

Simulated Tempering Method in the Infinite Switch Limit with Adaptive Weight Learning

Anton Martinsson*, Jianfeng Lu, Benedict Leimkuhler, Eric Vanden-Eijnden

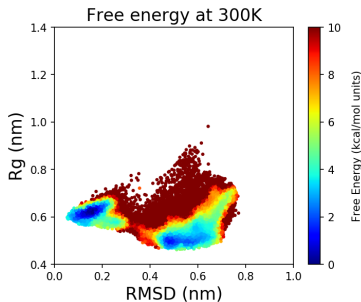
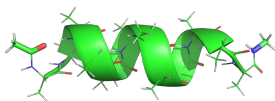
University of Edinburgh,
Maxwell Institute for Graduate Studies in Analysis and its Applications



THE UNIVERSITY of EDINBURGH
School of Mathematics



15th November



- Canonical distribution for q (positions) and p (momentum):

$$\rho_{\beta}(q, p) = Z^{-1}(\beta) e^{-\beta \frac{1}{2} p^T m^{-1} p - \beta V(q)}$$

where $V(q)$ is potential and $Z(\beta)$ is the normalisation constant,
 $\beta > 0$ reciprocal temperature

- Ergodic Average of observable $A(q)$,

$$\mathbb{E}_\beta [A] = \int A(q) \rho_\beta(q, p) dp dq = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T A(q(t)) dt$$

- Standard sampling methods for $\rho_\beta(q, p)$ such as e.g. Monte Carlo or Langevin Dynamics,

$$dq = m^{-1} p dt$$

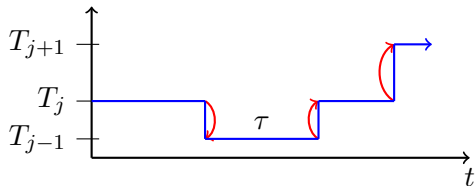
$$dp = -\nabla V(q) dt - \gamma p dt + \sqrt{2\gamma\beta^{-1}m} dW$$

struggle with energetic and entropic barriers

- **Accelerated sampling:** Simulated Annealing, Replica Exchange Molecular Dynamics, **Simulated Tempering**, Wang-Landau, Adaptive Force Biasing, Temperature-accelerated Molecular Dynamics, Hamiltonian Replica Exchange, ...

- Temperature “ladder” with M steps, with spacing ΔT

$$T_{\min} = T_1 < \dots < T_M = T_{\max}$$



- Let $\beta_i = (k_B T_i)^{-1}$ with assigned weight $\omega(\beta_i)$
- Switch every τ steps from $\beta_i \rightarrow \beta_j$ with probability α_{ij} ,

$$\alpha_{ij} = \min\left\{1, \frac{\omega(\beta_i)}{\omega(\beta_j)} e^{-(\beta_i - \beta_j)V(q)}\right\}$$

- Invariant measure with density,

$$\rho(q, p, \beta_i) = C^{-1}(\beta) \omega(\beta_i) e^{-\frac{1}{2} \beta p^T m^{-1} p - \beta_i V(q)}$$

$$\text{where } C(\beta) = \sum_{i=1}^M \int_{\mathcal{D} \times \mathbb{R}^d} \omega(\beta_i) e^{-\frac{1}{2} \beta p^T m^{-1} p - \beta_i V(q)} dq dp$$

- 1 Choice of switch period, τ ?
- 2 How do we justify $\omega(\beta_i) = Z_q^{-1}(\beta_i)$?

$$\mathbb{P}(\beta_i) = \frac{\omega(\beta_i) Z_q(\beta_i)}{\sum_{j=1}^M \omega(\beta_j) Z_q(\beta_j)}.$$

- 3 Can we learn $\omega(\beta_i)$ weights on-the-fly?

Remarks:

- Grid Spacing, ΔT . Is there a way to motivate the choice?
- How do we calculate an average from ρ_{β_i} i.e $\mathbb{E}_{\beta_i}[A]$?

Infinite Switch Limit of Simulated Tempering, $\tau \rightarrow 0$.

- Follows same large deviation approach as for REMD by Dupuis et. al.^{1 2}
- ST in $\tau \rightarrow 0$ limit for some continuous $\beta_c \in [\beta_{\min}, \beta_{\max}]$ is equivalent to sampling the averaged potential³,

$$\bar{V}(q) = -\beta^{-1} \log \int_{\beta_{\min}}^{\beta_{\max}} \omega(\beta_c) e^{-\beta_c V(q)} d\beta_c$$

for some prior known weights $\omega(\beta_c)$.

