

# **Convergence to equilibrium for a thin film equation**

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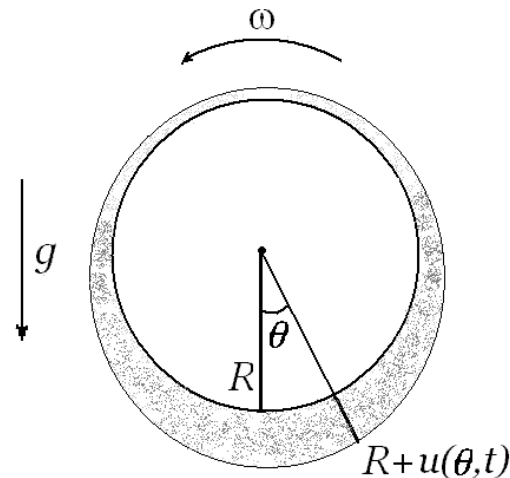
joint work with

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# The physical model

$$u_t + \partial_x(u^3(u_{xxx} + u_x - \sin x)) + \omega u_x = 0.$$

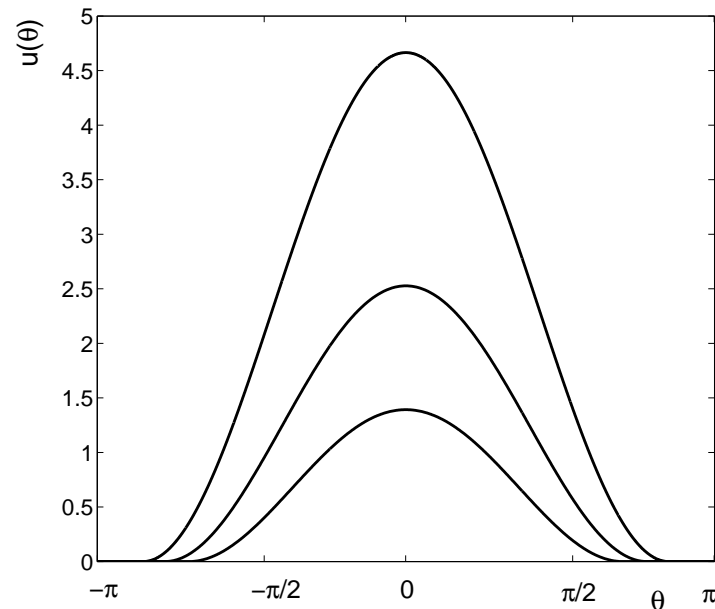


- $u \geq 0$ , periodic;
- $\partial_x(u^3(u_{xxx} + u_x))$  surface tension term;
- $-\partial_x(u^3 \sin x)$  gravitational drainage;
- rotation speed  $\omega$ .

[Moffatt 1976, Pukhnachev 1977, Benilov & al. 2000-date]

# Summary of results (Pukhnachev's model with $\omega = 0$ )

- For every mass there is a **unique energy minimizer**  $u^*$ ;
- $u^*$  is **globally attractive**;
- (no better than) **power-law decay**  $\|u(\cdot, t), u^*\|_{H^1} \geq Ct^{-\frac{2}{3}}$ .



