



Protecting Quantum Information with Optimal Control

Matthew Grace, Jason Dominy,
Wayne Witzel, and Malcolm Carroll

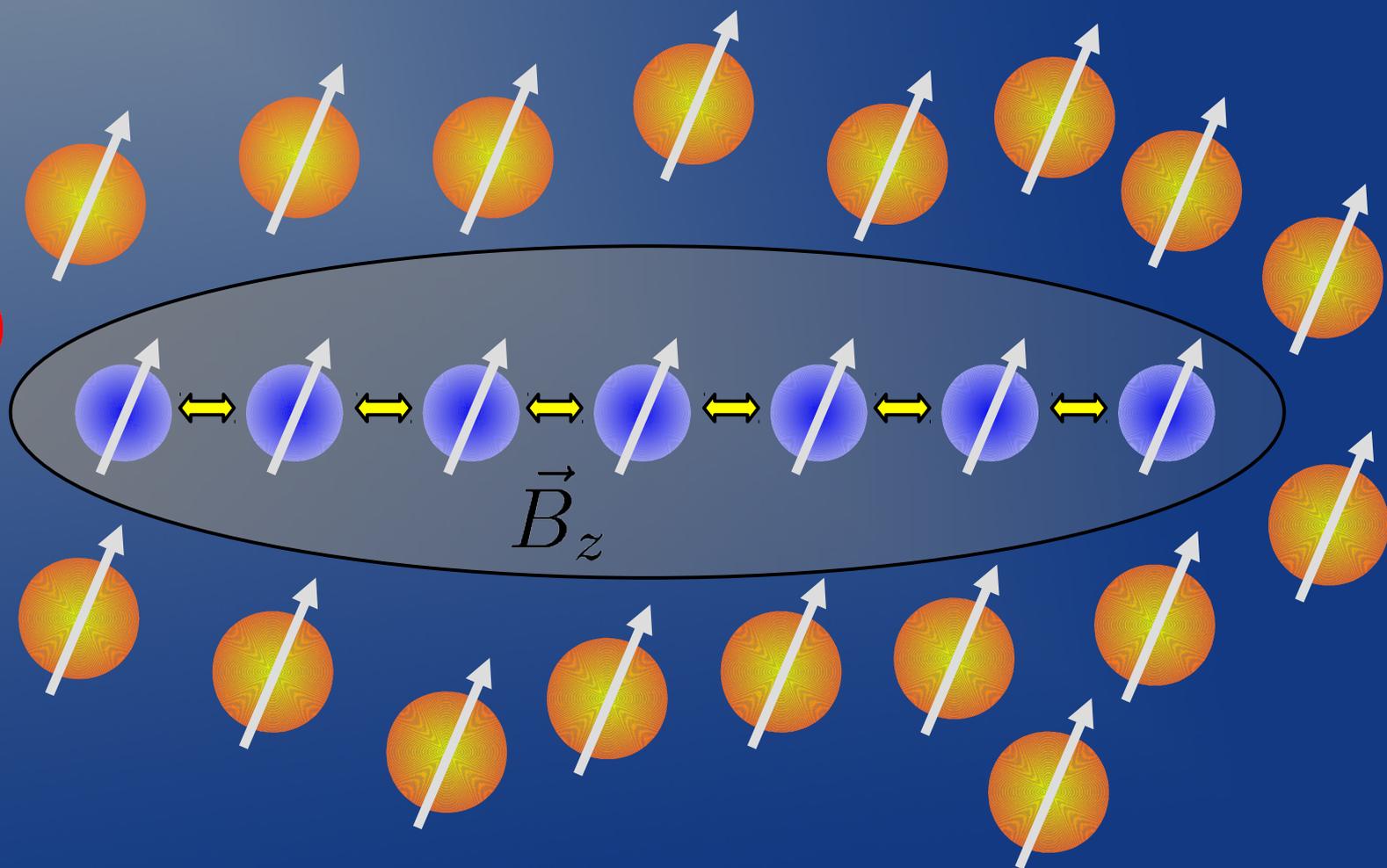
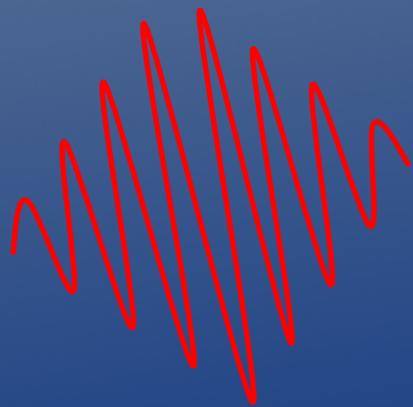


