

From stochastic thermodynamic to thermodynamic inference

Udo Seifert

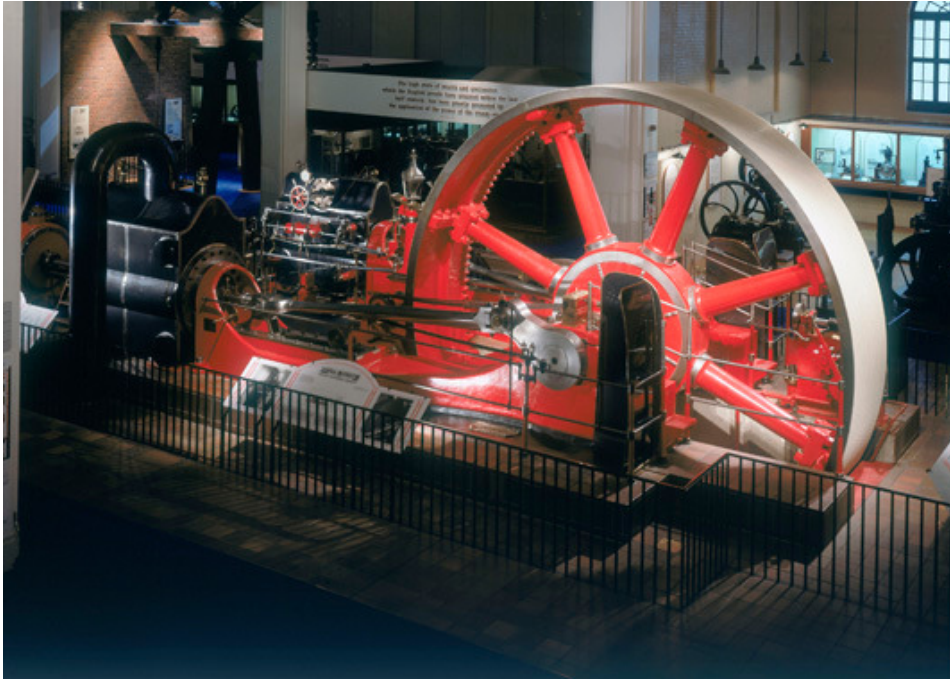
II. Institut für Theoretische Physik, Universität Stuttgart

- Principles of stochastic thermodynamics
- Molecular motors as paradigm
- Cost of temporal precision
- (Modelfree) inference from experimental data
- Efficiency of cellular information processing

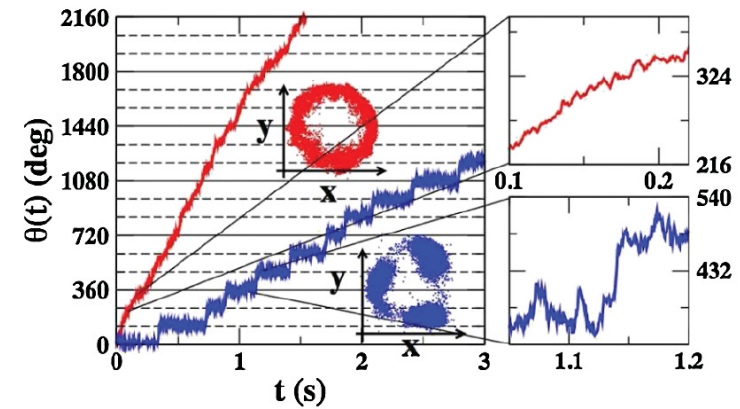
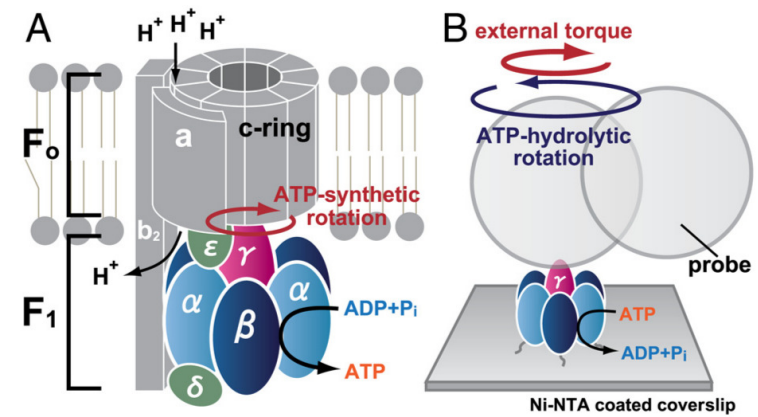
• From classical th'dynamics

to

stochastic th'dynamics



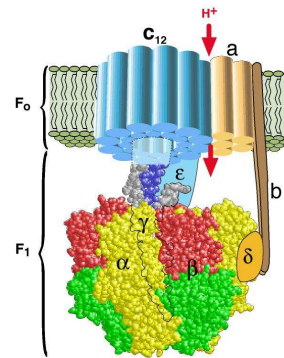
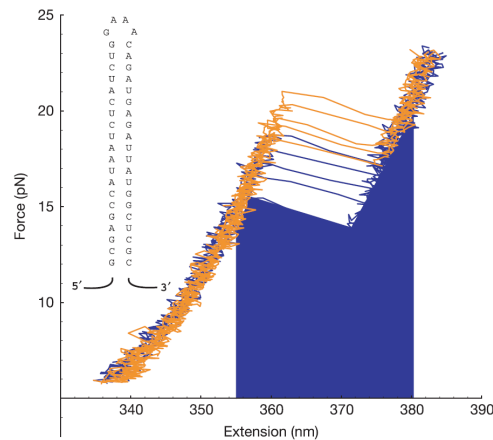
19th century steam engine



[K. Hayashi, ... H. Noji, PRL 104, 218103 (2010)]

21st century nano-engine: F₁ATP-ase

- **Stochastic thermodynamics** applies to systems for which
 - non-equilibrium is caused by mechanical or chemical forces
 - ambient solution provides a thermal bath of well-defined T and μ_i
 - fluctuations are relevant due to small numbers of involved molecules

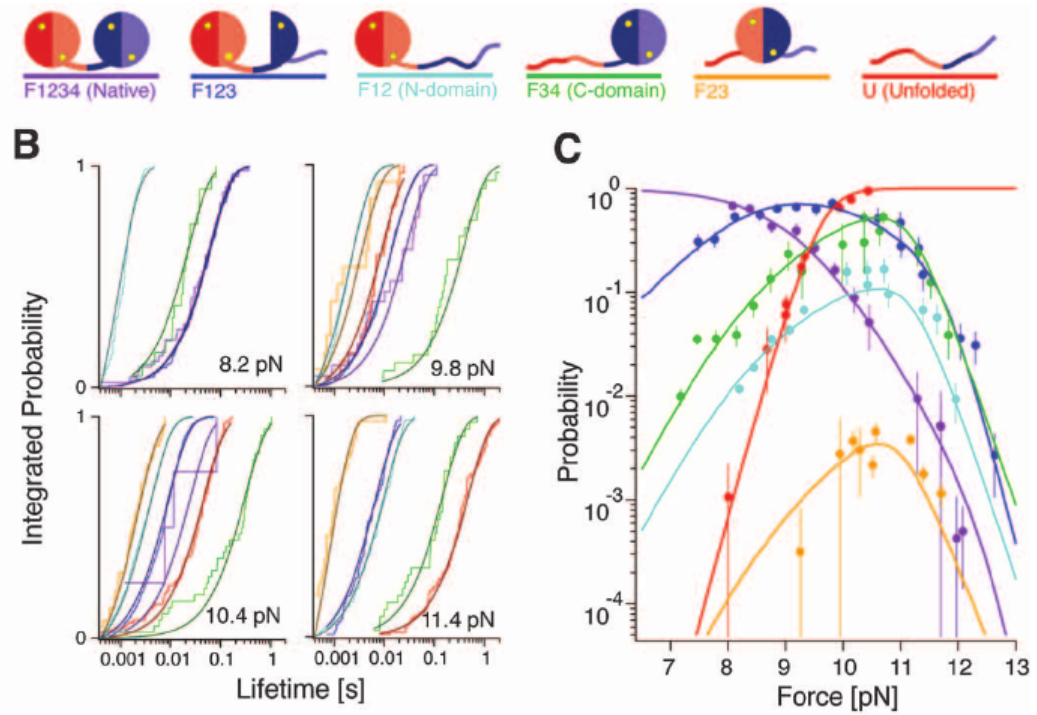
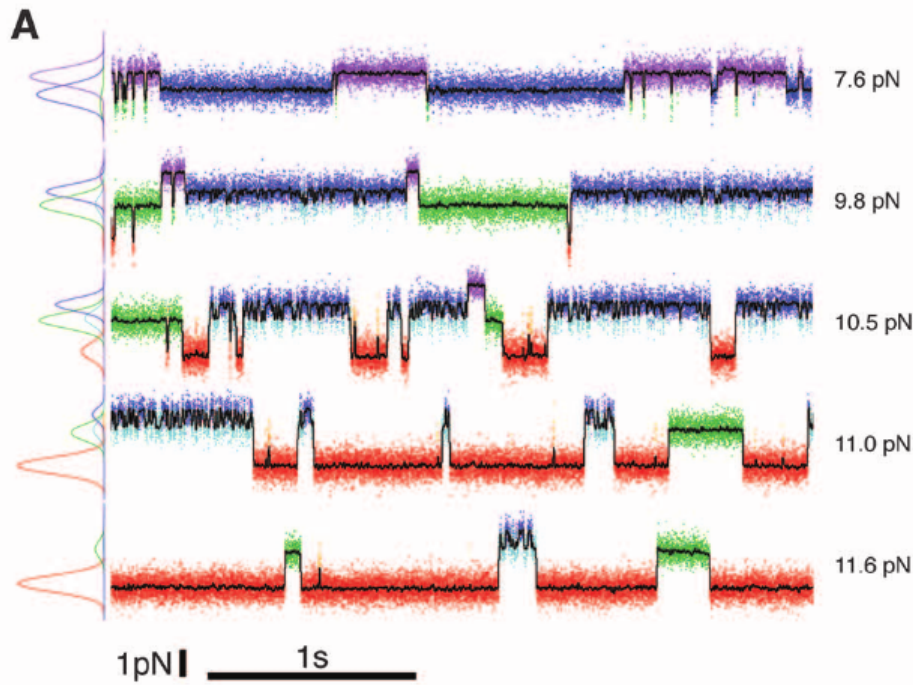
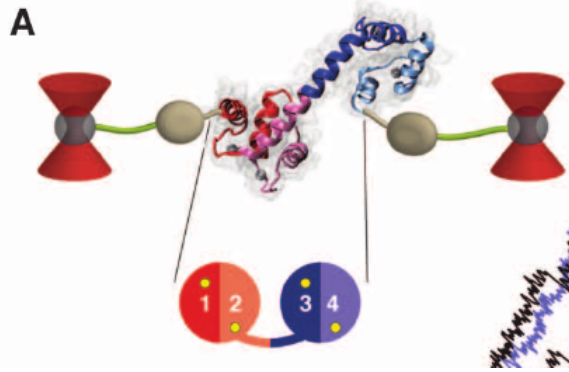


