

# On the rough-smooth interface in the two-periodic Aztec diamond

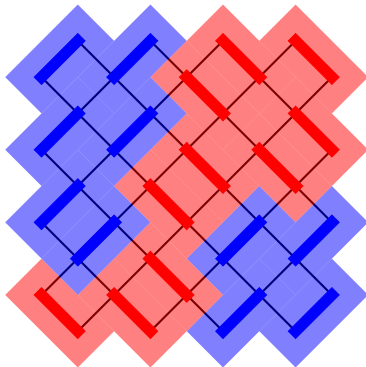
Kurt Johansson

(Joint work with Sunil Chhita and Vincent Beffara)

BIRS, November 21, 2019

# Aztec diamond

A **domino tiling** of an Aztec diamond shape corresponds to a **dimer configuration** on the Aztec graph.



# Probability measure

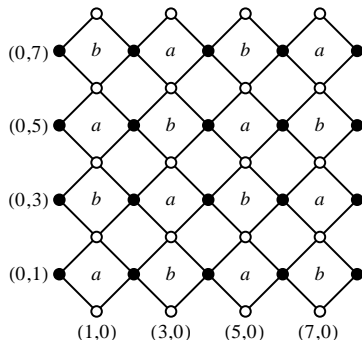
Let  $\nu(e) > 0$  be the **weight** of the edge  $e$  in the graph  $\mathcal{G}$ . The probability of a certain **dimer cover**  $C$ , i.e. each vertex is covered exactly once, is

$$\frac{1}{Z} \prod_{e \in C} \nu(e).$$

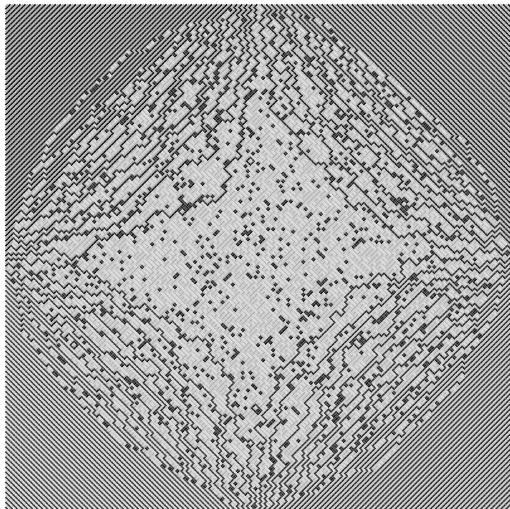
$Z$  is the **partition function**.

## Two Periodic Weighting

The **two-periodic weighting** of the Aztec diamond is defined in the following way. For a two-colouring of the faces, the edge weights around a particular coloured face alternate between  $a$  and  $b$ , we have  $a$ -edges and  $b$ -edges. E.g. for a size 4 Aztec diamond



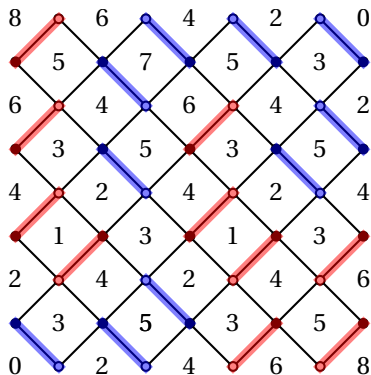
# Random tiling of a two-periodic Aztec diamond



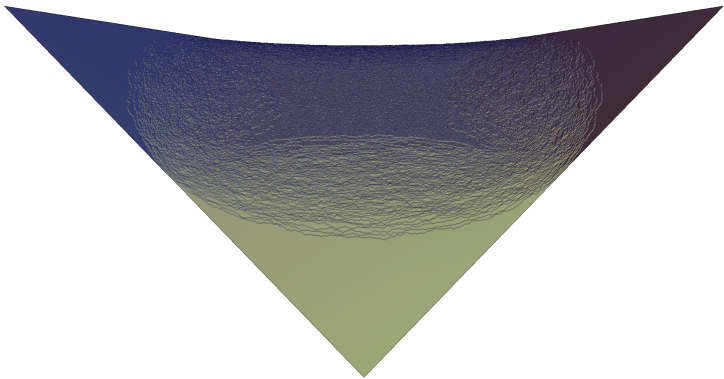
## Aztec diamond height function

To each tiling of an Aztec diamond we can associate a **height function**. The heights sit on the faces of the Aztec graph. The height differences between two faces are given by