Quantum Information & Decoherence

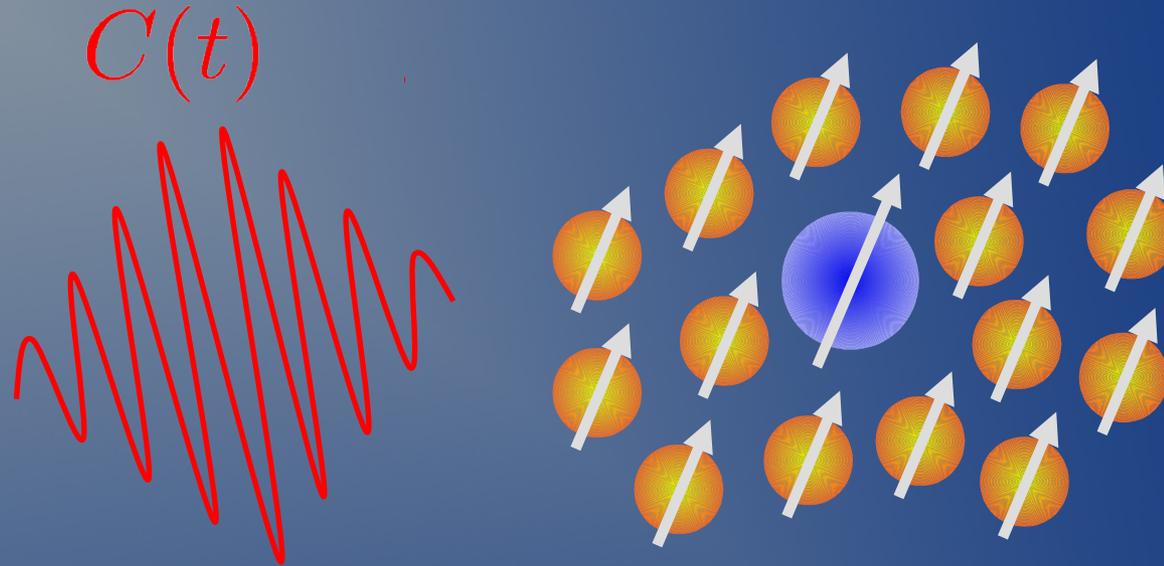
Realistic physical systems are (noisy) open systems
→ they interact with the surrounding environment.

Reality:

$$C(t) + \delta C(t)$$



Model Open System: Interacting Quantum Spins



$$H(t) = H_s(t) + H_{\text{int}} + H_e$$

Schrödinger equation: $\dot{U}(t) = -iH(t)U(t)$

Objective: Generate target system time evolutions

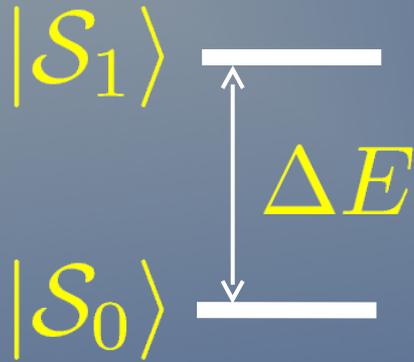
Model Open System: Interacting Quantum Spins

$$H = H_s + H_{\text{int}} + H_e$$

$$H = H_0 - \vec{\sigma} \cdot \vec{C}(t) - \vec{\sigma} \cdot \vec{\Gamma} + H_e$$

