

Automorphic representations and L -functions

Introductory talk, part II

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Theorem (Langlands 1970)

Let π be a cuspidal representation of $\mathrm{PGL}_2(\mathbb{A})$.

Assume that $L^S(s, \pi, \mathrm{Sym}^n)$ is holomorphic for $\Re s > 1$ for all $n \geq 1$.

Then $|a_p| \leq 2$ for all p .

Recall that a_p is the eigenvalue of the normalized Hecke operator, and writing $a_p = \alpha_p + \alpha_p^{-1}$

$$L(s, \pi, \mathrm{Sym}^n) = \prod_p \prod_{i=0}^n (1 - p^{-s} \alpha_p^{2i-n})^{-1}$$

Proof.

As representations of SL_2 we have

$$\mathrm{Sym}^n \otimes \mathrm{Sym}^n = \bigoplus_{i=0}^n \mathrm{Sym}^{2i}.$$

It follows that

$$L^S(s, \pi, \mathrm{Sym}^n \otimes \mathrm{Sym}^n) = \prod_{i=0}^n L^S(s, \pi, \mathrm{Sym}^{2i})$$

and hence, $L^S(s, \pi, \mathrm{Sym}^n \otimes \mathrm{Sym}^n)$ is holomorphic for $\Re s > 1$.

On the other hand, $L^S(s, \pi, \mathrm{Sym}^n \otimes \mathrm{Sym}^n)$ has non-negative coefficients.

Therefore, by **Landau's Lemma**, the Euler product converges absolutely.

In particular, for every p , $\det(1 - p^{-s}(\mathrm{Sym}^n \otimes \mathrm{Sym}^n)((\begin{smallmatrix} \alpha_p & \\ & \alpha_p^{-1} \end{smallmatrix})))^{-1}$ is holomorphic for $\Re s > 1$, where $a_p = \alpha_p + \alpha_p^{-1}$.

In other words, the eigenvalues of $(\mathrm{Sym}^n \otimes \mathrm{Sym}^n)((\begin{smallmatrix} \alpha_p & \\ & \alpha_p^{-1} \end{smallmatrix}))$ are $\leq p$ in absolute value. This means that $|\alpha_p|^{\pm 2n} \leq p$.

Since this is true for all n , we conclude that $|\alpha_p| = 1$. □

Eisenstein series

Let $\mathbb{H} = \{x + iy \mid y > 0\}$ be the hyperbolic upper half plane. The group $SL_2(\mathbb{R})$ acts on \mathbb{H} simply transitively by Möbius transformations.

The stabilizer of the point i is $K = SO(2)$.

The function $\text{Im } z$ is invariant under $N = \left\{ \begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix} \mid x \in \mathbb{R} \right\}$.

Let $\Gamma = SL_2(\mathbb{Z})$, $\Gamma_\infty = N \cap \Gamma$.

The simplest Eisenstein series (Maass, 1949) is

$$E(z; s) = \sum_{\gamma \in \Gamma_\infty \backslash \Gamma} \text{Im}(\gamma z)^{s+\frac{1}{2}} = \sum_{(m,n) \in \mathbb{Z}^2 \mid \gcd(m,n)=1} \frac{(\text{Im } z)^{s+\frac{1}{2}}}{|mz + n|^{2s+1}}$$

The series converges absolutely for $\Re s > 1$ and defines a function on $\Gamma \backslash \mathbb{H}$.

The constant term $\int_{\Gamma_\infty \backslash N} E(nz; s) \, dn$ is given by

$$(\text{Im } z)^{s+\frac{1}{2}} + \frac{\zeta^*(2s)}{\zeta^*(2s+1)} (\text{Im } z)^{-s+\frac{1}{2}}$$

In the 1960's **Langlands** computed the constant term of Eisenstein series induced from a cuspidal representation π of $M(\mathbb{A})$ where M is the Levi part of a maximal parabolic subgroup of G defined over \mathbb{Q} .