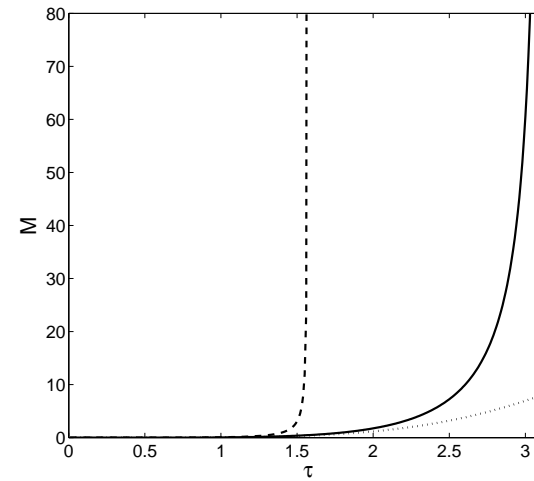
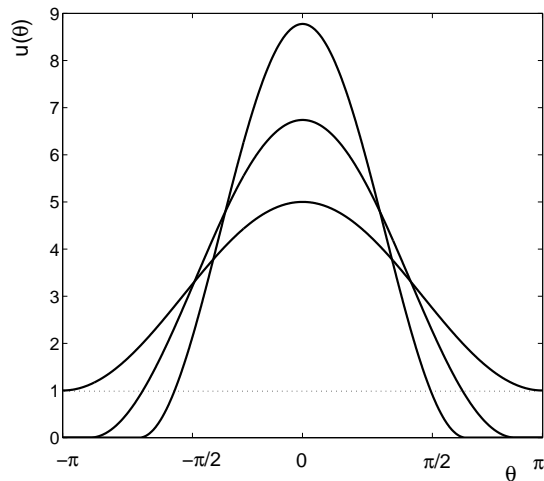
Energy minimizers for three values of the mass.

# Energy minimizers

For every value of  $\alpha, n, M > 0$ , functional

$$E(u) = \frac{1}{2} \int_{-\pi}^{\pi} u_x^2 - \alpha^2 u^2 dx - \int_{-\pi}^{\pi} u \cos x dx .$$

has a unique global minimizer of mass  $M$ .

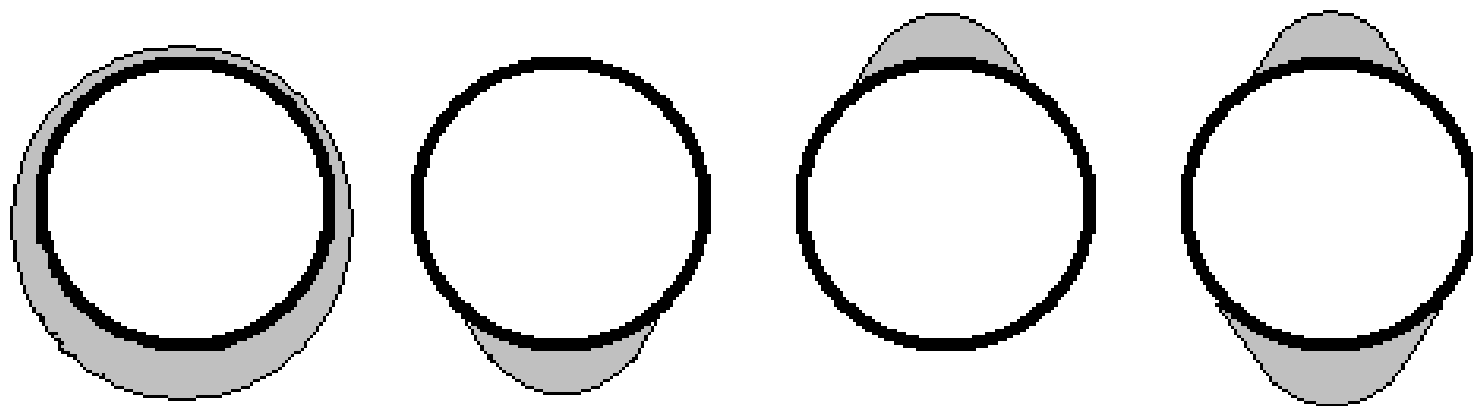


Minimizers for  $\alpha < 1$ ,  $\alpha = 1$ , and  $\alpha > 1$ .

*(Note that  $E$  is not convex when  $\alpha > 1$ .)*

## Other critical points of the energy

Droplet-shaped critical points have **zero contact angle**.  
For  $\alpha \leq 1$ , the global minimizer is the unique critical point.  
For  $\alpha > 1$ , there may be others (depending on the mass):