Control pulse area: $\theta(t) = \int_0^t \vec{C}(\tau) d\tau$

Qubit Dynamics & the Bloch Sphere

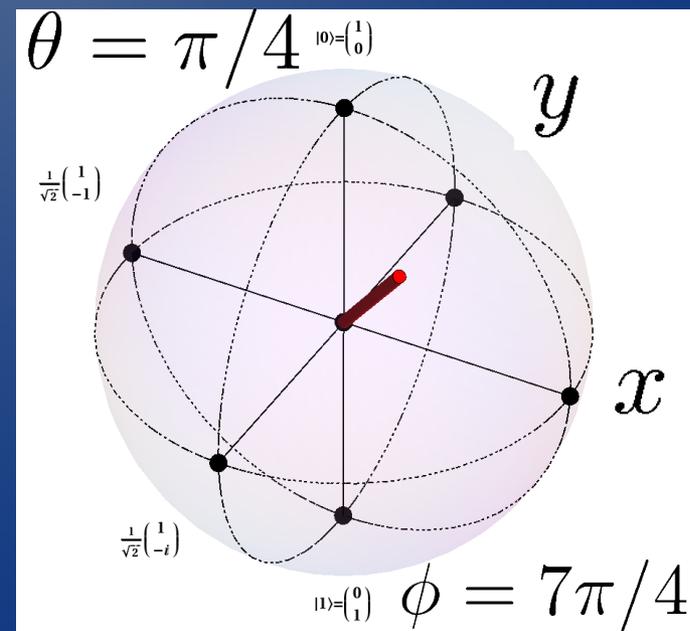
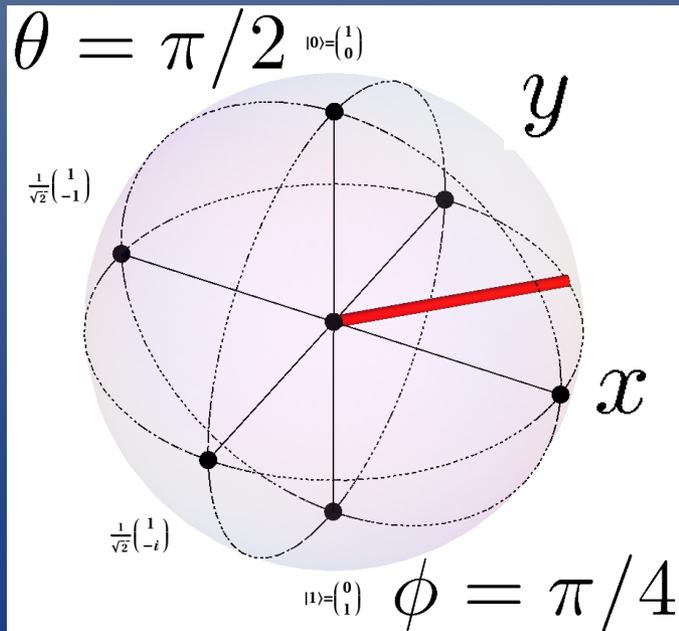


“Pure” spin states have the form:

$$|\psi\rangle = c_0|S_0\rangle + c_1|S_1\rangle,$$

where $|c_0|^2 + |c_1|^2 = 1$.

$$|\psi\rangle = \cos(\theta/2)|S_0\rangle + \exp(i\phi)\sin(\theta/2)|S_1\rangle.$$

