

Non-linear equations described by Sato, Segal-Wilson theory

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Integrable non-linear PDE

KdV	$\partial_t u = \frac{1}{4} \partial_x^3 u + \frac{3}{2} u \partial_x u$
Boussinesq	$\partial_t^2 u = -\frac{1}{3} \partial_x^2 (\partial_x^2 u + 2u^2)$
Non-linear Schrödinger	$i \partial_t u = \frac{1}{2} \partial_x^2 u \mp 4 u ^2 u$
Sawada-Kotera	$\partial_t u = \partial_x^5 u + 45u^2 \partial_x u - 15 \partial_x u \partial_x^2 u - 15u \partial_x^3 u$
Benjamin-Ono	$\partial_t u = -H(\partial_x^2 u) - u \partial_x u \quad (H: \text{Hilbert transf.})$
Sine-Gordon	$\partial_t^2 u = \partial_x^2 u - \sin u$
Lund-Regge	$\begin{cases} \partial_t^2 u = \partial_x^2 u + \frac{\sin u/2}{2 \cos^2 u/2} \left((\partial_t v)^2 - (\partial_x v)^2 \right) \\ \partial_t^2 v = \partial_x^2 v + 2 \frac{\partial_x u \partial_x v - \partial_t u \partial_t v}{\sin u} \end{cases}$

Lax pair

The above integrable systems are known to have the **Lax pairs**:

$$\text{Lax pair: } \begin{cases} L = \partial_x^\nu + q_{\nu-2}\partial_x^{\nu-2} + \cdots + q_1\partial_x + q_0 \\ P \text{ differential operators s.t. the order of } [P, L] \leq \nu - 2 \end{cases} .$$

Then, setting

$$[P, L] = f_{\nu-2}\partial_x^{\nu-2} + f_{\nu-3}\partial_x^{\nu-3} + \cdots + f_1\partial_x + f_0,$$

with f_j universal polynomials of $\{\partial_x^j q_k\}$, one has

$$\partial_t q_j = f_j \quad (0 \leq j \leq \nu - 2) \quad (\Leftrightarrow \partial_t L = [P, L]).$$

Example: For the KdV equation $\nu = 2$ ($' = \partial_x$)

$$\begin{cases} L = \partial_x^2 + q \\ P = \partial_x^3 + \frac{3}{2}q\partial_x + \frac{3}{4}q' \\ [P, L] = \frac{1}{4}q''' + \frac{3}{2}qq' \end{cases} .$$

- ① Let $C = \{z \in \mathbb{C}; |z| = r\}$ and define the **Hardy spaces** H_{\pm} by

$$H_+ = \left\{ \sum_{j \geq 0} a_j z^j \right\}, \quad H_- = \left\{ \sum_{j < 0} a_j z^j \right\}$$

- ② $\mathcal{A} = \{\mathbf{a}; \mathbf{a}(\lambda) = (a_1(\lambda), a_2(\lambda)) \text{ bounded on } C\}$. For $\mathbf{a} \in \mathcal{A}$ define

$$(\mathbf{a}u)(\lambda) = a_1(\lambda)u(\lambda) + a_2(\lambda)u(-\lambda)$$

Then $W_{\mathbf{a}} = \mathbf{a}H_+ \subset L^2(C)$ satisfies $z^2 W_{\mathbf{a}} \subset W_{\mathbf{a}}$.

- ③ Define projections on $L^2(C)$ by

$$L^2(C) \ni u \rightarrow p_{\pm}u(z) = \pm \frac{1}{2\pi i} \int_C \frac{u(\lambda)}{\lambda - z} d\lambda \in H_{\pm}, \quad (z \in D_{\pm}),$$

and **Toeplitz operator** $T(\mathbf{a})$ on H_+ by

$$T(\mathbf{a})u = p_+(\mathbf{a}u) \text{ for } u \in H_+.$$

- ④ $T(\mathbf{a}) : H_+ \rightarrow H_+$, $\mathcal{A}^{inv} = \left\{ \mathbf{a}; T(\mathbf{a})^{-1} \text{ exists} \right\}$

- $\Gamma = \{g = e^h; h \text{ anal. in a nhb of } D_+\} \Rightarrow g\mathcal{A} \subseteq \mathcal{A} \text{ but } g\mathcal{A}^{inv} \not\subseteq \mathcal{A}^{inv}$
- For $e_{t,x}(z) = e^{xz+tz^3} \in \Gamma$, $\mathbf{a} \in \mathcal{A}$ assume $e_{t,x}^{-1}\mathbf{a} \in \mathcal{A}^{inv}$ for $\forall x, t \in \mathbb{R}$, namely

$$p_+ : e_{t,x}^{-1}W_{\mathbf{a}} (= W_{e_{t,x}^{-1}\mathbf{a}}) \rightarrow H_+ \text{ is 1 to 1}$$