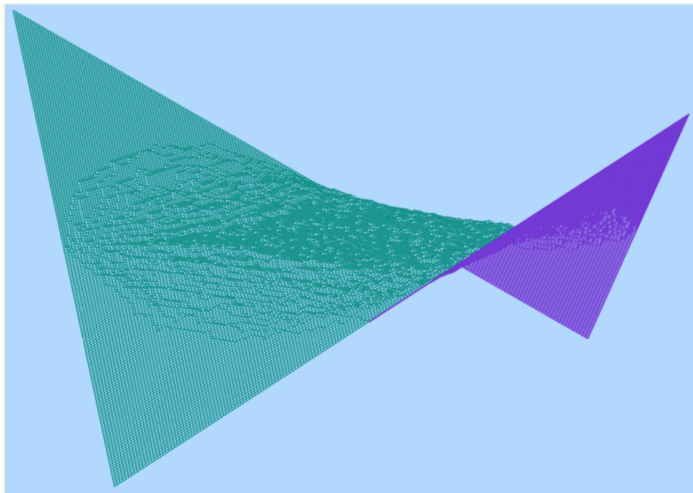
- $+3$  ( $-3$ ) if we cross a dimer with a white vertex to the right (left)
- $+1$  ( $-1$ ) if we do not cross a dimer and have a white vertex to the left (right)



# Two-periodic Aztec diamond height function



# Two-periodic Aztec diamond height function



Picture by V. Beffara



# Kasteleyn Matrix

We choose a **Kasteleyn sign**,  $s(e)$ ,  $|s(e)| = 1$ , for each edge with certain properties, and then define the **Kasteleyn matrix**  $\mathbb{K}$  with elements

$$\mathbb{K}(b_i, w_j) = s(b_i, w_j)\nu(b_i, w_j).$$

This is a signed weighted adjacency matrix for the graph.  
For the Aztec diamond graph we can take

$$\mathbb{K}(b, w) = \begin{cases} \nu(bw) & \text{if } e = (b, w) \text{ is horizontal} \\ i\nu(bw) & \text{if } e = (b, w) \text{ is vertical} \\ 0 & \text{otherwise (i.e. no edge between } b \text{ and } w) \end{cases}$$

# Kasteleyn's theorem

Let  $\mathbb{K}$  be a Kasteleyn matrix

Theorem (Kasteleyn)

$$\det(\mathbb{K}) = SZ,$$

where  $Z$  is the partition function, and  $|S| = 1$ .

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It follows from Kasteleyn's theorem that

## Theorem (Montroll-Potts-Ward, Kenyon)

If  $e_i = (b_i, w_i)$ , then the probability that  $e_1, \dots, e_m$  belong to a dimer cover is

$$\mathbb{P}(e_1, \dots, e_m) = \det (\mathbb{K}(b_i, w_i) \mathbb{K}^{-1}(w_i, b_j))_{1 \leq i, j \leq m}$$

This means that the dimers form a **determinantal point process** with correlation kernel  $K(e_i, e_j) = \mathbb{K}(b_i, w_i) \mathbb{K}^{-1}(w_i, b_j)$ ,  $e_i = (b_i, w_i)$ .