The main term is

$$\prod_{i=1}^m \frac{L(is, \pi, r_i)}{L(is + 1, \pi, r_i)}$$

where r_i are certain representations of the L -group of M .

This computation led Langlands to the notion of the **L -group**, the general notion of an automorphic L -function and finally to the **functoriality conjectures** (roughly: “every L -function is a standard L -function of GL_n ”).

Examples for r_i

- $G = \mathrm{GL}_{n_1+n_2}$, $M = \mathrm{GL}_{n_1} \times \mathrm{GL}_{n_2}$, $r_1 = \mathrm{St}_{n_1} \otimes \mathrm{St}_{n_2}$ (Rankin–Selberg convolution).
- G is an orthogonal group of rank n , $M = \mathrm{GL}_n$

$$r_1 = \begin{cases} \mathrm{Sym}^2 & \text{if } G = \mathrm{SO}(2n+1) \\ \wedge^2 & \text{if } G = \mathrm{SO}(2n) \end{cases}$$

- G – a classical group of rank n , $M = G' \times \mathrm{GL}_{n-k}$ where G' is a classical group (of the same type) of rank k , $r_1 = \mathrm{Can}_k \times \mathrm{St}_{n-k}$ where Can_k is the “canonical” representation of ${}^L G'$.
- $G = G_2$, $M = \mathrm{GL}_2$, $r_1 = \mathrm{Sym}^3$.
- $G = E_8$, $M = \mathrm{GL}_8$, $r_1 = \wedge^3$.
- $G = E_8$, $M \approx \mathrm{GL}_2 \times \mathrm{GL}_3 \times \mathrm{GL}_5$, $r_1 = \mathrm{St}_2 \times \mathrm{St}_3 \times \mathrm{St}_5$.

The theory of Eisenstein series was developed by Selberg and Langlands (1950's-60's) in connection with the spectral theory of $L^2(G(\mathbb{Q})\backslash G(\mathbb{A}))$.

Thanks to Bernstein, there is now a simple proof of the meromorphic continuation of Eisenstein series.

This gives the meromorphic continuation of $L(s, \pi, r_i)$, $i = 1, \dots, m$.

Unfortunately, by itself it gives neither the functional equation nor information on the poles.

In the important case where G is quasi-split and π is **generic**, **Shahidi's** work (1978–2000) gives a functional equation and *some* information about the poles.

This eventually leads to Sym^4 functoriality from GL_2 to GL_5 as well as functoriality from generic representations of classical groups to GL_N . (**Cogdell–Kim–Piatetskii–Shapiro–Shahidi**, 2002)

Basic Eisenstein series

Consider the right action of GL_n on row vectors with n entries.

For any $\Phi \in \mathcal{S}(\mathbb{A}^n)$ the Eisenstein series

$$\mathcal{E}(\Phi, g, s) = \int_{\mathbb{Q}^* \backslash \mathbb{I}} \sum_{\xi \in \mathbb{Q}^n \setminus \{0\}} \Phi(t\xi g) |\det(tg)|^{s+\frac{1}{2}} dt, \quad g \in [GL_n]$$

converges for $\Re s > \frac{1}{2}$ and can be meromorphically continued to \mathbb{C} .

Riemann (followed up by Hecke and Tate)

$$\begin{aligned} \mathcal{E}(\Phi, g, s) &= \int_{|t| \geq 1} \sum_{\xi \in \mathbb{Q}^n \setminus \{0\}} \Phi(t\xi g) |\det(tg)|^{s+\frac{1}{2}} dt - \frac{\Phi(0)}{n(s + \frac{1}{2})} \\ &+ \int_{|t| \geq 1} \sum_{\xi \in \mathbb{Q}^n \setminus \{0\}} \hat{\Phi}(t\xi g^*) |\det(tg^*)|^{\frac{1}{2}-s} dt + \frac{\hat{\Phi}(0)}{n(s - \frac{1}{2})} = \mathcal{E}(\hat{\Phi}, g^*, -s) \end{aligned}$$

$$\text{where } \hat{\Phi}(x) = \int_{\mathbb{A}^n} \Phi(y) \psi([x, y]) dy, \quad [xg, yg^*] = [x, y].$$