- $f(t, x, z) = (\mathbf{a}T(e_{t,x}^{-1}\mathbf{a})^{-1}1)(z) \in W_{\mathbf{a}}$ satisfies

$$\begin{cases} f = e^{xz+tz^3} (1 + \sum_{j \geq 1} r_j(t, x)z^{-j}) \in e_{t,x}(1 + H_-) \\ f'' = e^{xz+tz^3} (z^2 + r_1z + r_2 + 2r_1' + O(z^{-1})) \\ \Rightarrow e_{t,x}^{-1}(f'' - z^2f - 2r_1'f) = O(z^{-1}) \in H_- \end{cases}$$

- Since $f'' - z^2f - 2r_1'f \in W_{\mathbf{a}}$, applying p_+ one has

$$f'' - z^2f - 2r_1'f = 0 \Rightarrow -\partial_x^2 f + 2r_1'f = -z^2f \quad (\text{Schrödinger})$$

- Comparing the coefficients of z^{-j} , one has

$$2r_{j+1}' + r_j'' - 2r_1'r_j = 0 \implies r_j \text{ is determined by } \left\{ r_1^{(k)} \right\}_{1 \leq k \leq j}$$

- ① Take derivatives ∂_t, ∂_x in $f = e^{xz+tz^3} (1 + r_1 z^{-1} + r_2 z^{-2} + \dots)$

$$\begin{cases} \partial_t f = e_{t,x} (z^3 + r_1 z^2 + r_2 z + r_3 + O(z^{-1})) \\ f' = e_{t,x} (z + r_1 + O(z^{-1})) \\ f''' = e_{t,x} (z^3 + r_1 z^2 + r_2 z + r_3 + 3r_1' z + 3r_2' + 3r_1'' + O(z^{-1})) \end{cases}$$

$$\Rightarrow \partial_t f - f''' + 3r_1' f' - 3(r_2' + r_1'' - r_1 r_1') f \in e_{t,x}^{-1} H_-$$

$$\Rightarrow \partial_t f - f''' + 3r_1' f' - 3(r_2' + r_1'' - r_1 r_1') f = 0.$$

Computing the coefficient of z^{-1} gives

$$\partial_t r_1 - (3r_3' + 3r_2'' + r_1''') + 3r_1' (r_2 + r_1') - 3(r_2' + r_1'' - r_1 r_1') r_1 = 0$$

$$\implies \partial_t r_1 = \frac{1}{4} r_1''' - \frac{3}{2} (r_1')^2$$

② $u = -2a_1' \implies \partial_t u = \frac{1}{4} u''' + \frac{3}{2} u u'$ (KdV)

③ $e_{t,x}(z) = e^{xz+tz^n}$ (odd n) gives the **KdV hierarchy** $\{\text{KdV}_n\}_{\text{odd } n}$

④ $u(t, x) = -2\partial_x^2 \log \tau_a(e_{t,x})$ with $\tau_a(g) = \det(g^{-1} T(ga) T(a)^{-1})$.

- ① ($\nu = 3$) The underlying operator: $L = \partial_x^3 + q_1 \partial_x + q_0$
- ② $\mathcal{A} = \{\mathbf{a}; \mathbf{a}(\lambda) = (a_1(\lambda), a_2(\lambda), a_3(\lambda)) \text{ bdd on } C\}$. $\omega = e^{2\pi i/3}$

$$\begin{cases} (\mathbf{a}u)(\lambda) = a_1(\lambda)u(\lambda) + a_2(\lambda)u(\omega\lambda) + a_3(\lambda)u(\omega^2\lambda) \\ W_{\mathbf{a}} = \mathbf{a}H_+ \implies z^3 W_{\mathbf{a}} \subset W_{\mathbf{a}} \end{cases}$$
- ③ $f = \mathbf{a} \left(T(e_{t,x}^{-1} \mathbf{a})^{-1} \mathbf{1} \right) (z) = e_{t,x} \left(1 + \sum_{r \geq 1} r_j z^{-j} \right) (e_{t,x} = e^{xz+tz^2})$

$$\begin{cases} Lf = z^3 f \text{ with } q_0 = 3(r_1' r_1 - r_1'' - r_2'), q_1 = -3r_1' \\ \partial_t f - \partial_x^2 f + 2r_1' f = 0 \end{cases}$$
- ④ The coefficients $\{r_1, r_2\}$ satisfies

