_{< r})^\circ$, evaluated at $Y_{\geq r}$

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- Combine to obtain an expansion of the character of π' near $\gamma_{< r}$ in terms of Fourier transforms of orbital integrals on $C_{G'}(\gamma_{< r})^{\circ}$

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- Combine to obtain an expansion of the character of π' near $\gamma_{< r}$ in terms of Fourier transforms of orbital integrals on $C_{G'}(\gamma_{< r})^{\circ}$
- Use S 2018 to expand the character of π near γ_{<r} in terms of Fourier transforms of orbital integrals on C_G(γ_{<r})[°]

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- Use S 2018 to expand the character of π near γ_{<r} in terms of Fourier transforms of orbital integrals on C_G(γ_{<r})[°]
- (9) "Collapse" the expansion of the character of π near $\gamma_{< r}$ into an expansion near γ_0

The tricky part is the "collapse" step 6, which requires a delicate understanding of the relationship among the various expansions Making this relationship precise is notationally awkward, but here goes: The tricky part is the "collapse" step 6, which requires a delicate understanding of the relationship among the various expansions

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$$\begin{split} \Phi_{\pi'^{g}}(\gamma_{< r} \cdot \exp(Y'_{\geq r})) \quad equals \\ \sum_{\mathcal{O}' \in \mathcal{O}^{H'^{\circ}_{g}}(\mathcal{U}'_{g}^{*})} \widetilde{\mathfrak{G}}_{G'_{g}/H'_{g}}(\mathcal{O}', \gamma_{< r}) c_{\mathcal{O}'}(\pi'^{g}, \gamma_{< r}) \widehat{O}_{\mathcal{O}'}^{H'^{\circ}_{g}}(Y'_{\geq r}), \end{split}$$

then

$$\Phi_{\pi}(\gamma_{< r} \cdot \exp(Y_{\geq r})) \quad equals$$

$$\sum_{\substack{g \in G' \setminus G/H^{\circ} \\ \Gamma_{\pi,g} \in \text{Lie}^{*}(H)}} \sum_{\mathcal{O'}} \widetilde{\mathfrak{G}}_{G/H}(\mathcal{O'}, \gamma_{< r}) c_{\mathcal{O'}}(\pi'^{g}, \gamma_{< r}) \widehat{O}_{\mathcal{O'}}^{H^{\circ}}(Y_{\geq r}).$$

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Theorem (S 2021) If

$$\begin{split} & \widehat{O}_{\xi_g}^{G_g^0}(\gamma_{\leq 0} + Y_{>0}^0) \quad \text{equals} \\ & \sum_{\substack{g^0 \in G_g^0 \setminus G_g^0 / J_g^0 \,^{\circ} \, \mathcal{O}^0 \in \mathcal{O}^{J_g^0 \,^{\circ}}(\Gamma_{\xi,gg^0})}} \sum_{\mathfrak{G}_g^0 / J_g^0} (\mathcal{O}^0, \gamma_{\leq 0}) c_{\mathcal{O}^0}(\xi_{gg^0}, \gamma_{\leq 0}) \widehat{O}_{\mathcal{O}^0}^{J_g^0 \,^{\circ}}(Y_{>0}^0), \end{split}$$

then

$$\widehat{O}_{\xi}^{G}(\gamma_{\leq 0} + Y_{> 0}) \quad equals \\ \sum_{\substack{g \in G^{0} \setminus G/J^{\circ} \\ \Gamma_{\xi,g} \in \operatorname{Lie}^{*}(J)}} \sum_{\mathcal{O}^{0}} \mathfrak{G}_{G/J}(\mathcal{O}^{0}, \gamma_{\leq 0}) c_{\mathcal{O}^{0}}(\xi_{g}, \gamma_{\leq 0}) \widehat{O}_{\mathcal{O}^{0}}^{J^{\circ}}(Y_{> 0}).$$

$$\widehat{O}_{\xi}^{G}(\gamma) = \sum_{\substack{g \in G^{0} \setminus G/J^{\circ} \\ \Gamma_{\xi,g} \in \mathsf{Lie}^{*}(J)}} \sum_{\mathcal{O}^{0}} \mathfrak{G}_{G/J}(\mathcal{O}^{0}, \gamma_{\leq 0}) c_{\mathcal{O}^{0}}(\xi_{g}, \gamma_{\leq 0}) \widehat{O}_{\mathcal{O}^{0}}^{J^{\circ}}(Y_{>0})$$

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 \mathfrak{G} is a Gauss sum; \mathfrak{G} is essentially the same (upon using the exponential map to move between Lie(G)_{>0} and G_{>0}), but also has a contribution from the modified Weil representation that Fintzen–Kaletha–S use in Yu's construction, which depends on the topologically semisimple part γ_0 (see Jessica's talk)

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This causes no problems, since, in the "re-expansion" step 3, we have already passed to the topologically unipotent element $\gamma_{>0} = \exp(\gamma_{>0})$, for which the two quantities agree

Theorem (S 2021)

lf

$$\begin{split} \Phi_{\pi_0^g}(\gamma_0 \cdot \exp(Y_{>0}^0)) \quad equals \\ \sum_{\substack{g^0 \in G_g^0 \setminus G_g^0 / J_g^0 \\ \Gamma_{gg^0} \in \text{Lie}^*(J_g^0)}} \sum_{\mathcal{O}^0 \in \mathcal{O}^{J_g^0 \circ}(\Gamma_{\pi,gg^0})} \widetilde{\mathfrak{G}}_{G_g^0 / J_g^0}(\mathcal{O}^0,\gamma_0) c_{\mathcal{O}^0}(\pi_0^g,\gamma_0) \widehat{O}_{\mathcal{O}^0}^{J_g^0 \circ}(Y_{>0}^0), \end{split}$$

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If we are dealing with semisimple orbital integrals, then we may also recover a quantitative version of a result of Waldspurger:

Corollary (Waldspurger 1995 and S 2021)
If
$$\gamma_{\leq 0} \in \text{Lie}(G)$$
 is regular semisimple, then
 $\widehat{O}_{\xi}^{G}(\gamma_{\leq 0} + Y_{>0})$ equals $\sum_{\substack{g \in G^{0} \setminus G/J^{0} \\ \Gamma_{\xi,g} \in \text{Lie}^{*}(J)}} \mathfrak{G}_{G/G_{g}^{0}}(\Gamma_{\xi,g}, \gamma_{\leq 0}) \Lambda_{\xi_{g}}(\gamma_{\leq 0} + Y_{>0}).$

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They specialise to a result of Kaletha, which I regard as the group analogue of Waldspurger's result on Fourier transforms of orbital integrals:

If γ_0 is regular semisimple, then

$$\Phi_{\pi}(\gamma_{0} \cdot \exp(Y)) \quad equals \sum_{\substack{g \in G^{0} \setminus G/J^{0} \\ \Gamma_{\pi,g} \in \operatorname{Lie}^{*}(J)}} \widetilde{\mathfrak{G}}_{G/G_{g}^{0}}(\Gamma_{\pi,g},\gamma) \Phi_{\pi_{0}}(\gamma) \Lambda_{\Gamma_{\pi,g}}(Y).$$