# A simulation of the two-periodic Aztec diamond

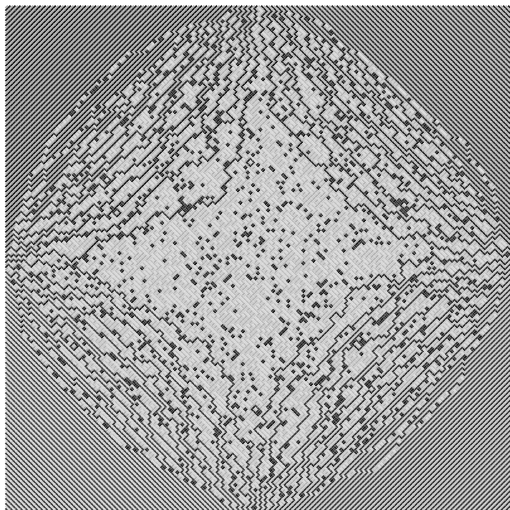
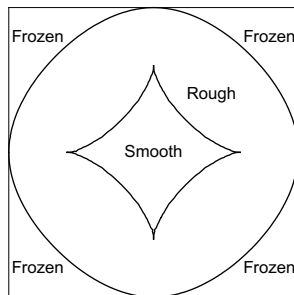


Figure:  $n = 200$ ,  $a = 0.5$ ,  $b = 1$  with 8 grayscale colors

# Phases



The curve in the picture is a degree 8 curve with two real components. We get three regions which are called **frozen**, **rough** and **smooth**.

# Phases

Kenyon, Okounkov and Sheffield have characterized the different **limiting translation invariant Gibbs measures** that are possible for bipartite dimer models on the plane.

There are three classes of Gibbs measures, **frozen**, **rough** and **smooth**, given by an appropriate infinite, translation-invariant full-plane inverse Kasteleyn matrix  $\mathbb{K}^{-1}$ .

# Phases

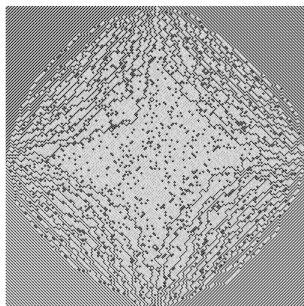
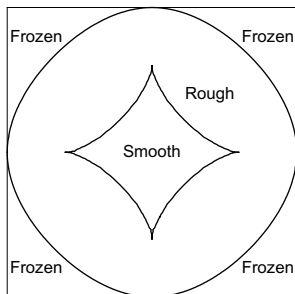
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*Correlations* between dominos decay polynomially with distance in the rough region, and exponentially in the smooth region.

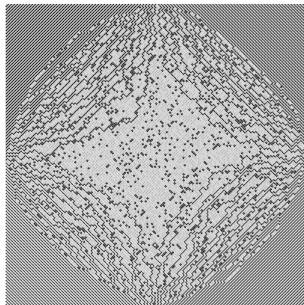
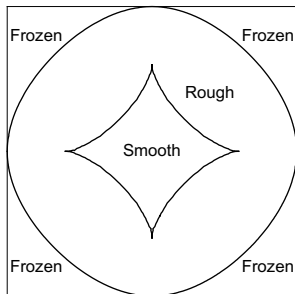
# Rough-smooth boundary

We now have two types of boundaries, the rough-frozen boundary and the rough-smooth boundary.



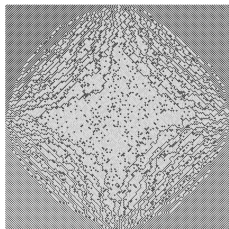


# Rough-smooth boundary



What can we say about the interface fluctuations at the rough-smooth boundary? What is actually the interface?

# Rough-smooth boundary

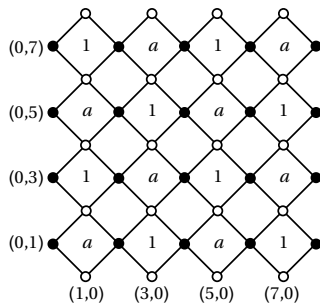


What can we say about the interface fluctuations at the rough-smooth boundary? What is actually the interface?

The rough-frozen interface is well-defined, the first place when the regular pattern is broken

At the rough-smooth boundary the situation is less clear. How should we define the interface combinatorially at the discrete level?

# Formula for the inverse Kasteleyn matrix in the two-periodic case



The coordinate system that we use is indicated in the figure.

