# Fractional powers of first order differential operators and inverse measures

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#### Theoretical and Applied Aspects for Nonlocal Models Banff International Research Station

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 Joint work with Martín Mazzitelli (Instituto Balseiro, Argentina) and José L. Torrea (Universidad Autónoma de Madrid), arXiv 2022.

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$$Lu = -\frac{d^2u}{dx^2} + 2x\frac{du}{dx}$$

is symmetric with respect to  $d\gamma_1$ :  $\int_{\mathbb{R}} (Lu) v \, d\gamma_1 = \int_{\mathbb{R}} u(Lv) \, d\gamma_1$ 

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$$H_n(x) = (-1)^n e^{x^2} \frac{d^n}{dx^n} (e^{-x^2}) \qquad n \ge 0$$

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- Hermite polynomials form an orthogonal basis of  $L^2(\mathbb{R}, d\gamma_1)$
- Gaussian harmonic analysis: dγ<sub>1</sub> is nondoubling and non-Ahlfors regular, so classical CZ theory on metric measure spaces does not directly apply.

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- Orthonormal basis of  $L^2(\mathbb{R}, d\gamma_{-1})$ ? Not polynomials anymore!
- Inverse Gaussian harmonic analysis: similar obstructions.

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A hypersurface  $\Sigma$  in  $\mathbb{R}^{n+1}$  is a self-expander if

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 $H_{\Sigma}$ : mean curvature vector on  $\Sigma$ ;  $x^{\perp}$ : orthogonal projection of x onto the normal.

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Self-expanders describe the asymptotic behavior and the local structure of mean curvature flow after singularities for short times.

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Self-expanders describe the asymptotic behavior and the local structure of mean curvature flow after singularities for short times.

A self-expander M is a critical point of the weighted volume functional

$$F(M) = \int_M e^{|\mathbf{x}|^2/4} \, d\mathcal{H}^{n-1}$$

Reference. T. Ilmanen, Lectures on mean curvature flow and related equations (1995)

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#### Fractional derivatives and integrals

For  $u: \mathbb{R} \to \mathbb{R}$ , let

$$D_{left}u(x) = \lim_{t \to 0^+} \frac{u(x) - u(x-t)}{t} = -\lim_{t \to 0^+} \frac{e^{-tD_{left}}u(x) - u(x)}{t}$$

where the left translation semigroup is  $e^{-tD_{left}}u(x) = u(x - t)$ .

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where the left translation semigroup is  $e^{-tD_{left}}u(x) = u(x - t)$ . For  $0 < \alpha < 1$ ,

$$(D_{left})^{\alpha}u(x) = \frac{1}{\Gamma(-\alpha)} \int_0^{\infty} \left(e^{-tD_{left}}u(x) - u(x)\right) \frac{dt}{t^{1+\alpha}}$$
$$= c_{\alpha} \int_{-\infty}^{x} \frac{u(x) - u(t)}{(x-t)^{1+\alpha}} dt$$

is the Marchaud fractional derivative;

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$$(D_{left})^{-\alpha}u(x) = \frac{1}{\Gamma(\alpha)} \int_0^\infty e^{-tD_{left}}u(x) \frac{dt}{t^{1-\alpha}}$$
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is the Weyl fractional integral.

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#### References. Bernardis-Martín-Reyes-S.-Torrea (2016); S.-Vaughan (2020).

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• Using

$$\int_{\mathbb{R}} ((D_{left})^{\alpha} u) v \, dx = \int_{\mathbb{R}} u((D_{right})^{\alpha} v) \, dx$$

and **one-sided test functions**, we can define  $(D_{left})^{\alpha}u$  in the sense of distributions.

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• Maximum principle. If  $(D_{left})^{\alpha} u \leq 0$  in (0, T] and  $u \leq 0$  in  $(-\infty, 0]$  then  $\sup_{(-\infty, T]} u = \sup_{(-\infty, 0]} u$ .

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- Extension problem. If U = U(x, y) solves

$$\begin{cases} -D_{left}U + \frac{1-2\alpha}{y}U_y + U_{yy} = 0 & \text{for } x \in \mathbb{R}, \ y > 0\\ U(x,0) = u(x) & \text{on } \mathbb{R} \end{cases}$$

then  $-d_{\alpha}y^{1-2\alpha}U_y(x,y)\big|_{y=0} = (D_{left})^{\alpha}u(x).$ 

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then  $-d_{\alpha}y^{1-2\alpha}U_y(x,y)\big|_{y=0} = (D_{left})^{\alpha}u(x).$ 

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• Harnack inequality. If  $u \ge 0$  in  $\mathbb{R}$ , and  $(D_{left})^{\alpha}u = 0$  in (0,1), then

$$\sup_{(1/4,1/3)} u \le C \inf_{(1/2,3/4)} u$$

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References. Bernardis-Martín-Reyes-S.-Torrea (2016); S.-Vaughan (2020).