$$\begin{cases} \partial_t r_1 = 2r_2' + r_1'' - 2r_1' r_1 \\ \partial_t r_2 = 2r_1 r_2' - r_2'' - \frac{2}{3} r_1''' + 2(r_1')^2 + 2(r_1'' - r_1' r_1) r_1 \end{cases}$$

$$\partial_t^2 r_1 = \left(-\frac{1}{3} r_1''' + 2(r_1')^2 \right)' \Rightarrow \partial_t^2 q_1 = -\frac{1}{3} (q_1'' + 2q_1^2)'' \text{ (Boussinesq)}$$
- ⑤ e^{xz+tz^n} (even integer n) generates the **Boussinesq hierarchy**

$$\textcircled{1} \mathbf{H}_{\pm} = H_{\pm} \times H_{\pm} \subset \mathbf{L}^2(C), \mathcal{A} = \left\{ A(\lambda) = (a_{ij}(\lambda))_{1 \leq i, j \leq 2}; a_{ij} \text{ bdd on } C \right\}$$

$$W_A = \mathbf{A}\mathbf{H}_+ \implies zW_A \subset W_A$$

$$\textcircled{2} \text{Group: } \Gamma = \left\{ G(g) = \begin{pmatrix} g & 0 \\ 0 & g^{-1} \end{pmatrix}; g = e^h \in \Gamma \right\}$$

$$\textcircled{3} e_{t,x}(z) = e^{i(xz + tz^2)}, I = \text{the identity } 2 \times 2 \text{ matrix}$$

$$F = AT \left(G(e_{t,x})^{-1} A \right)^{-1} I = G(e_{t,x}) \left(I + \sum_{j \geq 1} R_j z^{-j} \right)$$

$$\textcircled{4} \text{Dirac op.: } \mathcal{L} = i \begin{pmatrix} -\partial_x & 0 \\ 0 & \partial_x \end{pmatrix} + \begin{pmatrix} 0 & q_1 \\ q_2 & 0 \end{pmatrix} \text{ with}$$

$$R_1 = \frac{1}{2} \begin{pmatrix} * & q_2 \\ q_1 & * \end{pmatrix}$$

$$\mathcal{L}^t F = z^t F$$

① $R_1 = \frac{1}{2} \begin{pmatrix} * & q_2 \\ q_1 & * \end{pmatrix}$ satisfies $q_1 = 2r_2, q_2 = 2r_1$

$$\begin{cases} i\partial_t q_1 = \frac{1}{2}\partial_x^2 q_1 - q_2 q_1^2 \\ i\partial_t q_2 = -\frac{1}{2}\partial_x^2 q_2 + q_1 q_2^2 \end{cases}$$

② For $A = \begin{pmatrix} a & b \\ b & \bar{a} \end{pmatrix} \implies q_2 = \bar{q}_1 \implies$ **de focussing NLS**

$$\mathcal{L} = \mathcal{L}^* \quad (\text{self-adjoint})$$

③ For $A = \begin{pmatrix} a & b \\ b & -\bar{a} \end{pmatrix} \implies q_2 = -\bar{q}_1 \implies$ **focussing NLS**

$$J\mathcal{L}J = \mathcal{L}^* \quad (J\text{-self-adjoint}) \quad (J = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} C)$$

Open problems

- ① KdV: $L = -\partial_x^2 + q \implies$ Weyl functions m_{\pm}

$$m(z) = \begin{cases} -m_+(-z^2) & \text{if } \operatorname{Re} z > 0 \\ m_-(-z^2) & \text{if } \operatorname{Re} z < 0 \end{cases}$$

- ② $\mathcal{A}^{inv} \ni A = (1, m(z)/z)$ enables us to extend the theory to unbdd. C and more general initial data including smooth ergodic ones.
- ③ de focussing NLS: Set $\phi = \frac{i-m}{i+m}$ with the Weyl function m .

$$A = \begin{pmatrix} 1 & \phi \\ \phi & 1 \end{pmatrix} \in \mathcal{A}^{inv} \implies \text{Extend to general initial data ?}$$

- ④ Boussinesq eq., focussing NLS the underlying L is non-self adjoint
- ⑤ Construct the theory for the other integrable systems.

- [1] M. Sato: *Soliton Equations as Dynamical Systems on an Infinite Dimensional Grassmann Manifolds*, Suriken Koukyuroku **439** (1981), 30-46. (<http://www.kurims.kyoto-u.ac.jp/en/publi-01.html>)
- [2] G.Segal-G.Wilson: Loop groups and equations of KdV type, Publ. IHES, **61** (1985),5-65.

Thank you for attention !