Rankin–Selberg and Jacquet–Shalika integrals

Let $\varphi_i \in V_{\pi_i}$ be cusp forms of $GL_n(\mathbb{A})$. Then, (Rankin, 1939; Selberg, 1940; Jacquet–Piatetski-Shapiro–Shalika 1980s)

$$\int_{[GL_n]} \varphi_1(g)\varphi_2(g)\mathcal{E}_{\Phi}(g, s) dg \stackrel{=} L(s + \frac{1}{2}, \pi_1 \times \pi_2)$$

(i.e., up to local integrals).

Let $\varphi \in V_{\pi}$ be a cusp form on $GL_{2n}(\mathbb{A})$. Then, (Jacquet–Shalika, 1990)

$$\int_{[GL_n]} \int_{[Mat_n]} \varphi\left(\begin{pmatrix} I_n & X \\ & I_n \end{pmatrix} \begin{pmatrix} g & \\ & g \end{pmatrix}\right) \psi(\text{tr } X) dX \mathcal{E}_{\Phi}(g, s) dg \stackrel{=} L(s + \frac{1}{2}, \pi, \wedge^2)$$

Many other integrals involving Eisenstein series have been discovered over the years. (The uncontested champion here is **David Ginzburg**.)

The conceptual framework is still far from understood.

Godement–Jacquet integral

Now consider the space Mat_n of $n \times n$ -matrices.

Left and right matrix multiplication gives rise to the tensor product representation

$$\iota : \text{GL}_n \times \text{GL}_n \rightarrow \text{GL}_{n^2}$$

The pullback of $\mathcal{E}(\Phi, s)$, $\Phi \in \mathcal{S}(\text{Mat}_n(\mathbb{A}))$ to $[\text{GL}_n \times \text{GL}_n]$ is a **reproducing kernel** for the standard L -function in the sense that

$$(\mathcal{E}(\Phi, \iota(g, h), s))^{\text{cusp}} = \sum_{\pi} \sum_{\varphi} L(ns + \frac{1}{2}, \pi) \delta(f_{\Phi, s}) \varphi(g) \overline{\varphi(h)}$$

for a suitable test function $f_{\Phi, s}$ where π ranges over the irreducible cuspidal representations of $\text{GL}_n(\mathbb{A})$ and φ over an orthonormal basis in the space of π .

In other words,

$$\langle \mathcal{E}(\Phi, \iota(g, \cdot), s), \varphi \rangle_{[Z \backslash \text{GL}_n]} = L(ns + \frac{1}{2}, \pi) \delta(f_{\Phi, s}) \varphi(g)$$

Relation between Godement–Jacquet and Hecke zeta integrals for GL_2 (Ginzburg–Soudry 2020)

Identify the space Mat_2 of 2×2 -matrices over \mathbb{Q} with \mathbb{Q}^4 via

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} \leftrightarrow (a, b, c, d).$$

This gives an action of $GL_4(\mathbb{A})$ on $\mathcal{S}(\text{Mat}_2(\mathbb{A}))$ by right translation. In these coordinates, the tensor representation

$$\iota = \iota_L \times \iota_R : GL_2 \times GL_2 \rightarrow GL_4$$

$$g_1^{-1} \begin{pmatrix} a & b \\ c & d \end{pmatrix} g_2 \leftrightarrow (a, b, c, d) \iota_L(g_1) \iota_R(g_2)$$

is given explicitly as

$$\iota_R(g) = \text{diag}(g, g), \quad \iota_L\left(\begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix}\right)^{-1} = \begin{pmatrix} \alpha & \gamma & & \\ \beta & \delta & \gamma & \\ & & \beta & \delta \\ & & & \delta \end{pmatrix}.$$

For $\Phi \in \mathcal{S}(\mathbb{A}^4)$ let $\tilde{\Phi} \in \mathcal{S}(\text{Mat}_2(\mathbb{A}))$ be the partial Fourier transform

$$\tilde{\Phi}\left(\begin{pmatrix} u \\ v \end{pmatrix}\right) = \int_{\mathbb{A}^2} \Phi(w, v) \psi([w, u]) dw, \quad u, v \in \mathbb{A}^2$$

(in the top row) where $[(a, b), (c, d)] = ad - bc$.