H. Wang and G. Oster (1998). Nature 396:279-282.

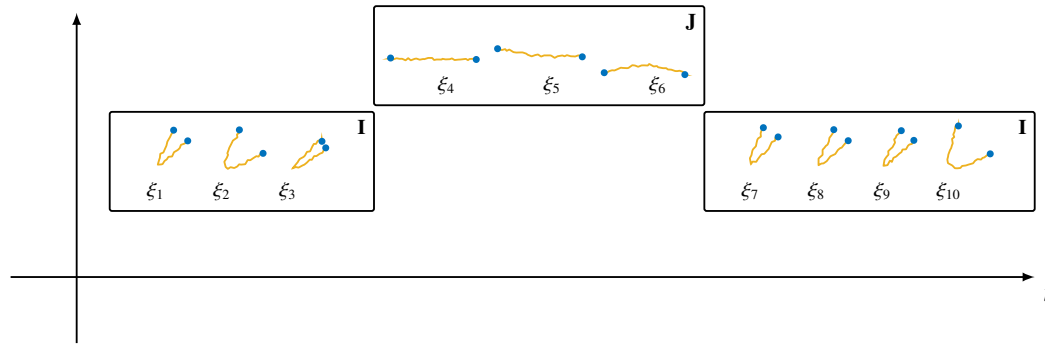
[Collin et al, Nature 437, 231, 2005]

- Main idea: Energy conservation (1^{st} law) and entropy production (2^{nd} law) along an individual stochastic trajectory
- Review: U.S., Rep. Prog. Phys. **75** 126001, 2012.

- Biomolecules in equilibrium: Meso-states of calmodulin [J. Stigler et al, Science **334** 512 (2011)]



- Thermodynamically consistent markovian dynamics on meso-states



- trajectory $I(t)$
- crucial time-scale separation:
 - * transitions between meso-states are slow
 - * transitions between the micro-states belonging to one meso-state are fast
- master equation

$$\partial_t P_I(t) = \sum_J [P_J(t)K_{JI} - P_I(t)K_{IJ}].$$

- local detailed balance condition on the rates $\{K_{IJ}\}$

$$\Rightarrow K_{IJ}/K_{JI} = P_J^e/P_I^e = \tau_J^e/\tau_I^e = \exp(-\beta\Delta_{IJ}F) = \exp(-\beta\Delta_{IJ}E + \Delta_{IJ}S)$$

- th'dyn potentials of meso-states operationally accessible from traj' data

- Thermodynamics along a trajectory $I(t)$ of biochemical/physical meso-states

[T. Schmiedl and U.S., J. Chem. Phys. 126, 044101 (2007)]

- internal energy $E(t) = E_{I(t)}$ becomes stochastic

- first law (Sekimoto 1998)

$$\Delta_{IJ}E \equiv E_J - E_I = -Q_{IJ}$$

- entropy of "system"

$$S_{\text{sys}}(t) \equiv S_{I(t)} - \ln[P_{I(t)}(t)]$$

- total entropy change in a transition from I to J at time t

$$\Delta_{IJ}S_{\text{tot}}(t) = \beta Q_{IJ} + \Delta_{IJ}S_{\text{sys}}(t) = \ln[P_I(t)K_{IJ}/P_J(t)K_{JI}]$$

- integral fluctuation theorem for total entropy production

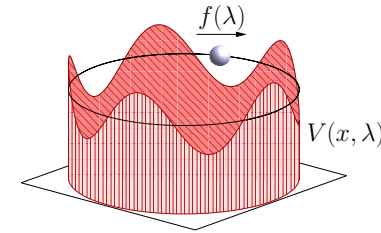
$$\langle \exp[-\Delta S_{\text{tot}}] \rangle = 1 \Rightarrow \langle \Delta S_{\text{tot}} \rangle \geq 0$$

- * any lengths t , any initial distribution $\{P_I^0\}$

[U.S., PRL 2005]

- second law on ensemble level (Schnakenberg 1976)

$$\langle \dot{S}_{\text{tot}}(t) \rangle \equiv \sum_{IJ} P_I(t) K_{IJ} \Delta_{IJ} S_{\text{tot}}(t) \geq 0$$



- Stochastic th'dynamics of NESS: Driven colloidal particle as paradigm

- Langevin dynamics $\dot{x} = \mu[-V'(x) + f] + \zeta$ with $\langle \zeta_1 \zeta_2 \rangle = 2\mu T \delta(t_2 - t_1)$

- first law [(Sekimoto, 1997)]:

$$dw = du + dq$$

- * applied work: $dw = f dx$

- * internal energy : $du = dV$

- * dissipated heat: $dq = dw - du = [-\partial_x V(x) + f]dx = T ds_b$

- total entropy as quantitative measure of broken time reversal symmetry

$$x(t) \rightarrow \tilde{x}(t) \equiv x(\mathcal{T} - t)$$

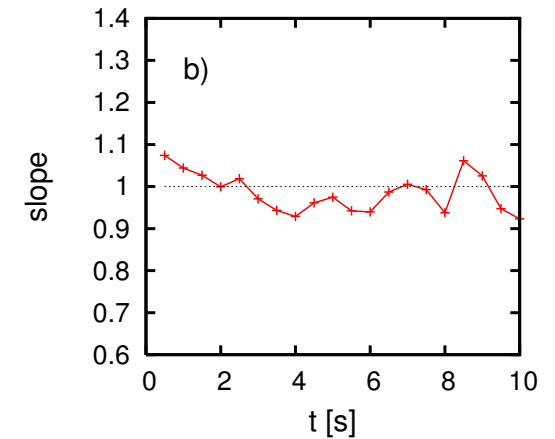
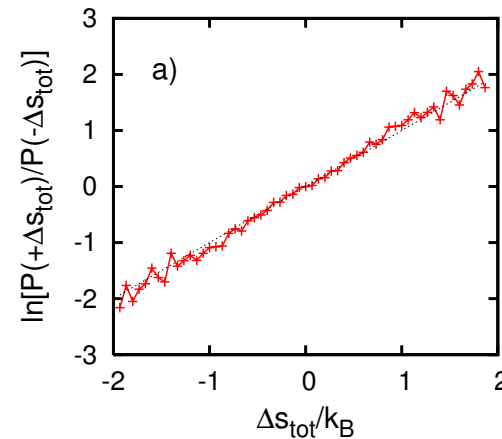
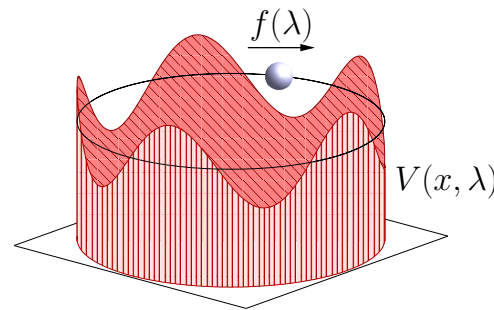
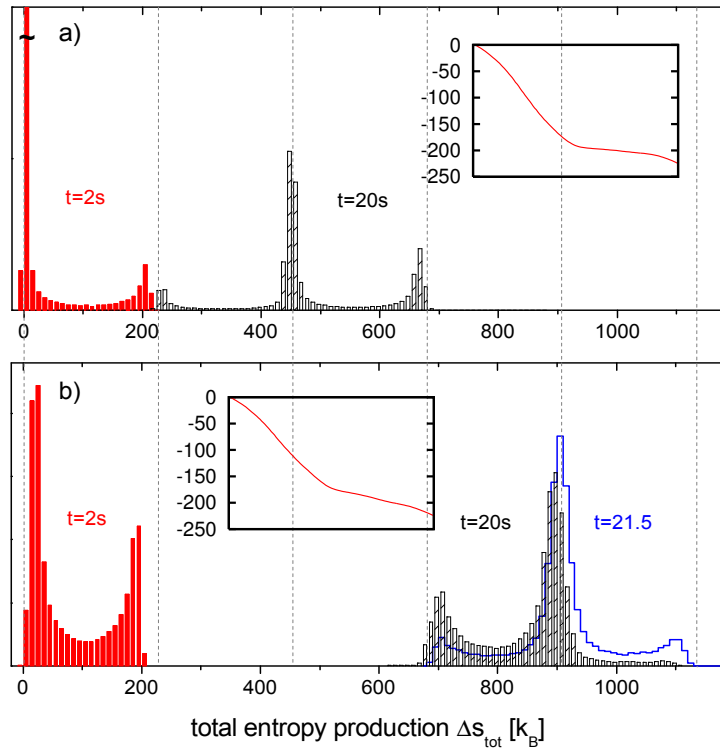
$$\Delta_{stot}[x(t)] \equiv \ln[p[x(t)]/p[\tilde{x}(t)]] = \Delta[-\ln p^s(x)] + q/T$$

- IFT for total entropy production $\langle \exp[-\Delta_{stot}] \rangle = 1 \Rightarrow \langle \Delta_{stot} \rangle \geq 0$ [U.S., PRL 2005]

- Fluctuation theorem $p(-\Delta s_{\text{tot}})/p(\Delta s_{\text{tot}}) = \exp(-\Delta s_{\text{tot}})$ in any NESS

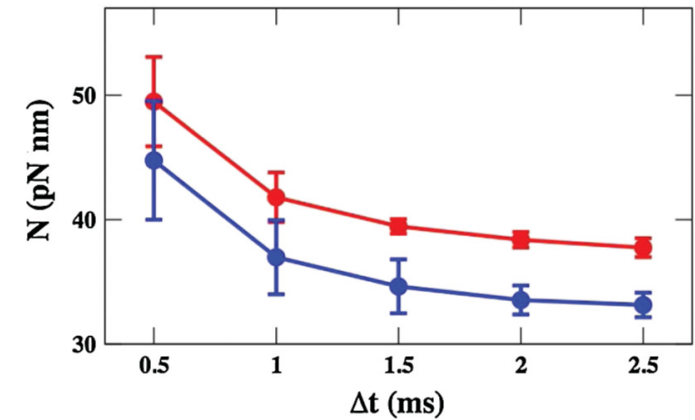
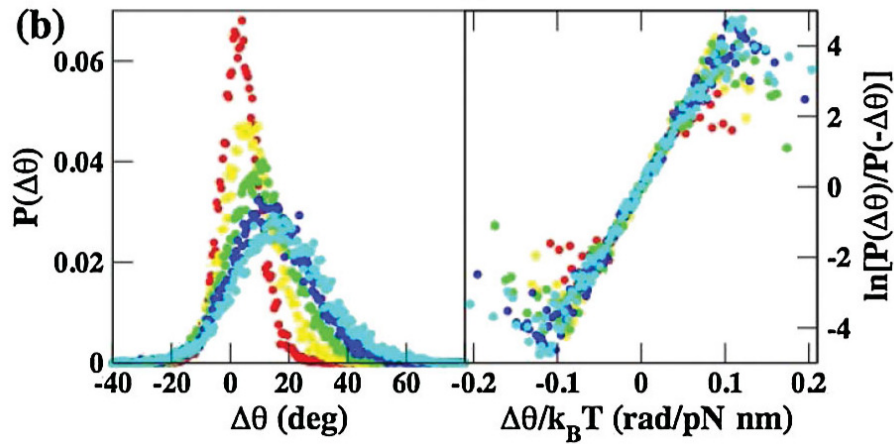
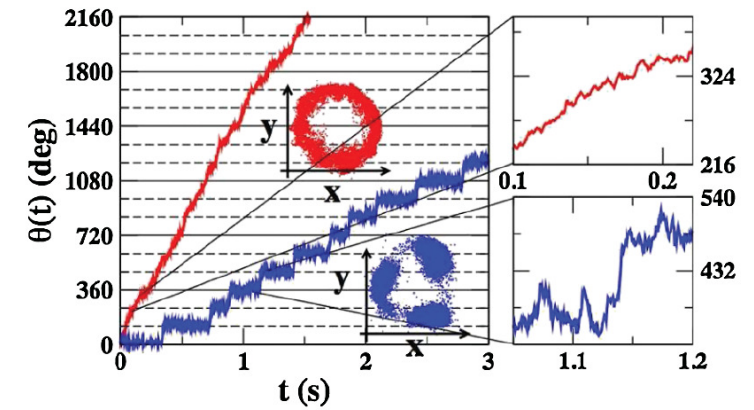
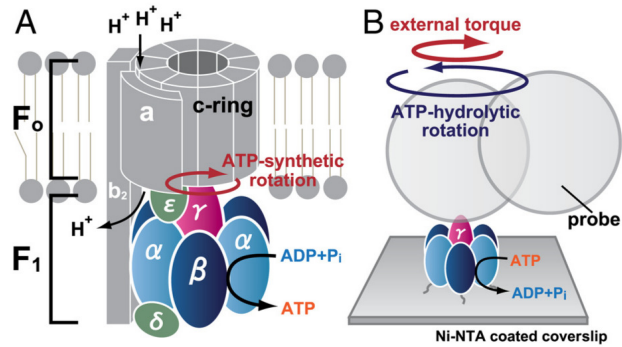
Evans et al (1993), Gallavotti & Cohen (1995), Kurchan (1998), Lebowitz & Spohn (1999), U.S. (2005)

- Experimental data [Speck, Blicke, Bechinger, U.S., EPL **79** 30002 (2007)]



FT-representation

- F1-ATPase and the fluctuation theorem [K. Hayashi et al, PRL 104, 218103 (2010)]



$$\Gamma \dot{\theta} = N + \zeta$$

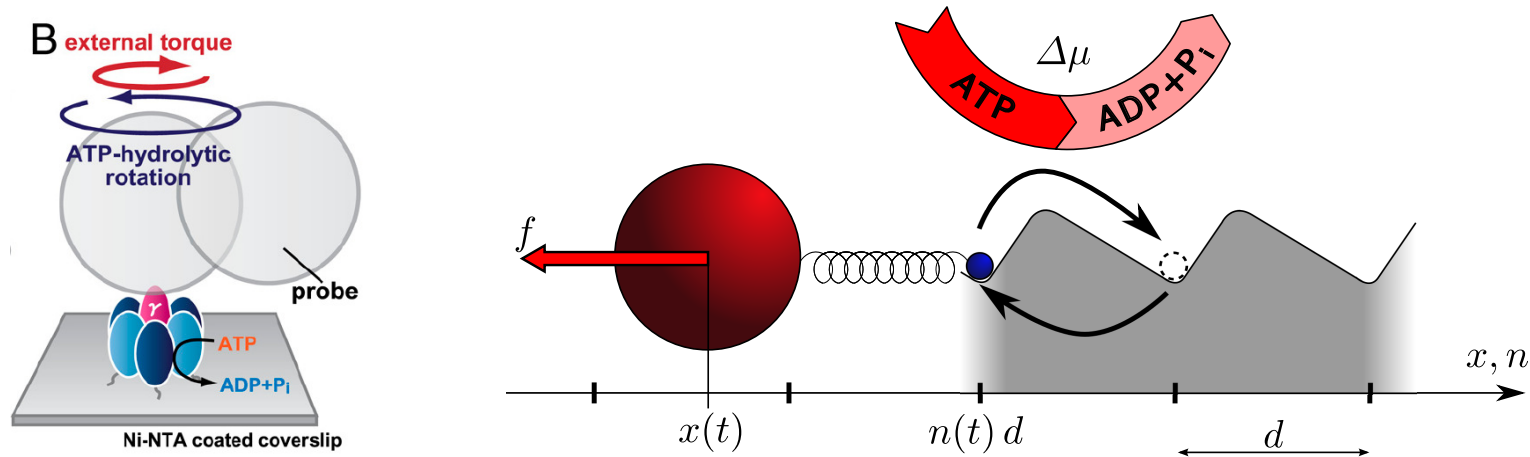
$$\Rightarrow \ln[p(\Delta\theta)/p(-\Delta\theta)] = N\Delta\theta/k_B T$$

independent of friction coefficient Γ

time-dependence?

torque N from $\Delta t \rightarrow \infty$?