# The inverse Kasteleyn Matrix

## Theorem (Chhita-J. based on Chhita-Young)

Consider an Aztec diamond of size  $n = 4m$  with the two-periodic weighting and let  $\mathbb{K}_m$  be its Kasteleyn matrix. Then,

$$\mathbb{K}_m^{-1}((x_1, x_2), (y_1, y_2)) = \mathbb{K}_{sm}^{-1}((x_1, x_2), (y_1, y_2)) - \sum_{i=1}^4 B_i((x_1, x_2), (y_1, y_2)),$$

where  $\mathbb{K}_{sm}^{-1}$  is the full-plane inverse Kasteleyn matrix for the smooth phase, which has an explicit double integral formula, and  $B_1, \dots, B_4$  are contributions also given by explicit double integral formulas.

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Recently a more systematic approach has been developed to get the inverse Kasteleyn matrix or, more specifically, a closely related correlation kernel for an associated particle process, see *The two periodic Aztec diamond and matrix valued orthogonal polynomials*, by Maurice Duits, Arno B.J. Kuijlaars and *Correlation functions for determinantal processes defined by infinite block Toeplitz minors*, by T. Berggren, M. Duits.

# Airy kernel point process

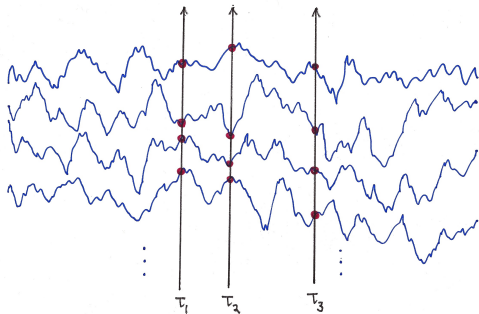


Figure: The **Airy line ensemble**. The top path is the **Airy process**.

# Airy kernel point process

The **extended Airy point process** is a determinantal point process on parallel lines  $\{\tau_q\} \times \mathbb{R}$ ,  $1 \leq q \leq L_1$  in  $\mathbb{R}^2$ . We can think of it as a random measure  $\mu_{\text{Ai}}$  defined via a Laplace transform. Let  $A_p$ ,  $1 \leq p \leq L_2$ , be disjoint intervals in  $\mathbb{R}$ ,  $w_{p,q} \in \mathbb{C}$ ,

$$\begin{aligned} & \mathbb{E} \left[ \exp \left( \sum_{p=1}^{L_2} \sum_{q=1}^{L_1} w_{p,q} \mu_{\text{Ai}}(\{\tau_q\} \times A_p) \right) \right] \\ &= \det \left( I + (e^\Psi - 1) K_{\text{extAi}} \right)_{L^2(\{\tau_1, \dots, \tau_q\} \times \mathbb{R})}, \end{aligned}$$

where

$$\Psi(x) = \sum_{q=1}^{L_1} \sum_{p=1}^{L_2} w_{p,q} \mathbb{I}_{\{\tau_q\} \times A_p}(x).$$

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where

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Recall that the **extended Airy kernel** is given by

$$K_{\text{extAi}}(\tau_1, \xi_1; \tau_2, \xi_2) = -\phi_{\tau_1, \tau_2}(\xi_1, \xi_2) + \tilde{K}_{\text{extAi}}(\tau_1, \xi_1; \tau_2, \xi_2),$$

where

$$\tilde{K}_{\text{extAi}}(\tau_1, \xi_1; \tau_2, \xi_2) = \int_0^\infty e^{-\lambda(\tau_1 - \tau_2)} \text{Ai}(\xi_1 + \lambda) \text{Ai}(\xi_2 + \lambda) d\lambda.$$

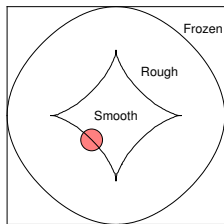


# Asymptotics for the inverse Kasteleyn matrix at the rough-smooth boundary

Let  $x = (x_1, x_2)$  be a white vertex and  $y = (y_1, y_2)$  a black vertex.  
Scaling around a point at the rough-smooth boundary:

$$(x_1, x_2) = (4[\rho m] + 2[c_1 \xi_1 m^{1/3}](1, 1) - 2[c_2 \tau_1 m^{2/3}](-1, 1),$$