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• One-sided fractional Sobolev spaces. For  $1 \le p < \infty$  and  $\omega \in A_p^-(\mathbb{R})$  (one-sided Sawyer weight)

$$W^{lpha, p}(\omega^p) = \{ u = (D_{left})^{-lpha} f : f \in L^p(\omega^p) \}$$

Characterizations with left fractional derivatives.

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Characterizations with left fractional derivatives.

• Bourgain-Brezis-Mironescu-type result. For

$$W^{1,p}(\omega) = \{ u \in L^p(\omega) : D_{left} u \in L^p(\omega) \}$$

we have

$$\begin{split} \lim_{\alpha \to 1} (D_{\mathit{left}})^{\alpha} u &= D_{\mathit{left}} u \qquad \text{in } L^p(\omega), \ 1$$

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• Fundamental Theorem of Fractional Calculus.

$$u(x) = \lim_{\varepsilon \to 0} (D_{left})^{\alpha}_{\varepsilon} (D_{left})^{-\alpha} u(x) \quad \text{in } L^{p}(\omega^{p}) \text{ and a.e.}$$

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For a(x) continuous, let

$$\mathfrak{D}_{left,a}u(x) = D_{left}u + a(x)u$$

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Fix  $x_0 \in \mathbb{R}$  and define

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It can be checked that

$$\mathfrak{D}_{left,a}u(x)=\mathcal{E}(x)D_{left}(\mathcal{E}^{-1}u)(x)$$

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#### Lemma (Semigroup)

$$e^{-t\mathfrak{D}_{left,a}}u(x) = \mathcal{E}(x)e^{-tD_{left}}(\mathcal{E}^{-1}u)(x)$$

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## Fractional powers of first order differential operators

**Positive fractional power** is a *modulation* of Marchaud:

$$\begin{split} (\mathfrak{D}_{left,a})^{\alpha} u(x) &= \frac{1}{\Gamma(-\alpha)} \int_0^{\infty} \left( e^{-t\mathfrak{D}_{left,a}} u(x) - u(x) \right) \frac{dt}{t^{1+\alpha}} \\ &= c_{\alpha} \mathcal{E}(x) \int_{-\infty}^{x} \frac{(\mathcal{E}^{-1}u)(x) - (\mathcal{E}^{-1}u)(t)}{(x-t)^{1+\alpha}} dt \\ &= \mathcal{E}(x) (D_{left})^{\alpha} (\mathcal{E}^{-1}u)(x) \end{split}$$

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Negative fractional power is a *modulation* of Weyl:

$$(\mathfrak{D}_{left,a})^{-\alpha}u(x) = \frac{1}{\Gamma(\alpha)} \int_0^\infty e^{-t\mathfrak{D}_{left,a}}u(x) \frac{dt}{t^{1-\alpha}}$$
$$= c_{-\alpha}\mathcal{E}(x) \int_{-\infty}^x \frac{(\mathcal{E}^{-1}u)(t)}{(x-t)^{1-\alpha}} dt$$
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# Fractional powers of first order differential operators

Positive fractional power is a *modulation* of Marchaud:

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► We can transfer results from (D<sub>left</sub>)<sup>±α</sup> to (D<sub>left,a</sub>)<sup>±α</sup> using modulation (semigroup, distributions, maximum principle, extension, BBM-type results, etc)

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Recall that, for  $\omega \in A^-_p(\mathbb{R})$  and  $a \equiv 0$ ,

$$W^{1,p}(\omega) = \left\{ u \in L^p(\omega) : D_{left} u \in L^p(\omega) \right\}$$

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Since, for any a(x),

$$\mathcal{E}(x)^{-1}\mathfrak{D}_{left,a}u = D_{left}(\mathcal{E}^{-1}u)(x)$$

then the correct definition of Sobolev space associated with  $\mathfrak{D}_{\mathit{left},a}$  is

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Similarly,

$$W^{\alpha,p}_{a}(\omega^{p}) = \left\{ u = (\mathfrak{D}_{left,a})^{-\alpha}f : f \in L^{p}(\mathcal{E}^{-p}\omega^{p}) \right\}$$