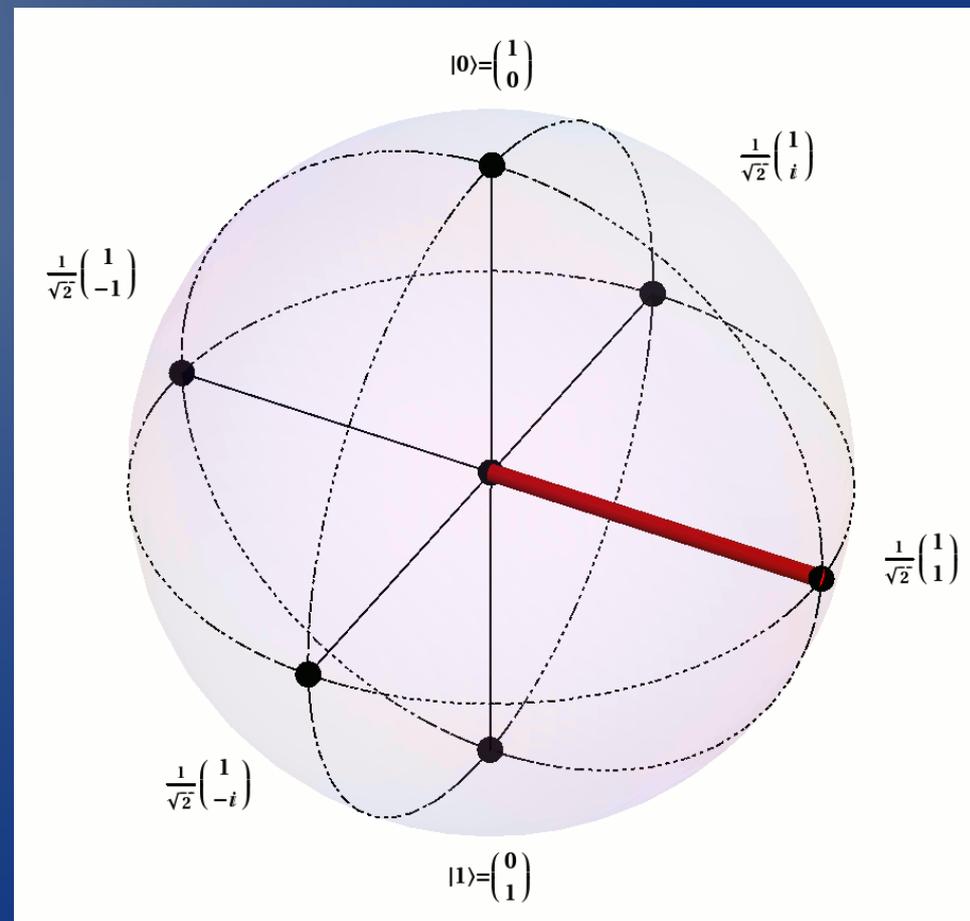
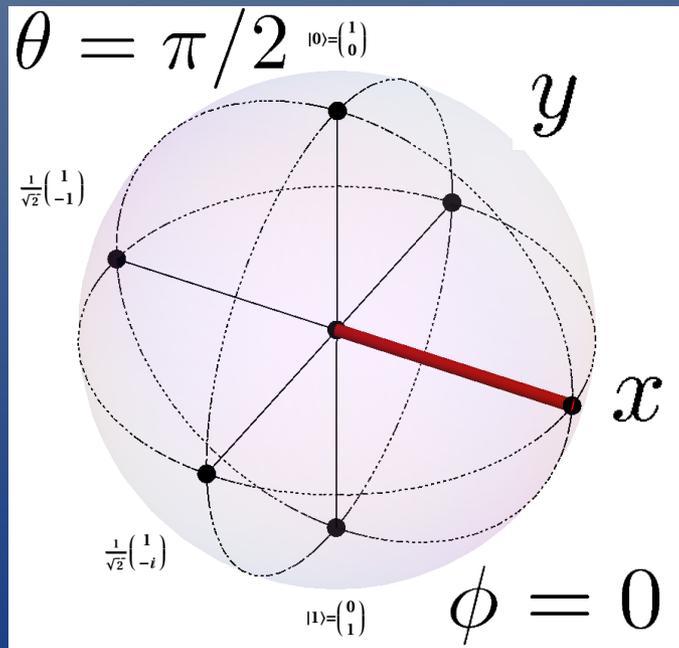


Qubit Dynamics & the Bloch Sphere: Memory Channels

Consider a decoherence process
for states in the xy -plane, i.e.,

$$|\psi\rangle = \frac{|\mathcal{S}_0\rangle + \exp(i\phi)|\mathcal{S}_1\rangle}{\sqrt{2}}.$$

Bloch vector dephasing
→ uncertainty in ϕ .



Qubit Dynamics & the Bloch Sphere: Memory Channels

With the “right” set of physical rotations,
this error can be corrected → “Hahn-echo”