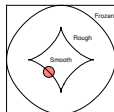
$$(y_1, y_2) = (4[\rho m] + 2[c_1 \xi_2 m^{1/3}](1, 1) - 2[c_2 \tau_2 m^{2/3}](-1, 1).$$



# Asymptotics for the inverse Kasteleyn matrix at the rough-smooth boundary

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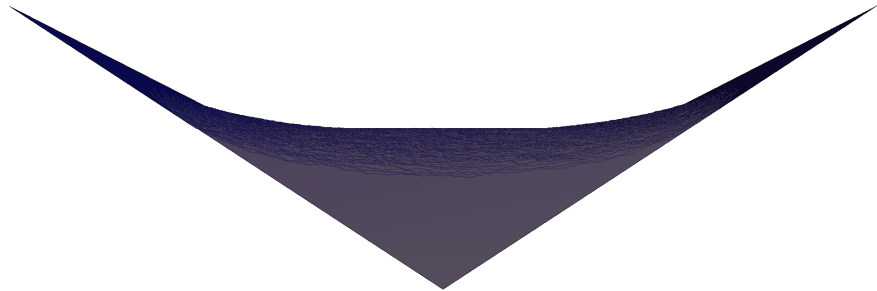


Asymptotics

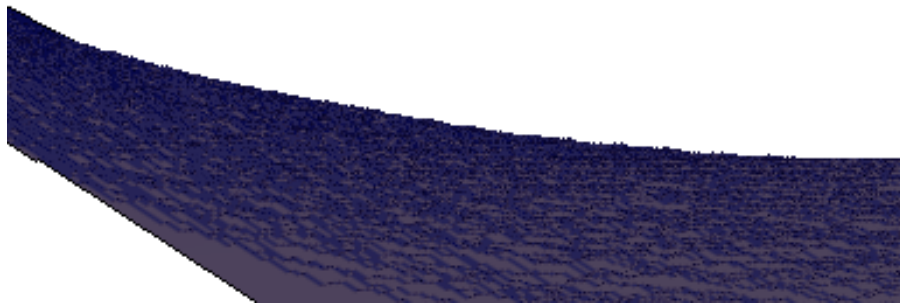
$$\mathbb{K}_m^{-1}(x, y) = \mathbb{K}_{sm}^{-1}(x, y) - m^{-1/3}(\text{pre-factor}) \tilde{K}_{\text{extAi}}(\tau_1, \xi_1 + \tau_1^2; \tau_2, \xi_2 + \tau_2^2)(1 + o(1))$$

as  $m \rightarrow \infty$ .

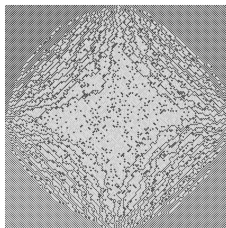
## Random measure from the height function



# Random measure from the height function



# Random measure



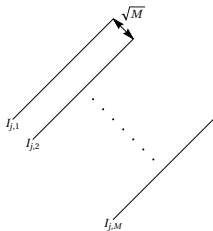
Height differences between two points in a vicinity of the rough-smooth boundary are due to two effects:

- Small and basically independent height fluctuations due to the "surrounding smooth phase".
- Long distance correlated effects due to the large scale structures that we see in the figure.
- By taking suitable averages of height differences we could hope to eliminate the small scale smooth phase effects. This is the idea behind the definition of a certain random signed measure.

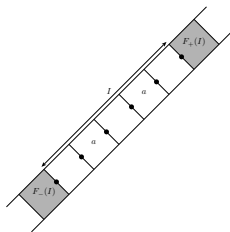
# Random measure

Consider a two-periodic Aztec diamond of size  $n = 4m$ .

We want to imbed the intervals  $A_p$  as discrete intervals of length  $\sim m^{1/3}$  in the Aztec diamond at the rough-smooth boundary. Consider only one interval,  $A = [a^l, a^r]$ . We want to imbed  $M = \lceil (\log m)^4 \rceil$  copies of it as discrete intervals starting and ending at a-faces a certain distance apart.