¹ Dupuis et. al. 2012.

² Plattner et. al 2011.

³ A.M. et. al. 2018.

- The rate functional $\mathbb{P}(\nu_T \approx \mu) \asymp \exp[-T^{-1}I_\tau(\mu)]$ of the ergodic dynamics is a monotonically decreasing function of τ , i.e if

$$\tau < \tau' \implies I_\tau(\mu) \geq I_{\tau'}(\mu)$$

For faster convergence to μ , let $\tau \rightarrow 0$.

- The dynamical equations in $\tau \rightarrow 0$ limit,

$$dq = m^{-1} p dt,$$

$$dp = -\beta^{-1} \bar{\beta}(V(q)) \nabla V dt - \gamma p dt + \sqrt{2\gamma\beta^{-1}m^{-1}} dW_p$$

which implies that we have averaged potential,

$$\bar{V}(q) = -\beta^{-1} \log \int_{\beta_{\min}}^{\beta_{\max}} \omega(\beta_c) e^{-\beta_c V(q)} d\beta_c$$

- 1 Choice of switch period, τ ?
- 2 How do we justify $\omega(\beta_i) = Z_q^{-1}(\beta_i)$?

$$\mathbb{P}(\beta_i) = \frac{\omega(\beta_i) Z_q(\beta_i)}{\sum_{j=1}^M \omega(\beta_j) Z_q(\beta_j)}.$$

- 3 Can we learn $\omega(\beta_i)$ weights on-the-fly?

Remarks:

- Grid Spacing, ΔT . Is there a way to motivate the choice?
- How do we calculate an average from ρ_{β_i} i.e $\mathbb{E}_{\beta_i}[A]$?

Justification of $\omega(\beta) \propto Z_q^{-1}(\beta)$.

- The density function of $V(q)$

$$\bar{\rho}(E) = \frac{\int_{\beta_{\min}}^{\beta_{\max}} e^{-\beta_c E} \omega(\beta_c) d\beta_c}{\int_{\beta_{\min}}^{\beta_{\max}} Z_q(\beta_c) \omega(\beta_c) d\beta_c} \Omega(E),$$

where $\Omega(E) = \int_{\mathcal{D}} \delta(V(q) - E) dq$.

- In that large system size limit,

$$\bar{\rho}(E) \asymp 1$$

when $\omega(\beta) \propto Z_q^{-1}(\beta)$.

- 1 Choice of switch period, τ ?
- 2 How do we justify $\omega(\beta_i) = Z_q^{-1}(\beta_i)$?

$$\mathbb{P}(\beta_i) = \frac{\omega(\beta_i) Z_q(\beta_i)}{\sum_{j=1}^M \omega(\beta_j) Z_q(\beta_j)}.$$

- 3 Can we learn $\omega(\beta_i)$ weights on-the-fly?

Remarks:

- Grid Spacing, ΔT . Is there a way to motivate the choice?
- How do we calculate an average from ρ_{β_i} i.e $\mathbb{E}_{\beta_i}[A]$?

Adaptive $\omega(\beta_c)$ adjustment as we learn $Z_q(\beta_c)$.

- For $\beta_c \in [\beta_{\min}, \beta_{\max}]$ we can construct

$$z(t, \beta_c) = \frac{1}{t} \int_0^t \frac{e^{-\beta_c V(q(s))}}{\int_{\beta_{\min}}^{\beta_{\max}} \omega(\beta'_c) e^{-\beta'_c V(q(s))} d\beta'_c} ds$$

which can be calculated for all $t > 0$.