By the Poisson summation formula, for any $\Phi \in \mathcal{S}(\mathbb{A}^4)$

$$\sum_{\xi \in \text{Mat}_2(\mathbb{Q})} \omega_\psi(g) \tilde{\Phi}(\xi) = |\det g|^{\frac{1}{2}} \sum_{\xi \in \mathbb{Q}^4} \Phi(\xi g), \quad g \in \text{GL}_4(\mathbb{A})$$

where the representation ω_ψ of $\text{GL}_4(\mathbb{A})$ on $\mathcal{S}(\text{Mat}_2(\mathbb{A}))$ is

$$\omega_\psi(g) \tilde{\Phi}(v) = |\det g|^{\frac{1}{2}} \widetilde{\Phi(\cdot g)}(v), \quad g \in \text{GL}_4(\mathbb{A}), \quad v \in \text{Mat}_2(\mathbb{A}).$$

For any $\Phi \in \mathcal{S}(\text{Mat}_2(\mathbb{A}))$ we have

- ① $\omega_\psi(\iota_R(g)) \Phi\left(\begin{pmatrix} u \\ v \end{pmatrix}\right) = \Phi\left(\begin{pmatrix} ug^* \\ vg^* \end{pmatrix}\right), \quad u, v \in \mathbb{A}^2, \quad g \in \text{GL}_2(\mathbb{A})$
- ② $\omega_\psi(\iota_L\left(\begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix}\right)) \Phi(Y) = \psi(x \det Y) \Phi(Y), \quad Y \in \text{Mat}_2(\mathbb{A}), \quad x \in \mathbb{A}$

Consider

$$\theta_{\Phi}^*(g) = \sum_{0 \neq \xi \in \text{Mat}_2(\mathbb{Q})} \omega_{\psi}(g)\Phi(\xi), \quad g \in \text{GL}_4(\mathbb{A}).$$

For any $\xi \in \text{Mat}_2(\mathbb{Q})$ we have

$$\int_{\mathbb{Q} \backslash \mathbb{A}} \omega_{\psi}(\iota_L\left(\begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix}\right))\Phi(\xi)\psi(x) dx = \begin{cases} \Phi(\xi) & \text{if } \det \xi = 1, \\ 0 & \text{otherwise.} \end{cases}$$

On the other hand, $\text{GL}_2(\mathbb{Q})$ acts transitively on

$$\left\{ \begin{pmatrix} u \\ v \end{pmatrix} \mid [u, v] = 1 \right\} = \{ \xi \in \text{Mat}_2(\mathbb{Q}) \mid \det \xi = 1 \}$$

by $\begin{pmatrix} u \\ v \end{pmatrix} \mapsto \begin{pmatrix} ug^* \\ vg^* \end{pmatrix}$. It follows that

$$\int_{\mathbb{Q} \backslash \mathbb{A}} \theta_{\Phi}^*(\iota_L\left(\begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix}\right)g)\psi(x) dx = \sum_{\gamma \in T(\mathbb{Q}) \backslash \text{GL}_2(\mathbb{Q})} \omega_{\psi}(\iota_R(\gamma)g)\Phi(I)$$

where $T = \left\{ \begin{pmatrix} * & \\ & 1 \end{pmatrix} \right\}$ is the stabilizer of $\begin{pmatrix} (1,0) \\ (0,1) \end{pmatrix}$.