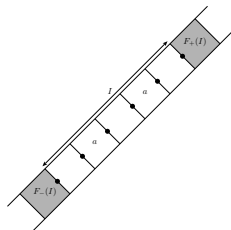
# Random measure



The **height change** along a discrete interval

$$\Delta h(I) = h(F_+(I)) - h(F_-(I)).$$

# Random measure



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**Random signed measure**

$$\mu_m(\{\beta\} \times A) = \frac{1}{4M} \sum_{k=1}^M \Delta h(I_k).$$



# Random measure

## Random signed measure

$$\mu_m(\{\beta\} \times A) = \frac{1}{4M} \sum_{k=1}^M \Delta h(l_k).$$

## Theorem (Beffara, Chhita, J., 18)

*The random signed measure  $\mu_m$  converges to  $\mu_{Ai}$ .*

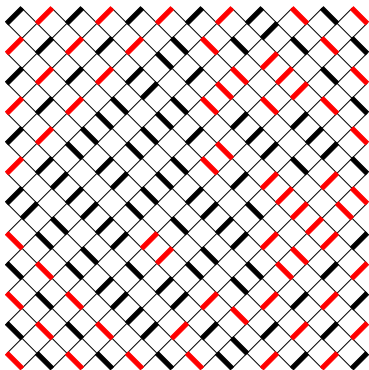
For one interval the result is

$$\lim_{m \rightarrow \infty} \mathbb{E}[e^{w\mu_m(A)}] = \mathbb{E}[e^{w\mu_{Ai}(A)}],$$

for  $w \in \mathbb{C}$ ,  $|w| < R$ .

# Squishing

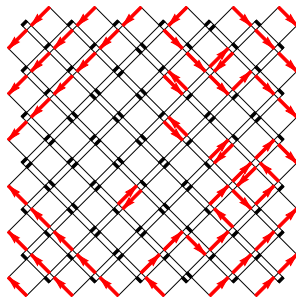
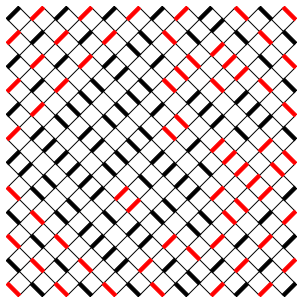
An  $a$ -**dimer** is a dimer that covers an  $a$ -edge. They are **oriented** from white to black.



**Figure:** The red dimers are  $a$ -dimers, and the black  $b$ -dimers.

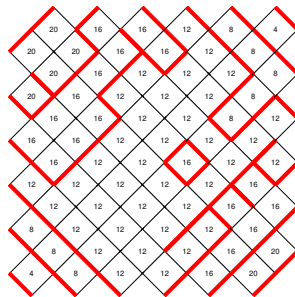
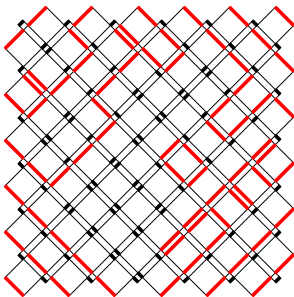
# Squishing

We let the  $b$ -faces become smaller, go to zero in size.



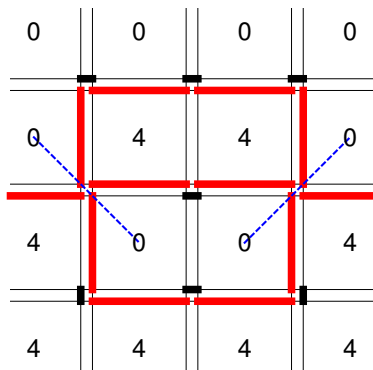
# Squishing

We get **double edges**, **loops** and **paths**.



# Paths and Loops

To get a unique split between paths and loops and get well-defined loops we need a convention. We use **mirrors**.



# Paths

The paths go between the boundaries.

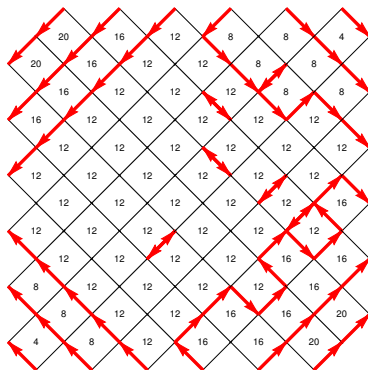


Figure: After squishing.