- In the limit as $t \rightarrow \infty$,

$$\lim_{t \rightarrow \infty} z(t, \beta_c) = \frac{Z_q(\beta_c)}{\int_{\beta_{\min}}^{\beta_{\max}} Z_q(\beta'_c) \omega(\beta'_c) d\beta'_c},$$

- This implies that we learn ratios of $Z_q(\beta_c)$ regardless of knowledge of $\omega(\beta_c)$,

$$\lim_{t \rightarrow \infty} \frac{z(t, \beta_c)}{z(t, \beta'_c)} = \frac{Z_q(\beta_c)}{Z_q(\beta'_c)}$$

- Define $z(t, \beta_c)$ as,

$$z(t, \beta_c) = \frac{1}{t} \int_0^t \frac{e^{-\beta_c V(q(s))}}{\int_{\beta_{\min}}^{\beta_{\max}} \omega(\beta'_c) e^{-\beta'_c V(q(s))} d\beta'_c} ds$$

- Adjusting the weights according to,

$$\kappa \dot{\omega}(t, \beta_c) = z^{-1}(t, \beta_c) - \lambda(t) \omega(t, \beta_c), \text{ with } \lambda(t) = \int_{\beta_{\min}}^{\beta_{\max}} z^{-1}(t, \beta_c) d\beta_c$$

will ensure that $\omega(t, \beta_c) \propto Z_q^{-1}(\beta_c)$ as $t \rightarrow \infty$.

- 1 Choice of switch period, τ ?
- 2 How do we justify $\omega(\beta_i) = Z_q^{-1}(\beta_i)$?

$$\mathbb{P}(\beta_i) = \frac{\omega(\beta_i) Z_q(\beta_i)}{\sum_{j=1}^M \omega(\beta_j) Z_q(\beta_j)}.$$

- 3 Can we learn $\omega(\beta_i)$ weights on-the-fly?

Remarks:

- Grid Spacing, ΔT . Is there a way to motivate the choice?
- How do we calculate an average from ρ_{β_i} i.e $\mathbb{E}_{\beta_i}[A]$?

- The equations of motion are,

$$dq = m^{-1} p dt,$$

$$dp = -\beta^{-1} \hat{\beta}(t, V(q)) \nabla V dt - \gamma p dt + \sqrt{2\gamma\beta^{-1}m} dW_p$$

- We need to calculate the following force re-scaling

$$\hat{\beta}(t, V(q)) = \frac{\int_{\beta_{\min}}^{\beta_{\max}} \beta_c \omega(t, \beta_c) e^{-\beta_c V(q)} d\beta_c}{\int_{\beta_{\min}}^{\beta_{\max}} \omega(t, \beta_c) e^{-\beta_c V(q)} d\beta_c}$$

- We therefore discretise $[\beta_{\min}, \beta_{\max}]$ with the aim of calculating $\hat{\beta}(t, V(q))$

- In standard Simulated Tempering

$$\mathbb{E}_{\beta_c} [A] = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T A(q(t)) \mathbb{1}_{\beta_c} dt$$

- We instead use a reweighting of the entire trajectory (importance sampling),

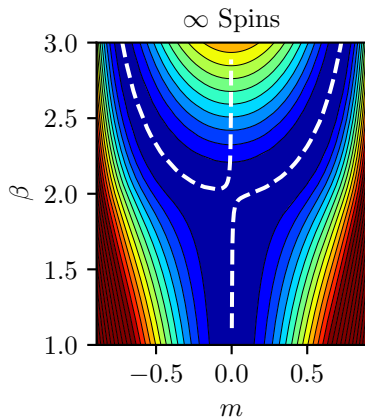
$$\mathbb{E}_{\beta_c} [A] = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T A(q(t)) W_{\beta_c}(t, q(t)) dt$$

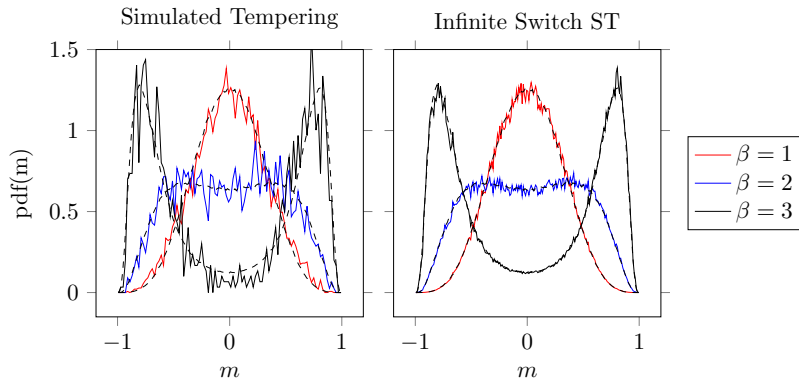
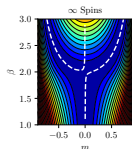
where,

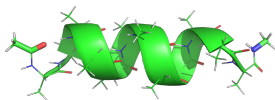
$$W_{\beta_c}(t, q) = \frac{z^{-1}(t, \beta_c)}{\int_{\beta_{\min}}^{\beta_{\max}} \omega(t, \beta'_c) e^{-(\beta'_c - \beta_c)V(q)} d\beta'_c}$$

where $z(t, \beta_c)$ is the estimate of $Z_q(\beta_c)$ at time t .