- Hybrid model [E. Zimmermann and U.S., New J. Phys. 14, 103023, 2012]



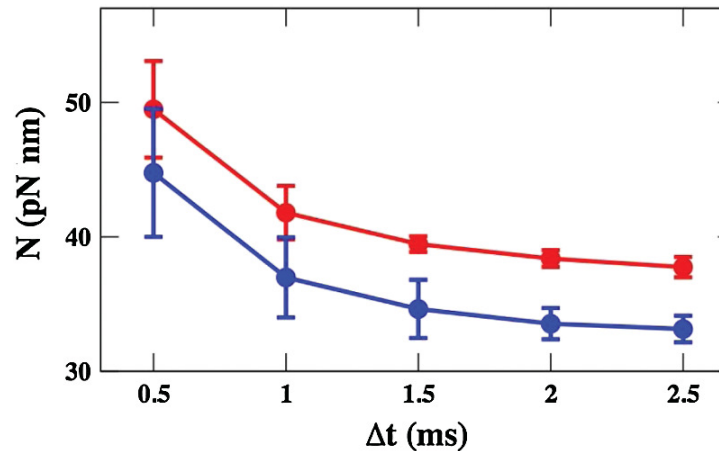
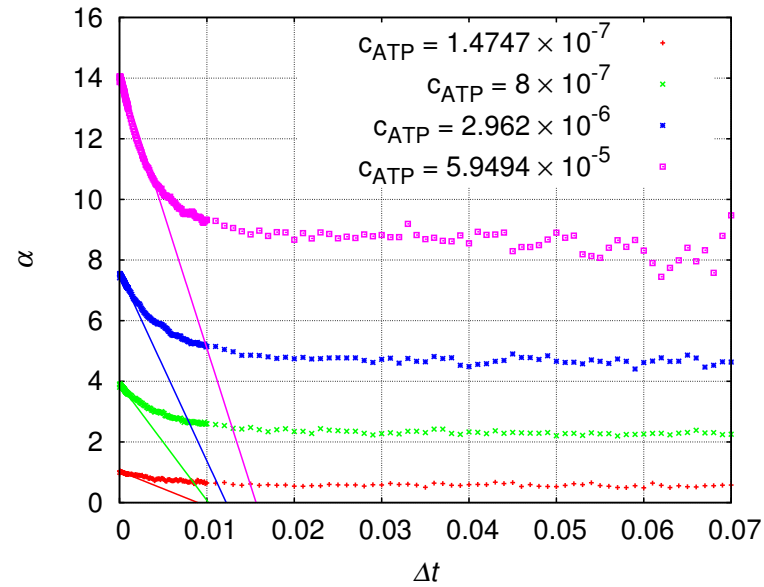
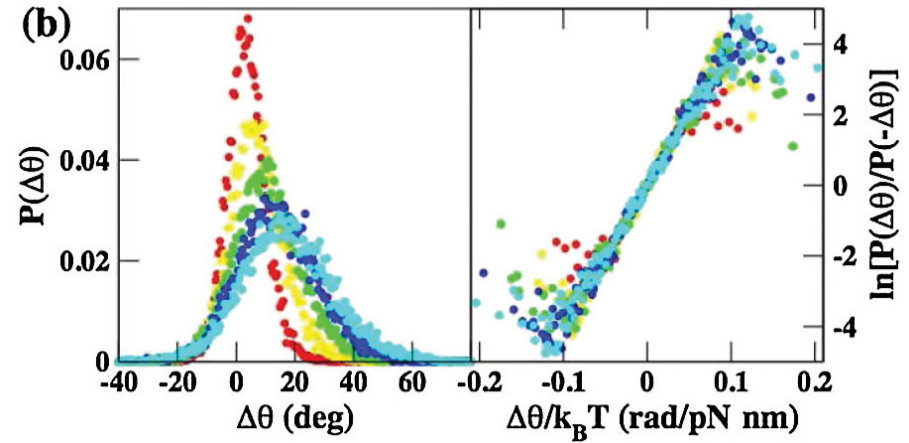
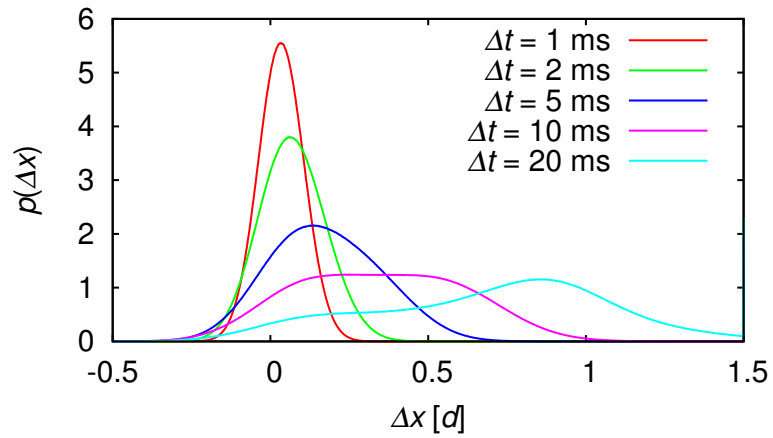
– probe particle

$$* \dot{x} = \mu(-\partial_y V(y) + f^{\text{ex}}) + \zeta \quad \text{with} \quad y(\tau) \equiv n(\tau) - x(\tau)$$

– motor

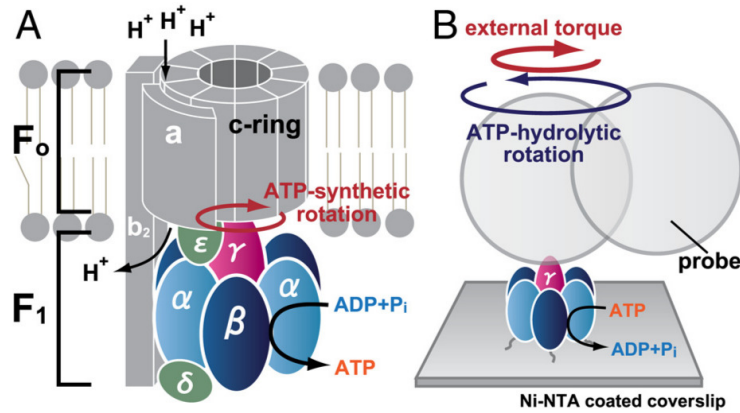
$$* w^+ / w^- = \exp[\Delta\mu - V(n + d, x) - V(n, x)]$$

● FT-slope from simulations vs experiment



$\Delta t \rightarrow 0$ limit yields average force/torque

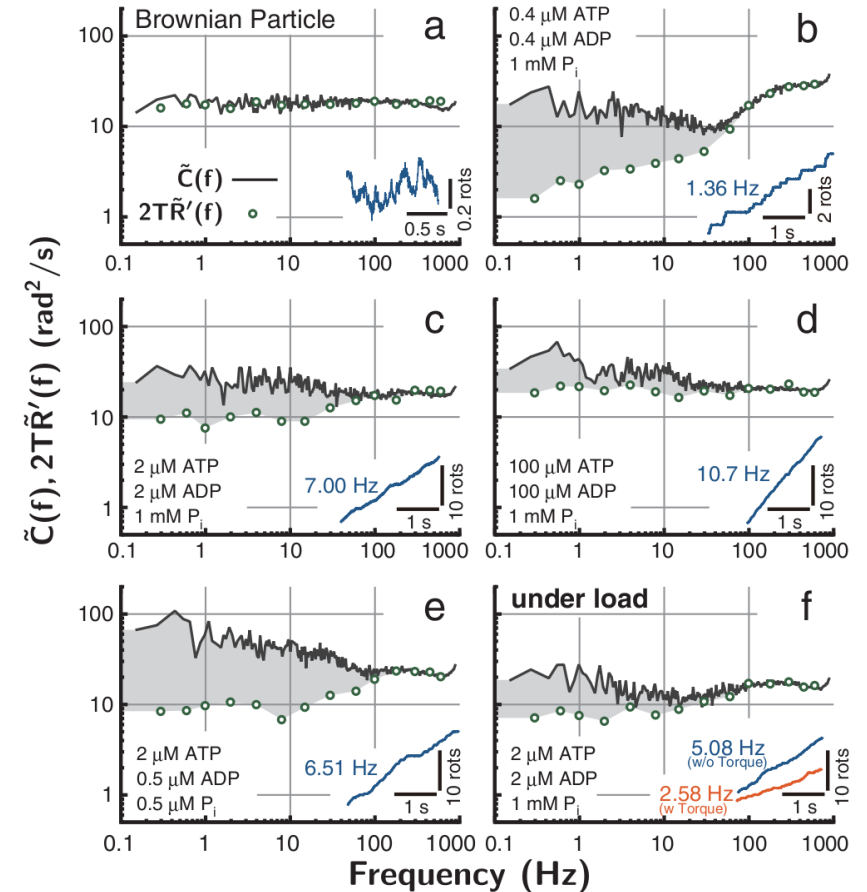
- Inferring the efficiency of a molecular motor [S. Toyabe et al, PRL 104, 198103 (2010)]