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**Conclusion.** The natural  $L^{p}$  space for analysis of  $(\mathfrak{D}_{left,a})^{\pm \alpha}$  is

$$L^{p}(\mathcal{E}^{-p}) = \{ u : \mathbb{R} \to \mathbb{R} : \mathcal{E}^{-1}u \in L^{p}(\mathbb{R}) \}$$

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### An important example

When a(x) = x,

$$\mathfrak{D}_{left,a}u = D_{left}u + xu$$

• This is the natural *derivative* in harmonic analysis of Hermite expansions:

$$-\frac{d^2}{dx^2} + x\frac{d}{dx} = (D_{left} + x)(D_{right})$$

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By taking  $x_0 = 0$ ,  $\mathcal{E}(x)^{-1} = \exp\left(\int_0^x a(y) \, dy\right) = e^{x^2/2}$ 

and we end up with the inverse Gaussian space

$$L^2(\mathcal{E}^{-2}) = \{u: \mathbb{R} \to \mathbb{R}: e^{x^2/2}u \in L^2(\mathbb{R})\} = L^2(\mathbb{R}, d\gamma_{-1})$$

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### Theorem (Mazzitelli–S.–Torrea, 2022)

The polynomials given by the Rodrigues formula

$$\mathcal{H}(x) = (-1)^n e^{-x^2} \frac{d^n}{dx^n} (e^{x^2}) \qquad n \ge 0$$

are eigenfunctions of  $\mathcal{L} = \frac{d^2}{dx^2} + 2x \frac{d}{dx}$  with  $\mathcal{LH}_n = 2n\mathcal{H}_n$ . Generating function formula:  $e^{t^2 - 2xt} = \sum_{n=0}^{\infty} \mathcal{H}_n(x) \frac{t^n}{n!}$ Three-term recurrence relation:  $\mathcal{H}_{n+1}(x) + 2x\mathcal{H}_n(x) - 2n\mathcal{H}_{n-1} = 0$   $n \ge 1$ . There is a unique moment functional  $\mathcal{F}$  such that

$$\mathcal{F}(1)=1, \quad \mathcal{F}(\mathcal{H}_n\mathcal{H}_m)=0 \quad \textit{and} \quad \mathcal{F}(\mathcal{H}_n^2)\neq 0.$$

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 Notice that H(x) are not orthogonal with respect to dγ<sub>-1</sub>, but with respect to a moment functional F.

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# Orthogonal basis for inverse Gaussian

#### Theorem

Consider the classical Hermite polynomials given by the Rodrigues formula

$$H_n(x) = (-1)^n e^{x^2} \frac{d}{dx} (e^{-x^2}) \qquad n \ge 0.$$

Then

$$H_n^*(x) = e^{-x^2} H_n(x)$$

forms an orthogonal basis of  $L^2(\mathbb{R}, d\gamma_{-1})$ . Moreover,

$$\mathcal{L}H_n^*=-(2n+2)H_n^*.$$

#### Remark

We also prove  $L^2(\mathbb{R}, d\gamma_{-1})$  boundedness of singular integrals associated to  $\mathcal{L} = \frac{d^2}{dx^2} + 2x \frac{d}{dx}$  (maximal semigroup operator, Riesz transform  $\frac{d}{dx} \mathcal{L}^{-1/2}$ , Littlewood–Paley square functions). The idea is to conjugate with the corresponding operators related to  $L = -\frac{d^2}{dx^2} + 2x \frac{d}{dx}$ .

• Laguerre polynomials: orthonormal basis of  $L^2((0,\infty), x^{\alpha}e^{-x} dx), \alpha > -1$ .

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• Laguerre polynomials: orthonormal basis of  $L^2((0,\infty), x^{\alpha}e^{-x} dx)$ ,  $\alpha > -1$ . Take the Laguerre derivative

$$\mathfrak{D}_{left,a}u = D_{left}u - \left(rac{lpha}{x} - 1
ight)u,$$

choose  $x_0 = 1$  and compute

$$\mathcal{E}^{-1}(x) = \exp\left[-\int_1^x \left(\frac{\alpha}{y}-1\right) dy\right] = \frac{1}{e}x^{-\alpha}e^x.$$

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We found the eigenpolynomials for the corresponding Laplacian:

$$\mathcal{L}_{\alpha,n}(x) = rac{x^{lpha}e^{-x}}{n!}rac{d^n}{dx^n}(x^{n-lpha}e^x) \qquad n \ge 0.$$

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• We also construct eigenpolynomials for the inverse Jacobi measure

$$(1-x)^{-lpha}(1+x)^{-eta}\,dx \qquad {
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• We obtain boundedness of singular integrals in  $L^2$  of the inverse measures.

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### Thank you for your attention!