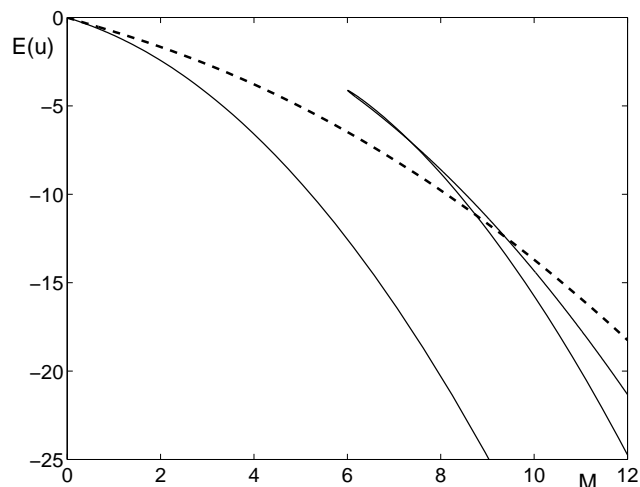


Are these critical points all the steady states?

Does Lyapunov's principle apply?

*“The  $\omega$ -limit set of an orbit under a gradient flow consists of critical points of the Lyapunov function.”*

# Bifurcation diagram



Energy levels of critical points, depending on the mass.

$\alpha = 1$  (dashed line),  $\alpha = \sqrt{2}$  (solid line).

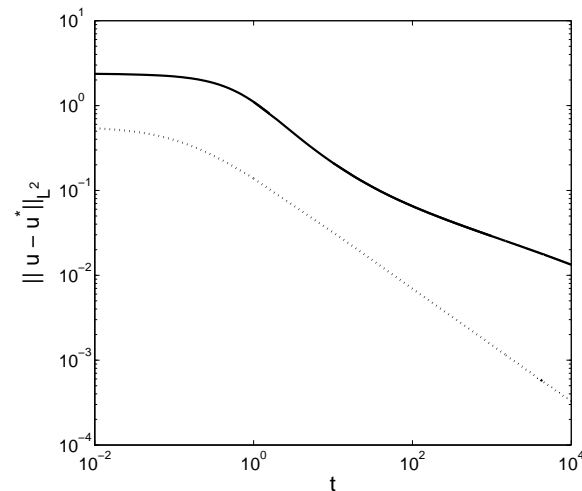
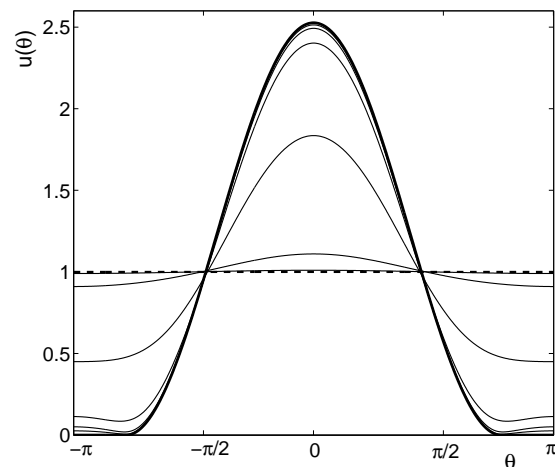
# Convergence to the energy minimizer

Let  $u(\cdot, t)$  be a solution of finite entropy, and let  $u^*$  be the global energy minimizer of the same mass. (If  $\alpha > 1$ , assume also that no other critical points have energy below  $E(u)$ .) Then:

- The solution  $u(t)$  converges to  $u^*$  as  $t \rightarrow \infty$ .  
(Proof: An energy-entropy compactness argument.)
- For  $n > \frac{3}{2}$ , the distance from a droplet cannot decay faster than a power law.  
(Proof: Entropy grows at most linearly in  $t$ .)
- If  $u^*$  is positive, then  $u(\cdot, t)$  converges exponentially.  
(Proof: Compare the dissipation with the energy.)

# Open questions

- What is the **rate of convergence** really? (Perhaps  $t^{-\frac{1}{3}}$  ?)  
How can we linearize around a droplet? [*Slepčev 2008*]



- Do all solutions **converge to equilibrium** (even when there are many steady states?)
- How to take advantage of the **gradient flow** structure? [*Otto 1998, ..., Ambrosio-Gigli-Savaré (book), ..., ..., Kamalinejad 2012,...*]