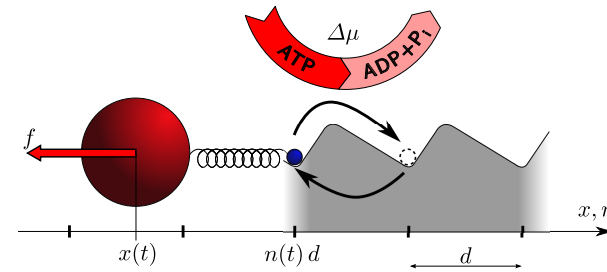
– Harada-Sasa relation [PRL 2006]

$$\mu \dot{Q}_P = v^2 + \int d\omega [C_{\dot{x}}(\omega) - 2k_B T \text{Re}R_{\dot{x}}(\omega)]$$

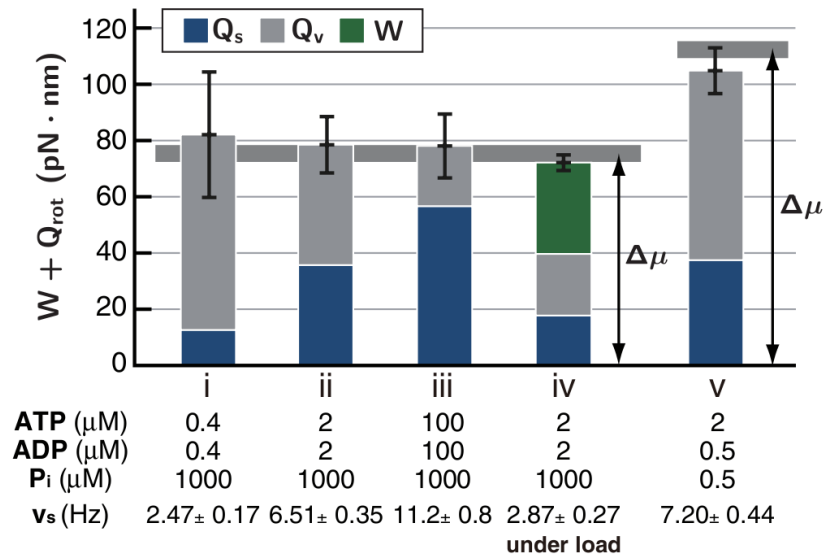
heat from "violation" of fluc-diss-theorem



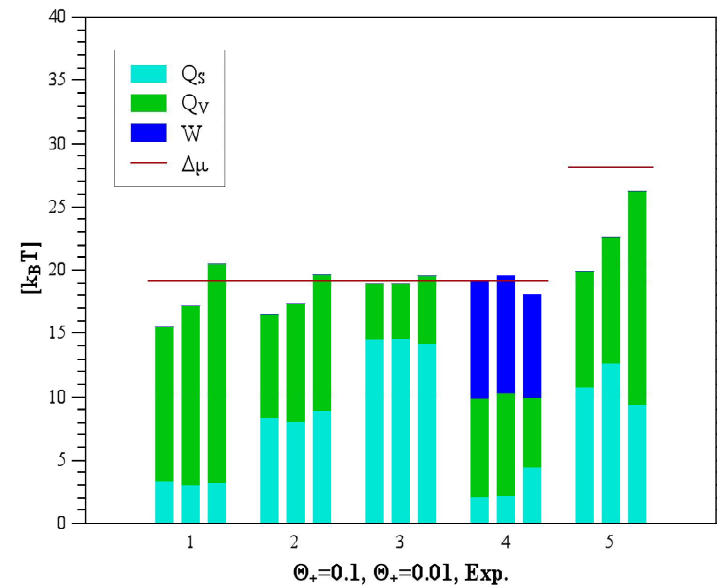
- Comparison:



experiment



theory



[S. Toyabe et al, PRL 104, 198103 (2010)]

[E. Zimmermann and US, NJP 2012]

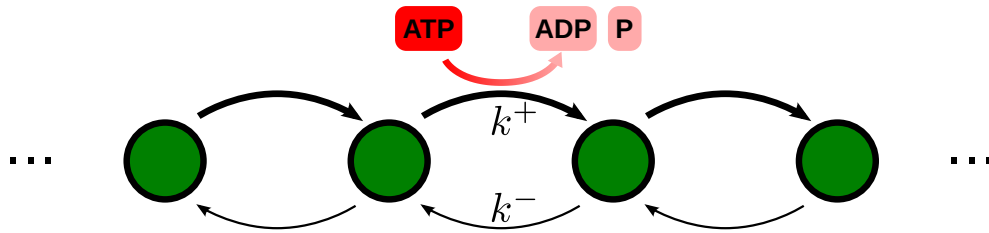
- Temporal precision in a finite temperature environment



at 300 K, a precision of 1 sec/day requires at least 6×10^{-11} J/day

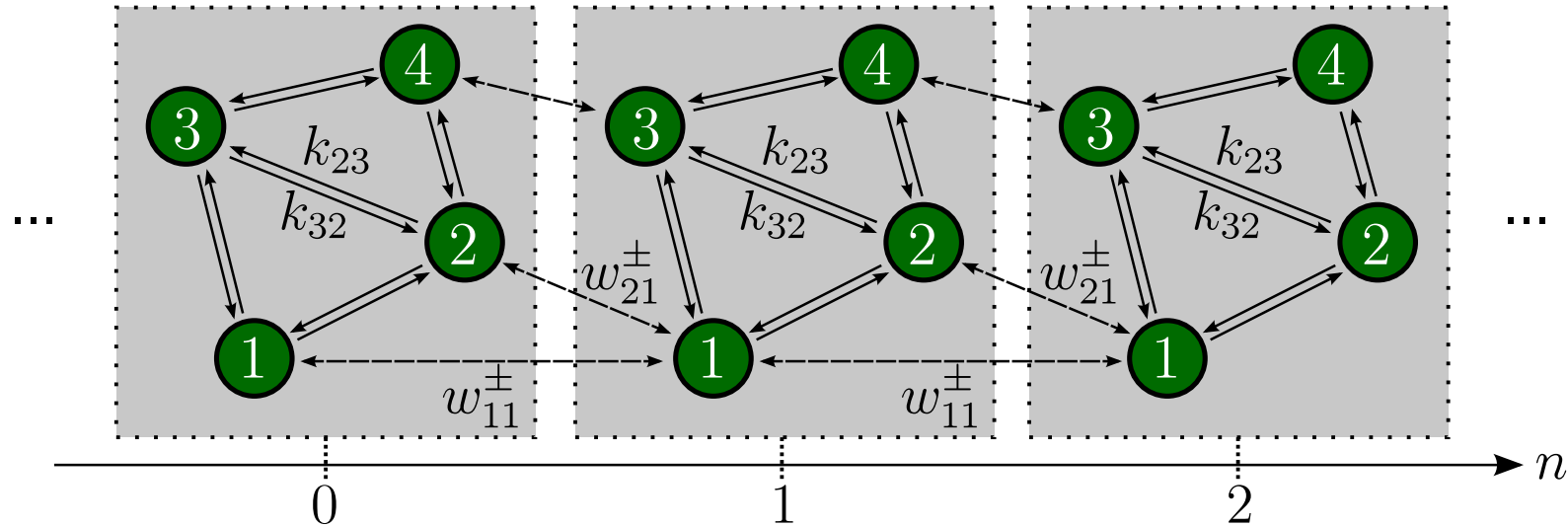
- Cost of running a simple clock: ARW

[AC Barato and US, Phys. Rev. Lett. 114, 158101, 2015]



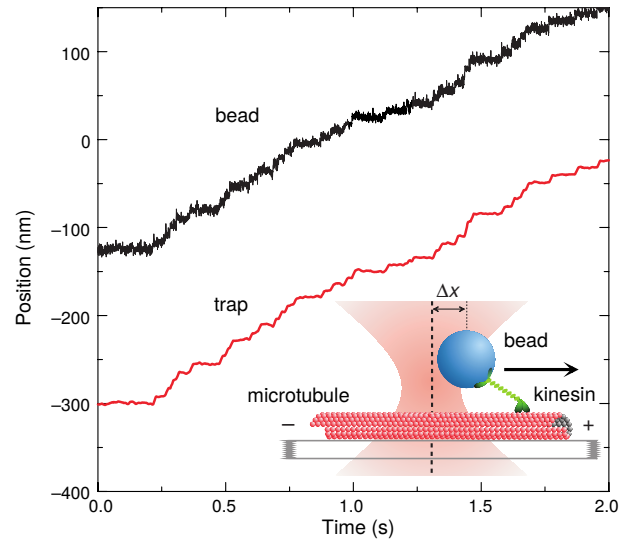
- output $n(t)$ with $\langle n \rangle = Jt = (k^+ - k^-)t$
- variance $\langle (n(t) - \langle n \rangle)^2 \rangle = 2Dt = (k^+ + k^-)t$
- uncertainty $\epsilon^2 \equiv \text{var}/\text{output}^2 = 2D/J^2t$
- th'dyn cost $\mathcal{C} = \sigma t = (k^+ - k^-) \ln(k^+/k^-)t$ with $\sigma \equiv$ rate of entropy production
- with affinity $\mathcal{A} = k_B T \ln(k^+/k^-) = \mu_{\text{ATP}} - \mu_{\text{ADP}} - \mu_{\text{P}}$
- $\boxed{\mathcal{C}\epsilon^2 = 2\sigma D/J^2 = \mathcal{A} \coth[\mathcal{A}/2k_B T] \geq 2k_B T}$ independent of run time t