# Paths

The paths go between the boundaries.

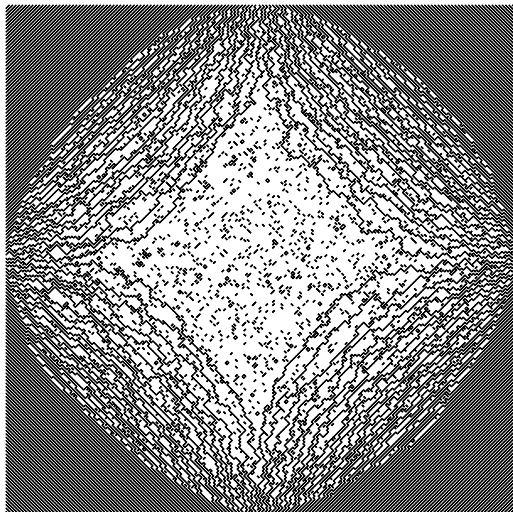
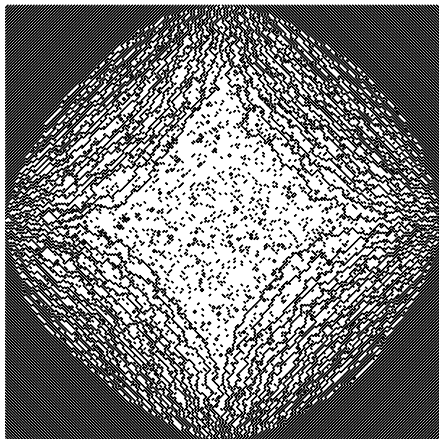


Figure: After squishing,  $n = 300$ ,  $a = 0.5$ .

## What we would like to prove

With high probability, if we go along the main diagonal there is a last path in the third quadrant close to the asymptotic rough-smooth boundary and this path converges to the Airy process.



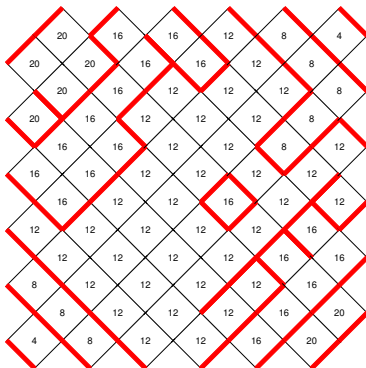


## What we can prove

Let  $h(f)$  be the height at the  $a$ -face  $f$  in the Aztec diamond. Then we can split it into two parts:

$$h(f) = h_\ell(f) + h_c(f)$$

where  $h_\ell(f)$  is the **loop height** and  $h_c(f)$  is the **corridor height**.



## What we can prove

Assume that  $a < 1/3$ . Imbed the interval  $A$  as a discrete interval of length  $\sim m^{1/3}$  in the Aztec diamond at the rough-smooth boundary. Define the **new random signed measure**

$$\kappa_m(\{\beta\} \times A) = \frac{1}{4}(h_c(F_+) - h_c(F_-)),$$

where  $F_+$  and  $F_-$  are the end-faces of the discrete imbedded interval. Then  $\kappa_m(\{\beta\} \times A)$  converges weakly to  $\mu_{\text{Ai}}(\{\beta\} \times A)$  as  $m \rightarrow \infty$ , where  $\mu_{\text{Ai}}$  is the Airy kernel point process.

We expect that with high probability  $\kappa_m$  is actually a positive measure. We should think of  $\kappa_m$  as counting the number of paths between the two faces.

# Ingredients in the proof

- Show that the averaging in the random signed measure  $\mu_m$  can be done along a single line instead of on parallel lines.
- We control the size of the loops; in a box of size  $L$  they are no bigger than  $C \log L$ . This uses a Peierls' type argument and requires  $a < 1/3$ .
- In a not too large region at the rough-smooth boundary the two-periodic Aztec measure can be replaced by the full-plane smooth measure.
- There are no bi-infinite paths in the full-plane smooth phase.

Thank you for listening!