Curie Weiss Magnet – mean field Isings model





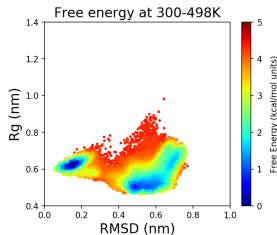
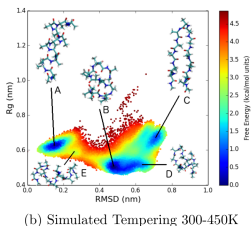
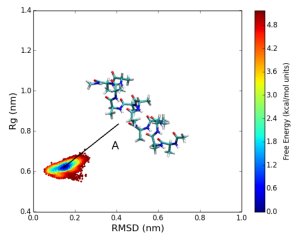
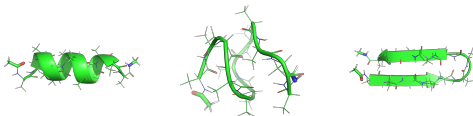


- Molecular Integration Simulation Toolkit (MIST)⁴ integrates with Amber 14, GROMACS 5.0.2, NAMD-Lite 2.0.3, LAMMPS
- In vacuum, using Gromacs with MIST⁵
- Measure RMSD and radius of gyration, r_g , from initial helix state

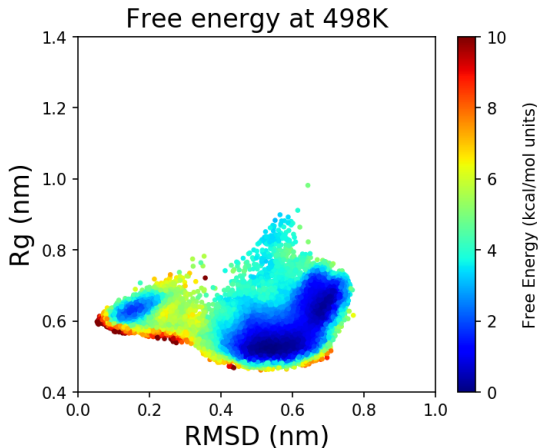
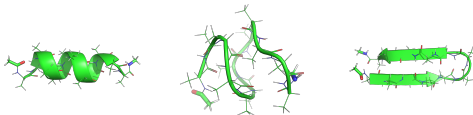
⁴**Bethune et. al.** *MIST: A Simple and Efficient Molecular Dynamics Abstraction Library for Integrator Development*, **Computer Physics Communications**, 2018

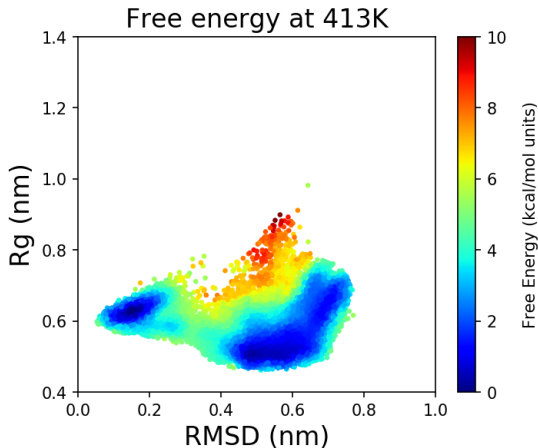
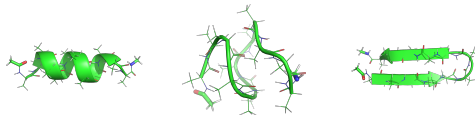
⁵<https://bitbucket.org/extasy-project/mist/wiki/Home>