- Thermodynamic uncertainty relation holds for general multicyclic processes
 AC Barato and US, Phys. Rev. Lett. 114, 158101, 2015; proof by Gingrich et al, PRL 2016



- $\mathcal{C} \geq 2k_B T / \epsilon^2$ for any th'dyn consistent process at finite T
- a precision of 1% costs at least 20.000 $k_B T$
- inevitable, universal cost of temporal precision (within stationary Markov processes)
- for any current $j = \sum_{ij} d_{ij} n_{ij}$ $\sigma \geq j^2 / D_j$

- Thermodynamic inference: Efficiency of a molecular motor



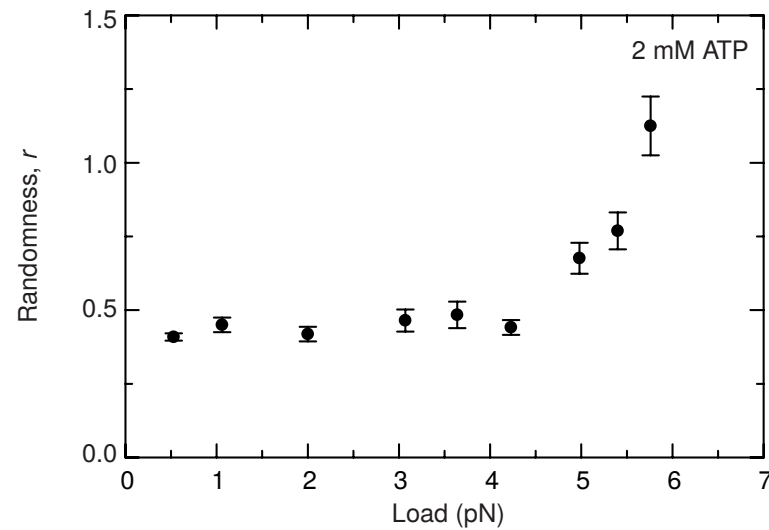
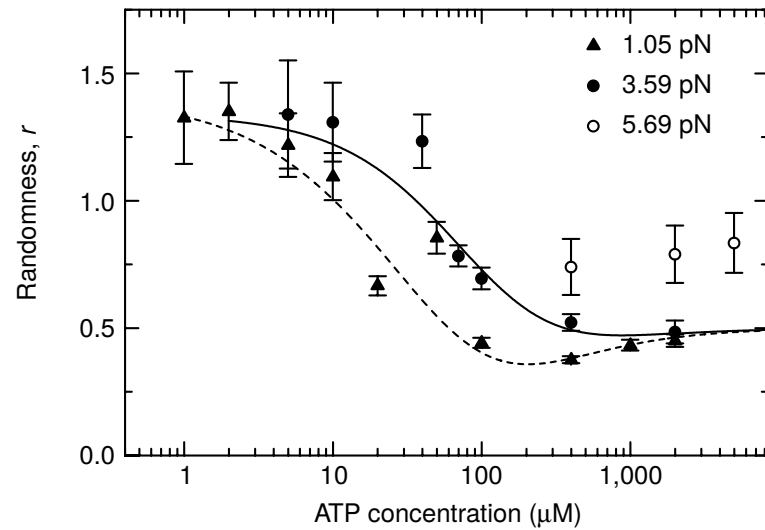
[Visscher et al, Nature, 1999]

– experimental data on

- * velocity v

- * diffusion constant D

- * randomness parameter $r \equiv 2D/vl$



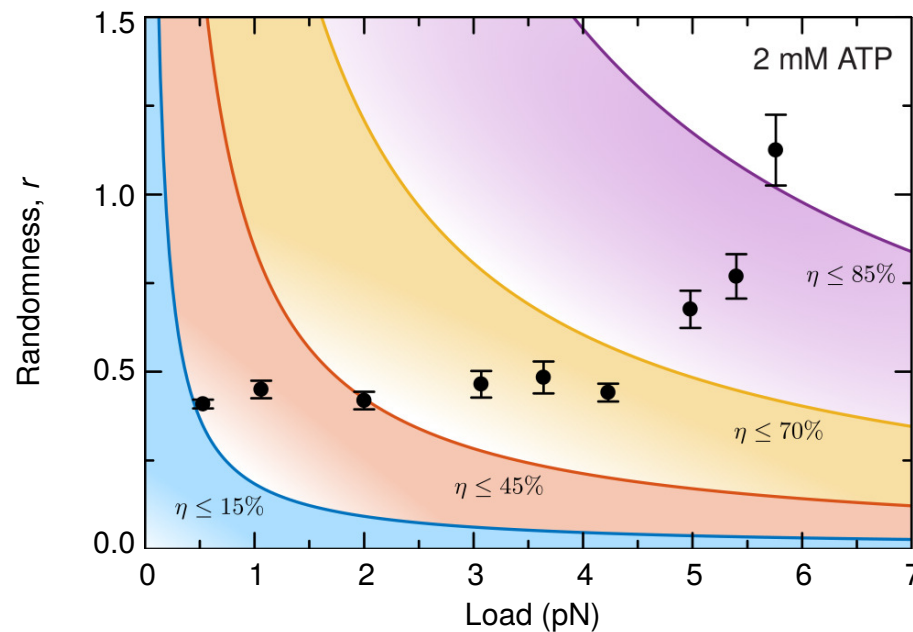
- Thermodynamic inference: Universal bound on the efficiency of molecular machines

P. Pietzonka, AC Barato, U.S., J Stat Mech, 124004, 2016; U.S., Physica A 504, 176, 2018

- efficiency

$$\eta \equiv \frac{P^{\text{out}}}{P^{\text{in}}} = \frac{fv}{\text{unknown}} = \frac{fv}{fv + \sigma} \leq \frac{1}{1 + vk_B T / (Df)}$$

- entropy production rate $\sigma = P^{\text{in}} - P^{\text{out}} = \text{"chem energy"} - fv \geq v^2/D$



- independent of the specific chemo-mechanical cycles and of $\Delta\mu$

- Th'dyn' uncertainty relation(s) beyond NESSs?

- periodic driving with period Ω

- * TUR "fails" in general

AC Barato and US, PRX 2016

- * for a symmetric protocol $k_{ij}(t) = k_{ij}(\mathcal{T} - t)$

$$\frac{\exp[\Delta S_{\text{per}}] - 1}{t_{\text{per}}} D_j > j^2$$

Proesmans and van den Broeck, EPL 2017

- * a number of "technical" bounds

Barato, Chetrite, Faggionato, Gabrielli, NJP 2018, JStatMech 2019

Proesmans and Horowitz, JStatMech 2019

Koyuk and U.S., JPA 2019

- * operationally accessible version

$$\sigma(\Omega) D(\Omega) \geq [J(\Omega) - \Omega \partial_{\Omega} J(\Omega)]^2$$

Koyuk and U.S., PRL 2019

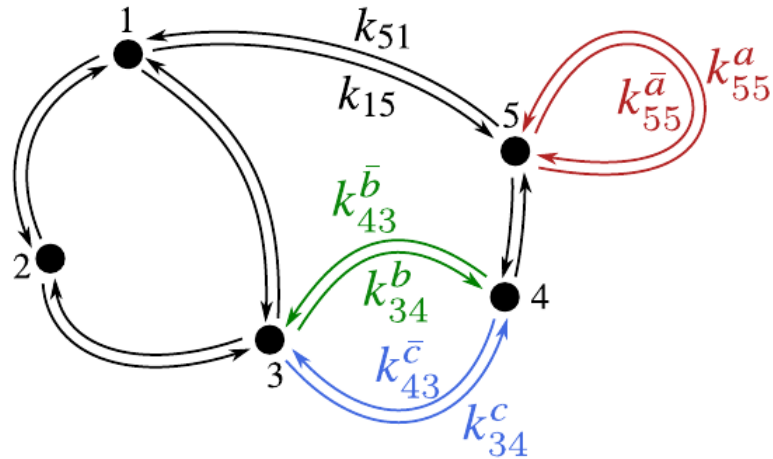
- relaxations towards equilibrium or a NESS