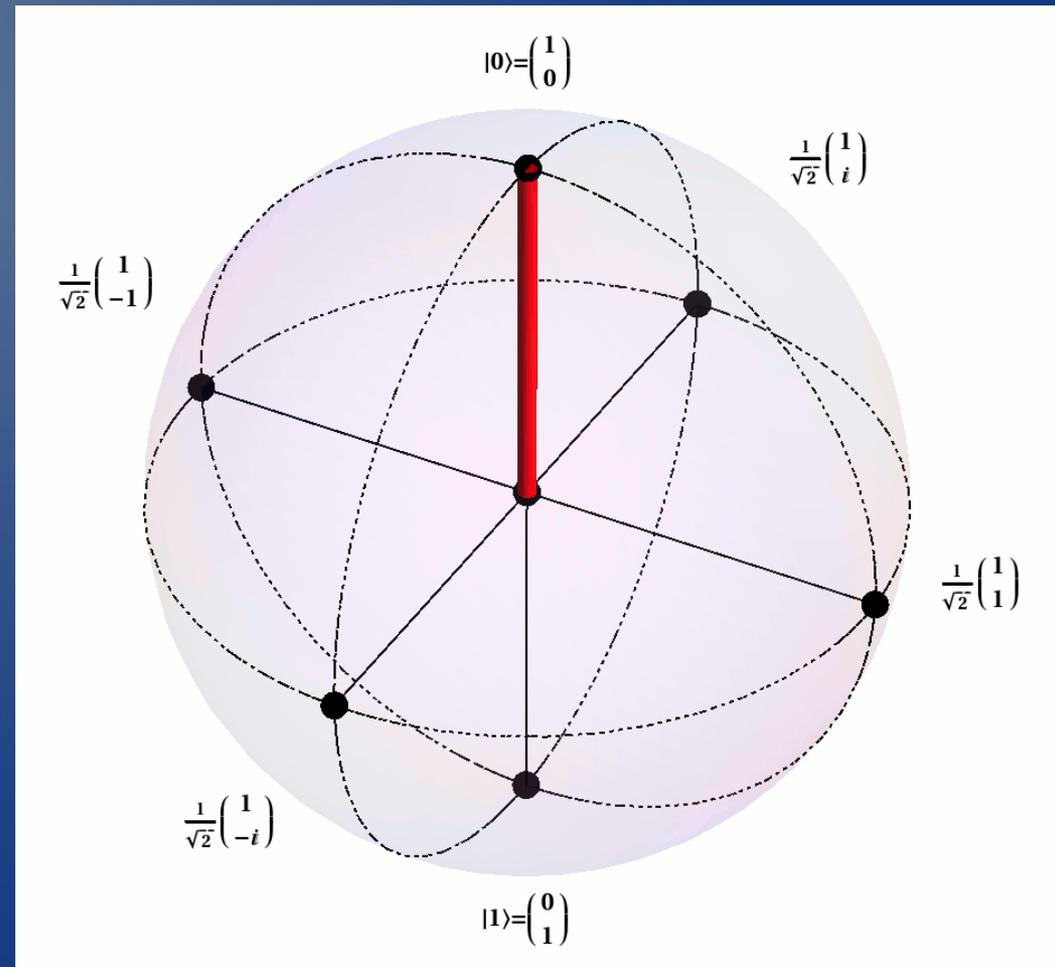
Hahn-echo pulse sequence:

$$\tau \rightarrow Y_{\pi} \rightarrow \tau \rightarrow Y_{\pi}$$

↑
Free evolution

↑
 π -rotation about
the y -axis

$$\pi\text{-pulses} \Rightarrow \theta(t_f) = \pi$$



Quantum Memory Channels: Dynamical-Decoupling Pulse Sequences

$$\prod_i^N U_i \approx \mathcal{I}, \quad \text{where } U_i \text{ represent } \pi \text{ and } \pi/2$$

rotations and free evolutions.

This is an approximation to \mathcal{I} because

- $\{U_i\}$ and N are finite
- Non-unitary evolution is corrected with unitary “time-reversal” operations

Dynamical-Decoupling Pulses

Aside from satisfying geometric and pulse area constraints, e.g.,

$$\theta(t_f) = \int_0^{t_f} \vec{C}(\tau) d\tau,$$

What other control-field conditions can improve gate fidelity?