Taking Mellin transform,

$$\begin{aligned}
 & \int_{\mathbb{Q} \setminus \mathbb{A}} \mathcal{E}(\Phi, \iota_L\left(\begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix}\right)g, s) \psi(x) dx \\
 &= \int_{\mathbb{Q} \setminus \mathbb{A}} \int_{\mathbb{Q}^* \setminus \mathbb{I}} \theta_{\tilde{\Phi}}^*(\iota_R(tI) \iota_L\left(\begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix}\right)g) |\det tg|^s dt \psi(x) dx \\
 &= \sum_{\gamma \in T(\mathbb{Q}) \setminus \mathrm{GL}_2(\mathbb{Q})} \int_{\mathbb{Q}^* \setminus \mathbb{I}} \omega_\psi(\iota_R(t\gamma)g) \Phi(I) |\det tg|^s dt.
 \end{aligned}$$

Thus, for any cusp form ϕ on $[Z \setminus \mathrm{GL}_2]$

$$\begin{aligned}
 & \int_{\mathbb{Q} \setminus \mathbb{A}} \langle \mathcal{E}(\Phi, \iota\left(\begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix}\right), \cdot, s), \phi \rangle_{[Z \setminus \mathrm{GL}_2]} \psi(x) dx \\
 &= \int_{T(\mathbb{Q}) \setminus \mathrm{GL}_2(\mathbb{A})} \omega_\psi(\iota_R(g)) \Phi(I) |\det g|^{2s} \phi(g) dg \\
 &= \int_{T(\mathbb{A}) \setminus \mathrm{GL}_2(\mathbb{A})} \omega_\psi(\iota_R(g)) \Phi(I) \int_{\mathbb{Q}^* \setminus \mathbb{I}} \phi\left(\begin{pmatrix} t & \\ & 1 \end{pmatrix}g\right) |t|^{2s} dt |\det g|^{2s} dg.
 \end{aligned}$$

The function $\omega_\psi(\iota_R(g)) \Phi(I)$ is a test function on $T(\mathbb{A}) \setminus \mathrm{GL}_2(\mathbb{A})$.

Recall

$$\langle \mathcal{E}(\Phi, \iota(g, \cdot), s), \phi \rangle_{[Z \backslash \mathrm{GL}_2]} \text{ “=” } L(2s + \frac{1}{2}, \pi) \phi(g).$$

Taking Whittaker coefficients of both sides we get

$$\begin{aligned} & \int_{\mathbb{Q} \backslash \mathbb{A}} \langle \mathcal{E}(\Phi, \iota(\begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix}, \cdot), s), \phi \rangle_{[Z \backslash \mathrm{GL}_2]} \psi(x) dx \\ & \text{ “=” } L(2s + \frac{1}{2}, \pi) \int_{\mathbb{Q} \backslash \mathbb{A}} \phi(\begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix}) \psi(x) dx. \end{aligned}$$

On the other hand,

$$\begin{aligned} & \int_{\mathbb{Q} \backslash \mathbb{A}} \langle \mathcal{E}(\Phi, \iota(\begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix}, \cdot), s), \phi \rangle_{[Z \backslash \mathrm{GL}_2]} \psi(x) dx \\ & = \int_{T(\mathbb{A}) \backslash \mathrm{GL}_2(\mathbb{A})} \omega_\psi(\iota_R(g)) \Phi(I) \int_{\mathbb{Q}^* \backslash \mathbb{I}} \phi(\begin{pmatrix} t & \\ & 1 \end{pmatrix} g) |t|^{2s} dt |\det g|^{2s} dg. \end{aligned}$$

Thus, the Hecke integral is obtained from the Godement–Jacquet integral by taking a Whittaker coefficient.