Dechant and Sasa, JStatMech 2018, Liu, Gong and Ueda, arxiv 2019

- A few further generalizations and ramifications
 - underdamped Langevin dynamics
 - * TUR "fails" for a finite time \mathcal{T}
 - * conjectured but not proven for $\mathcal{T} \rightarrow \infty$ yet
 - role of a magnetic field Brandner et al PRL 2018, Chun et al, PRE 2019
 - "generalized" TURs from the fluctuation theorem Hasegawa and Van Vu, PRL 2019, Timpanaro et al PRL2019
 - optimal current? (or even the optimal observable?) Polettoni et al PRE2016, Busiello and Pigolotti, PRE 2019
Falasco et al arxiv 2019, Manikandan et al PRL 2020
 - bounds on variance of time-symmetric observables Maes PRL 2017, Nardini and Touchette EPJB 2018, Terlizzi and Baiesi, JPA 2019
 - bounds on variance of first-passage times Garrahan PRE 2017, Gingrich and Horowitz, PRL 2017
 - open quantum systems

- Generalization to time-dependent driving

[T Koyuk and U.S., arXiv 2005.02312]

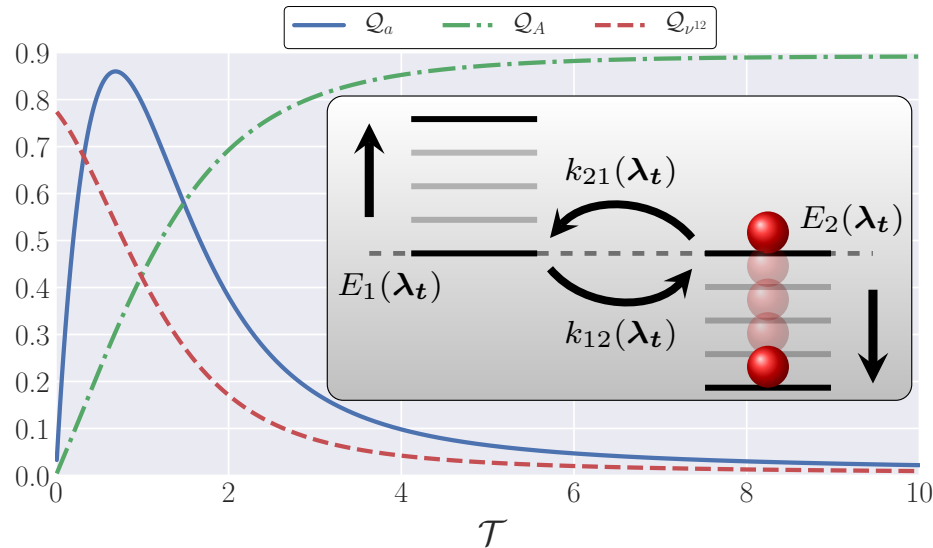


- network with rates $k_{ij}(\lambda)$ that depend on a driving protocol $\lambda = \lambda(t)$
- protocol $\lambda(t) = \lambda(vt)$ depends on an experimentally controllable speed parameter v
- system is driven for a total (observation) time $t = \mathcal{T}$
- mean current $J(\mathcal{T}, v)$

$$[J(\mathcal{T}, v) + \Delta J(\mathcal{T}, v)]^2 / D_J(\mathcal{T}, v) \leq \sigma(\mathcal{T}, v) \quad \text{with} \quad \Delta \equiv \mathcal{T} \partial_{\mathcal{T}} - v \partial_v$$

- Generalization to state observables

- Example: driven two-state system $k_{12}(vt)$

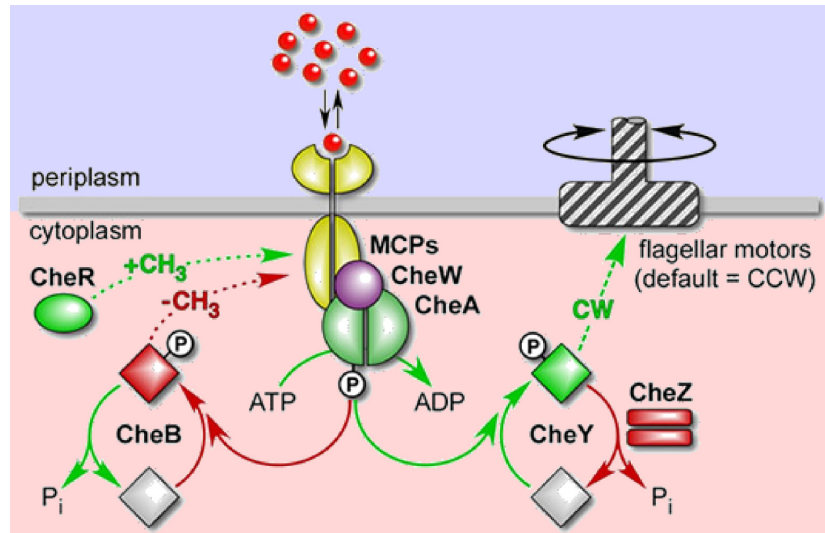


- observable $a(\mathcal{T}, v) \equiv \delta_{n(\mathcal{T})2}$ at final time [blue]
- observable $A(\mathcal{T}, v) \equiv \tau_2/\mathcal{T}$ total time spent in state 2 [green]

$$[\Delta X(\mathcal{T}, v)]^2/D_X(\mathcal{T}, v) \leq \sigma(\mathcal{T}, v) \text{ for } X(\mathcal{T}, v) \in \{a(\mathcal{T}, v), A(\mathcal{T}, v)\} \text{ with } \Delta \equiv \mathcal{T}\partial_{\mathcal{T}} - v\partial_v$$

- current $J(\mathcal{T}, v) = [n_{12}(\mathcal{T}, v) - n_{21}(\mathcal{T}, v)]/\mathcal{T}$ [red]

- Information and entropy production in cellular sensing



– Lan et al, Nature Phys. 2012;
Mehta and Schwab PNAS 2012

Information about an external process $\{X(t)\}$ (conc' of a nutrient) is recorded by an **internal variable** $\{Y(t)\}$ (conc' of phosph'd protein)

Q: Is a rate of information (which one?) acquired about $\{X(t)\}$ related to the thermodynamic cost σ of maintaining the sensory network?

[AC Barato, D Hartich, U.S., PRE 87, 042104, 2013]

[cf: Horowitz & Sandberg NJP 2014, Lang, Fisher, Mora & Mehta PRL 2014,

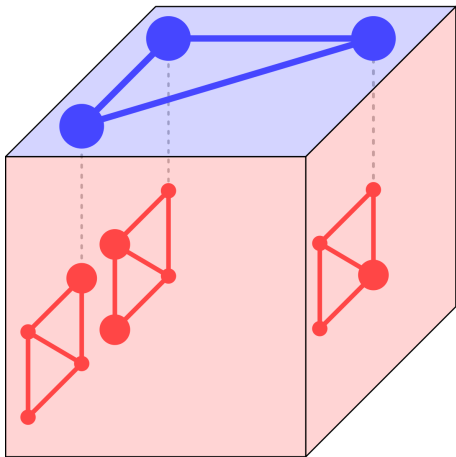
Bo, Del Giudice & Celani JSM 2015,...]

- Framework: Stochastic thermodynamics of bipartite systems

[D. Hartich, AC Barato, U.S., J Stat Mech, P02016, 2014]

- (external) x -jumps $w^{\alpha\beta}$

independent of y



- (internal) y -jumps w_{ij}^{α}

affected by x

- thermodynamic entropy production

$$\sigma = \sigma_x + \sigma_y \geq 0$$

- conditional Shannon entropy

$$H[x|y] = - \sum_{i,\alpha} P_i P(\alpha|i) \ln P(\alpha|i)$$

- learning rate

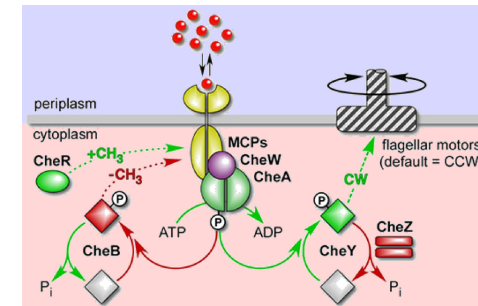
$$l_y \equiv - \frac{d}{dt} \Big|_{y\text{-jumps}} H[x|y] \leq \sigma_y \leq \sigma$$

- efficiency of learning

$$\eta \equiv l_y / \sigma_y \leq 1$$

- rate of mutual information not bounded by σ_y !