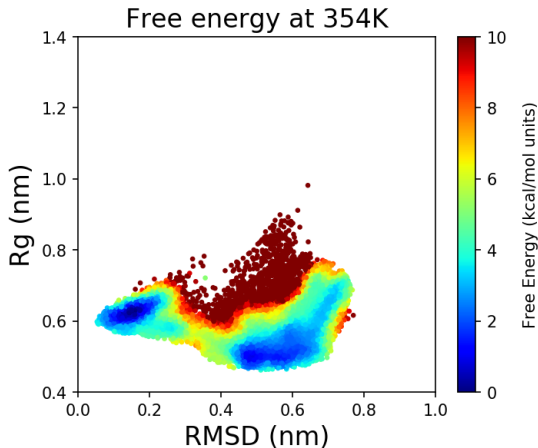
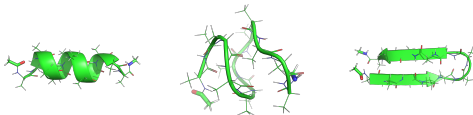
Alanine-12: LD vs ST vs ISST

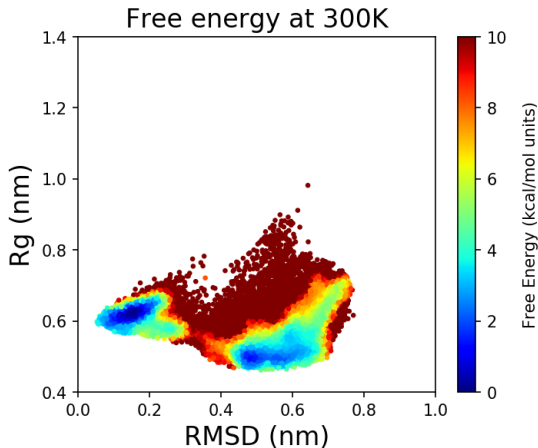
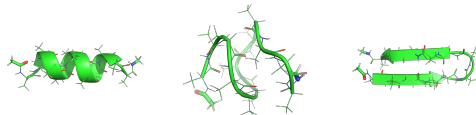


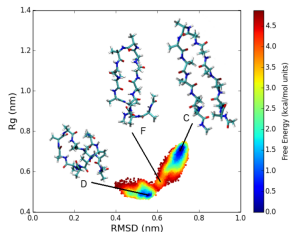
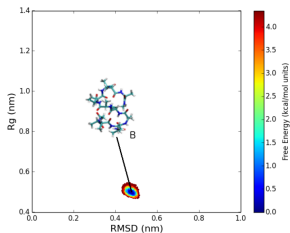
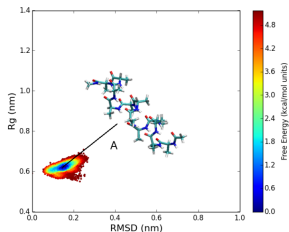
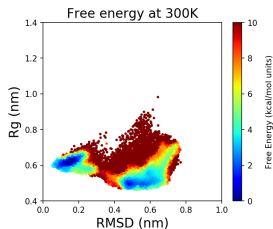
⁶Bethune et. al., 2018















⁷Bethune et. al., 2018

- Operate ST in the Infinite Switch limit, with adaptive weight learning – without temperature being a dynamical variable ⁸
- We have presented work on MD, future work is looking into applying the method in data science and ML
- We have implemented the method in MIST and it is available to download at <https://bitbucket.org/extasy-project/mist/wiki/Home>

⁸Anton Martinsson, Jianfeng Lu, Ben Leimkuhler, Eric Vanden-Eijnden, *Simulated Tempering Method in the Infinite Switch Limit with Adaptive Weight Learning*, **to appear in JSTAT**, 2018

-  Paul Dupuis, Yufei Liu, Nuria Plattner, and J. D. Doll.
On the Infinite Swapping Limit for Parallel Tempering.
Multiscale Modeling & Simulation, 10(3):986–1022, 2012.
-  Nuria Plattner, J D Doll, Paul Dupuis, Hui Wang, Yufei Liu, and J E Gubernatis.
An infinite swapping approach to the rare-event sampling problem.
The Journal of chemical physics, 135(13):134111, oct 2011.
-  Iain Bethune, Ralf Banisch, Elena Breitmoser, et. al.
MIST: A Simple and Efficient Molecular Dynamics Abstraction
Library for Integrator Development
Computer Physics Communications , 2018
-  Anton Martinsson, Jianfeng Lu, Ben Leimkuhler, Eric
Vanden-Eijnden,
Simulated Tempering Method in the Infinite Switch Limit with
Adaptive Weight Learning,
to appear in JSTAT, 2018