Dynamical-Decoupling Pulses

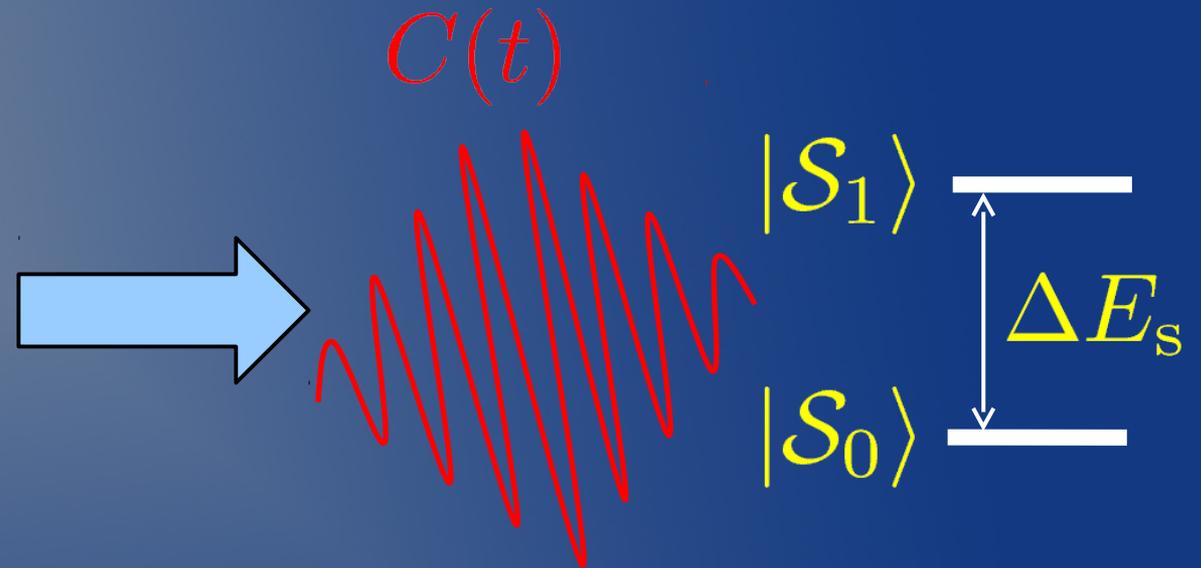
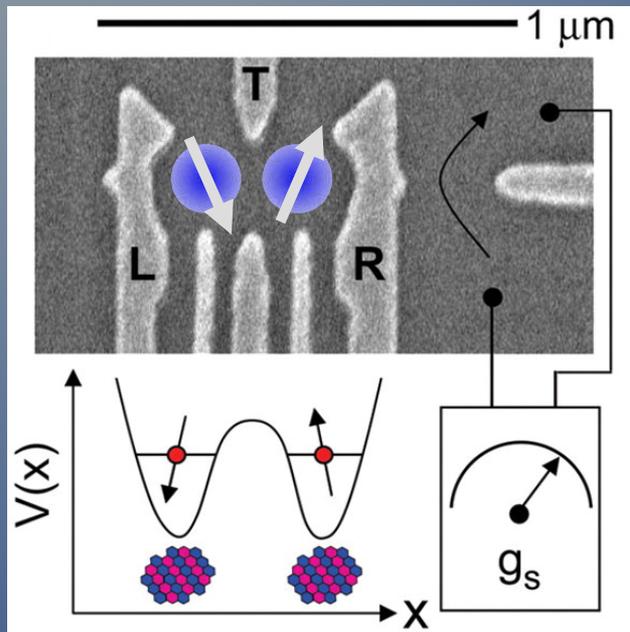
To remove 1st- and 2nd-order errors in π and $\pi/2$ z -axis rotations when $\vec{\Gamma} = \Gamma_x \Rightarrow \vec{\eta} = 0$:

$$\eta_1 = \int_0^{t_f} \sin[\theta(t)] dt, \quad \eta_2 = \int_0^{t_f} \cos[\theta(t)] dt,$$

$$\eta_3 = \int_0^{t_f} t \sin[\theta(t)] dt, \quad \eta_4 = \int_0^{t_f} t \cos[\theta(t)] dt,$$

$$\eta_5 = \int_0^{t_f} \int_0^{t_f} \sin[\theta(t_1) - \theta(t_2)] \text{sign}(t_1 - t_2) dt_1 dt_2,$$