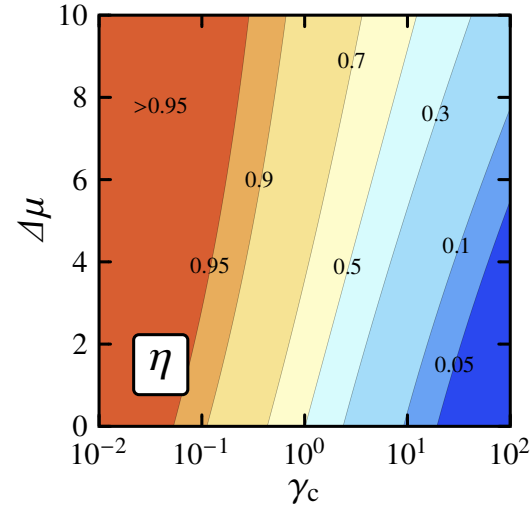
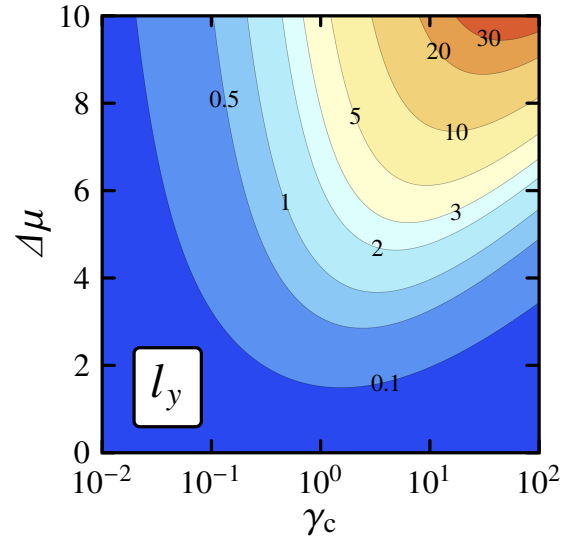
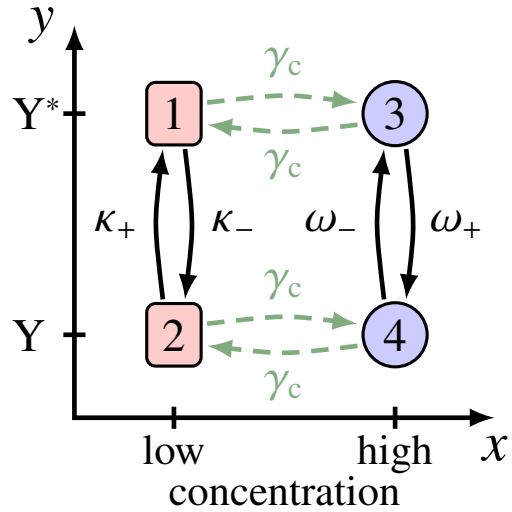
- Efficiency of cellular sensing



[AC Barato, D Hartich, U.S., NJP **16** 103024, 2014]

Model	ext conc	ligand	activity of rec'r	methyl	Che Y/Y^*
minimal 4-state	< - - - - -	X	- - - - - >	-	Y
eq receptor	X	Y_b	Y_a	-	-
coarse-grained	X	fast	Y_a	-	-
with adaptation	X	fast	Y_a	Y_m	-

- Minimal four state model [AC Barato, D Hartich, U.S., NJP **16** 103024, 2014]



$$\Delta\mu = k_B T \ln(\kappa_+ \omega_+ / \kappa_- \omega_-)$$

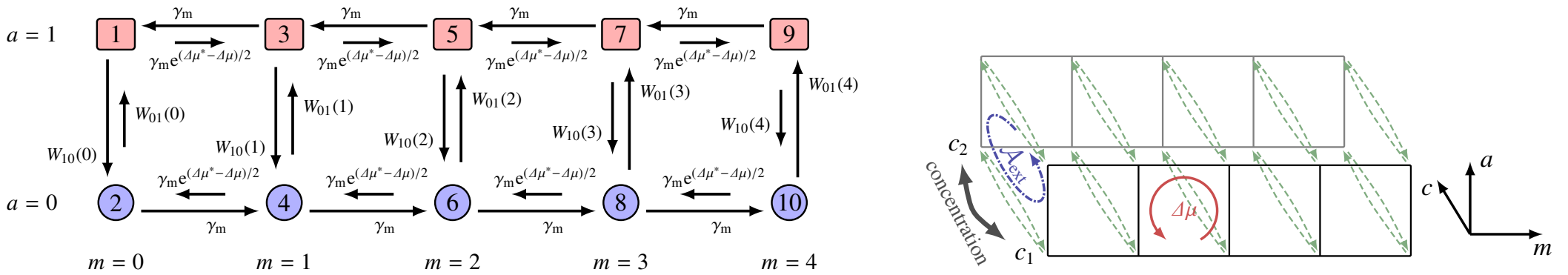
$$\sigma = J \Delta\mu$$

$$l_y = J \ln[P_1 P_4 / P_2 P_3]$$

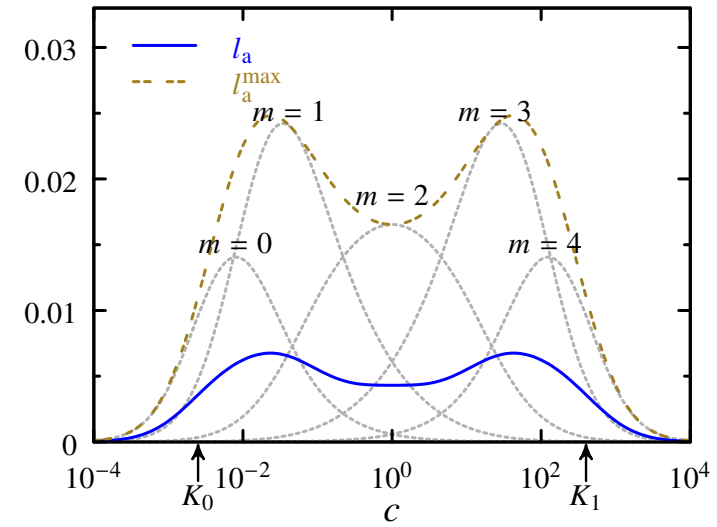
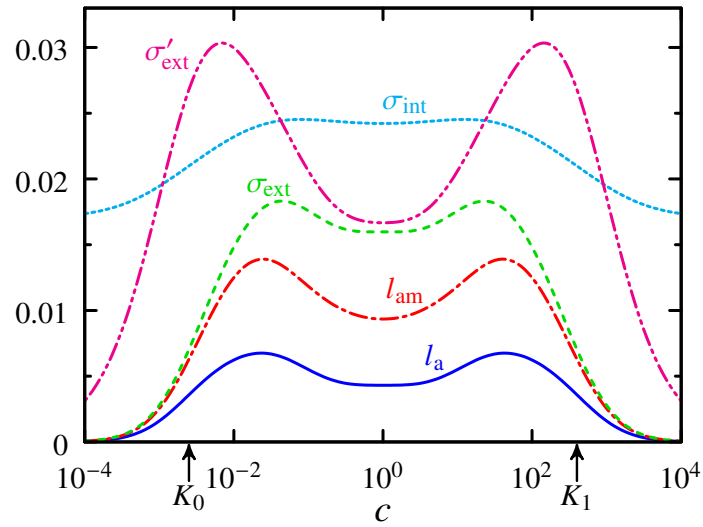
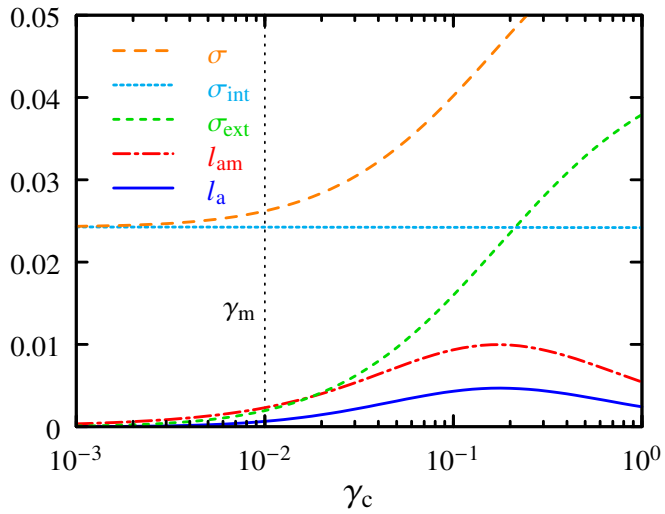
- l_y increases with $\Delta\mu$
- adiabatic case : $\eta \rightarrow 1$
- fast external changes: $\eta \rightarrow 0$

• Model with adaptation

[AC Barato, D Hartich, U.S., NJP **16** 103024] 2014



increasing methylation level \longrightarrow



- Stochastic thermodynamics as
 - . universal, thermodynamically consistent, quantitative framework
- Universal inequalities (bounds) for thermodynamic inference
 - thermodynamic uncertainty relation provides constraints on
 - * ... cost of any process with given precision
 - * ... efficiency of any molecular motor (complex)
 - generalization to arbitrary time-dependent driving
- Bipartite approach to information processing in sensing
- Acknowledgments

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David Hartich (MPI-BPC Göttingen), Timur Koyuk (Stuttgart)
- Review U.S., Annu. Rev. Cond. Mat. Phys. **10** 171 (2019)