Similar analysis shows that for any cuspidal representation π of $G(\mathbb{A})$, $G = \mathrm{GL}_{2n}$, $n \geq 1$ and $\phi \in V_\pi$ we have

$$\int_{[\mathrm{GL}_n]} \phi\left(\begin{pmatrix} g & \\ & h \end{pmatrix}\right) |\det g|^s dg$$

$$= L\left(s + \frac{1}{2}, \pi\right) \int_{[\mathrm{Mat}_n]} \phi\left(\begin{pmatrix} I & X \\ 0 & I \end{pmatrix} \begin{pmatrix} h & \\ & h \end{pmatrix}\right) \psi(\mathrm{tr} X) dX, \quad h \in [\mathrm{GL}_n].$$

This can also be proved using unfolding (Bump–Furusawa–Ginzburg). Note that the Shalika functional on the right-hand side is not unique. Integrating against an Eisenstein series (in the h variable) we get a relation

$$\int_{[\mathrm{GL}_n \times \mathrm{GL}_n]} \phi\left(\begin{pmatrix} g & \\ & h \end{pmatrix}\right) \mathcal{E}(h, s') |\det g|^{s - \frac{1}{2}} dg dh$$

$$= L\left(s + \frac{1}{2}, \pi\right) \int_{[\mathrm{GL}_n]} \int_{[\mathrm{Mat}_n]} \phi\left(\begin{pmatrix} I & X \\ 0 & I \end{pmatrix} \begin{pmatrix} h & \\ & h \end{pmatrix}\right) \psi(\mathrm{tr} X) dX \mathcal{E}(h, s') dh$$

$$= L\left(s + \frac{1}{2}, \pi\right) L\left(s' + \frac{1}{2}, \pi, \wedge^2\right) \int_{[M]} \phi(u) \psi_N(u) du$$

(Jacquet–Shalika, Bump–Friedberg, Friedberg–Jacquet)

More reproducing kernels for L -functions – the doubling method of Piatetski-Shapiro–Rallis

Let V be a finite-dimensional vector space over \mathbb{Q} and h a non-degenerate symmetric or alternating bilinear form on V .

Let $G = \text{Isom}(V, h)$, an orthogonal or symplectic group.

Consider the “doubled” group $H = \text{Isom}(V \oplus V, h \oplus (-h))$.

The diagonal

$$V^\Delta = \{(v, v) \mid v \in V\}$$

is a maximal isotropic subspace of $V \oplus V$ defined over \mathbb{Q} . Its stabilizer P is a maximal parabolic subgroup of H whose Levi part is isomorphic to $\text{GL}(V)$.

In particular, H is split (even if G isn't).

We can form an Eisenstein series $\mathcal{E}(f, s)$ induced from the character $|\det_V|^s$ of $P(\mathbb{A})$.

We have an embedding

$$\iota : G \times G \hookrightarrow H$$

Doubling method (Piatetski-Shapiro–Rallis, 1987)

The restriction of $\mathcal{E}(f, s)$ to $[G \times G]$ is a reproducing kernel for the L -function pertaining to the canonical representation

$$\text{Can} : {}^L G \rightarrow \text{GL}(N),$$

where

$$N = \dim V + \begin{cases} 1 & \text{if } h \text{ is alternating} \\ 0 & \text{if } h \text{ is symmetric and } \dim V \text{ is even} \\ -1 & \text{if } h \text{ is symmetric and } \dim V \text{ is odd} \end{cases}$$

In other words, for any cuspidal $\phi \in V_\pi$

$$\langle \mathcal{E}^*(f, \iota(\cdot, g), s), \phi \rangle_{[G]} = L(s + \frac{1}{2}, \pi, \text{Can})\phi(g).$$

- Unlike the case of the general linear group, the meromorphic continuation of $\mathcal{E}(f, s)$ is not merely a consequence of Poisson summation formula, but rather of the general theory of Eisenstein series.
- A related issue is the normalization $\mathcal{E}^*(f, s)$ of $\mathcal{E}(f, s)$. Currently, there is no satisfactory geometric way to do it.
- The doubling method gives meromorphic continuation and functional equation of $L(s, \pi, \text{Can})$. Moreover, the poles are controlled by those of $\mathcal{E}^*(f, s)$ (which have to be studied separately but they do not depend on π – work by Kudla and Rallis in the 1990s).
- Just as with the Godement–Jacquet integral, many other integral representations of $L(s, \pi, \text{Can})$ can be derived by taking appropriate models (Ginzburg, Soudry, Adrianov, Piatetski-Shapiro–Rallis). This was only realized in retrospect (rather recently) by Ginzburg–Soudry.