Double Quantum Dot: Effective One-Qubit Model



J. Petta, *et al.*, Science, **309** (2005)

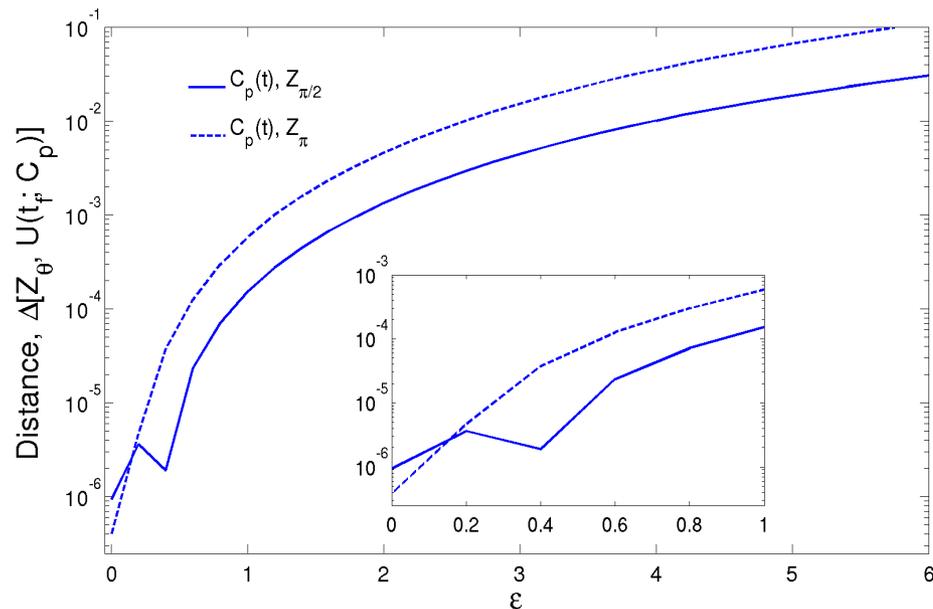
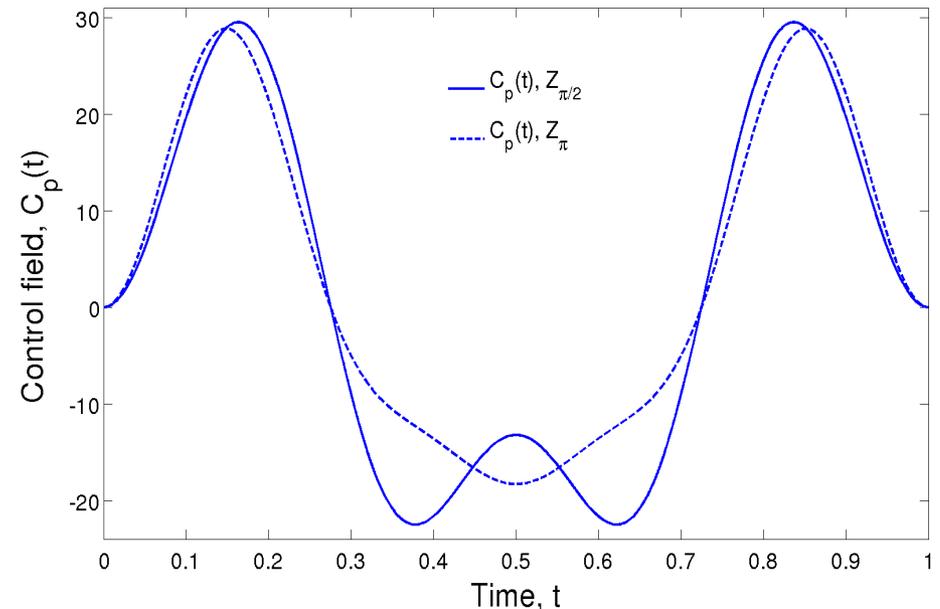
$$H = H_0 - \vec{\sigma} \cdot \vec{C}(t) - \vec{\sigma} \cdot \vec{\Gamma} + H_e$$

$$\rightarrow H = C_z(t) \sigma_z + \epsilon \sigma_x$$

Dynamical-Decoupling Pulses

S. Pasini, *et al.*, Phys. Rev. A, **80** (2009)

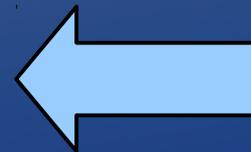
Feasible controls
satisfying $\vec{\eta} = 0$



Gate distance:

$$\Delta(Z_{\theta}, U_{t_f}) =$$

$$\sqrt{1 - \frac{1}{2} \left| \text{Tr} \left(Z_{\theta}^{\dagger} U_{t_f} \right) \right|}$$



Dynamical Decoupling + Optimal Control Pulses

By searching the space of controls satisfying

1. $\vec{\eta} \approx 0$

2. $\Delta(Z_\pi, U_{t_f}) = \sqrt{1 - \frac{1}{2} \left| \text{Tr} \left(Z_\pi^\dagger U_{t_f} \right) \right|} \approx 0,$

and incorporating parameter estimates for ϵ ,
we improve control fidelities for Z_π .