Generalized doubling (Cai–Friedberg–Ginzburg–Kaplan, 2017)

Let W be another finite-dimensional vector space, $n = \dim W$. Consider the symmetric bilinear form β on $W \oplus W^\vee$ such that

$$\beta(w, w^\vee) = \langle w^\vee, w \rangle, \quad \beta|_{W \times W} \equiv 0, \quad \beta|_{W^\vee \times W^\vee} \equiv 0.$$

Let

$$H = \text{Isom}(V \otimes (W \oplus W^\vee), h \otimes \beta).$$

For $n = 1$, this is isomorphic to $(V \oplus V, h \oplus (-h))$ considered before. Note that $V \otimes W$ is a maximal isotropic subspace. Let P be the maximal parabolic subgroup stabilizing $V \otimes W$. Its Levi subgroup M is isomorphic to $\text{GL}(V \otimes W) \simeq \text{GL}(n \cdot \dim V)$.

Fix a cuspidal representation τ of $\text{GL}(W)(\mathbb{A}) \simeq \text{GL}_n(\mathbb{A})$.

Let σ be the corresponding Speh representation of $M(\mathbb{A})$.

Let $\mathcal{E}(\cdot, s)$ be the (suitably normalized) Eisenstein series induced from σ .

Fix a complete flag

$$0 = W_0 \subsetneq W_1 \subsetneq \cdots \subsetneq W_n = W$$

Let $Q = LU$ be the parabolic subgroup of H stabilizing the flag

$$0 \subsetneq V \otimes W_1 \subsetneq \cdots \subsetneq V \otimes W_{n-1}.$$

Thus, $L \simeq \mathrm{GL}(V) \times \cdots \times \mathrm{GL}(V) \times \mathrm{Isom}(V \oplus V, h \oplus (-h))$.

The stabilizer of a suitable generic character ψ_U of U is the image of

$$\iota : G \times G \hookrightarrow L, \quad (g, h) \mapsto \{(g, \dots, g, (g, h))\}.$$

Generalized doubling (Cai–Friedberg–Ginzburg–Kaplan, 2017)

As a function on $[G \times G]$, the Fourier coefficient

$$\mathcal{E}^{\psi_U}(f, \cdot, s) = \int_{[U]} \mathcal{E}(f, u \cdot, s) \psi_U(u) du$$

is a reproducing kernel for the L -function $L(s, \cdot \otimes \tau, \mathrm{Can} \otimes \mathrm{St})$ i.e. $\forall \phi \in V_\pi$

$$\left\langle \mathcal{E}^{\psi_U}(\iota(\cdot, g), s), \phi \right\rangle_{[G]} = L\left(s + \frac{1}{2}, \pi \otimes \tau, \mathrm{Can} \otimes \mathrm{St}\right) \phi(g).$$

- A similar construction gives a reproducing kernel for the L -function

$$L(s, \pi \otimes \tau)L(s, \pi^\vee \otimes \tau)$$

where π ranges over the cuspidal representations of $\mathrm{GL}_m(\mathbb{A})$.
However, I'm not aware of an explicit reproducing kernel for

$$L(s, \pi \otimes \tau)$$

itself.

- (Double descent, Ginzburg–Soudry, 2022) Taking $n = N$, the residue

$$\mathrm{Res}_{s=\frac{1}{2}} \mathcal{E}^{\psi_U}(f, \cdot, s)$$

is a reproducing kernel for the space of cusps forms that functorially lift to τ , i.e.

$$\mathrm{Res}_{s=\frac{1}{2}} \mathcal{E}^{\psi_U}(f, \iota(g, h), s) \text{ "="} \sum_{\phi} \phi(g) \overline{\phi(h)}.$$