Systematic searching \rightarrow Optimal control theory

Quantum Optimal Control Theory

- Define an objective:

$$\Delta(Z_\theta, U_{t_f}) = \sqrt{1 - \frac{1}{2} \left| \text{Tr} \left(Z_\theta^\dagger U_{t_f} \right) \right|}$$

- Incorporate constraints:

- Schrödinger's equation
- Experimental limitations of the control field

$$\mathcal{J} = \Delta \circ U_{t_f} + \frac{\alpha}{2} \int_0^{t_f} \left\| \vec{C}(t) \right\|^2 dt$$

- Optimize iteratively

- Evolutionary algorithms
- Gradient-based methods

$$\nabla_c \mathcal{J} = 0$$

DD+OC Optimization Procedure

After calculating $\nabla_c \mathcal{J}$, all gradient directions $\nabla_c \eta_i$ are removed:

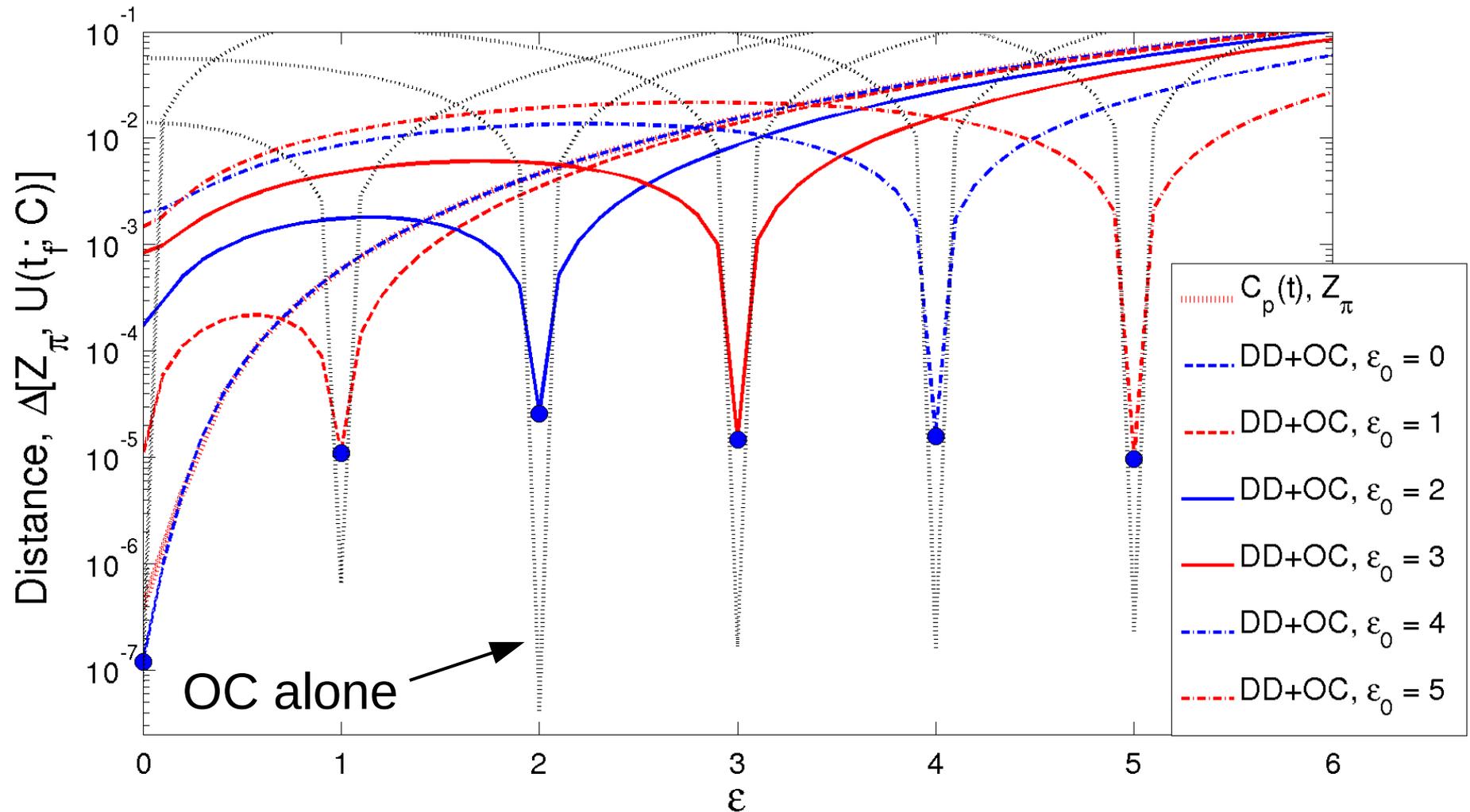
$$\nabla_c \mathcal{J} \longrightarrow \nabla_c \mathcal{J} - \sum_i p_i \nabla_c \eta_i,$$

$$\text{so } \langle \nabla_c \mathcal{J}, \nabla_c \eta_i \rangle = 0,$$

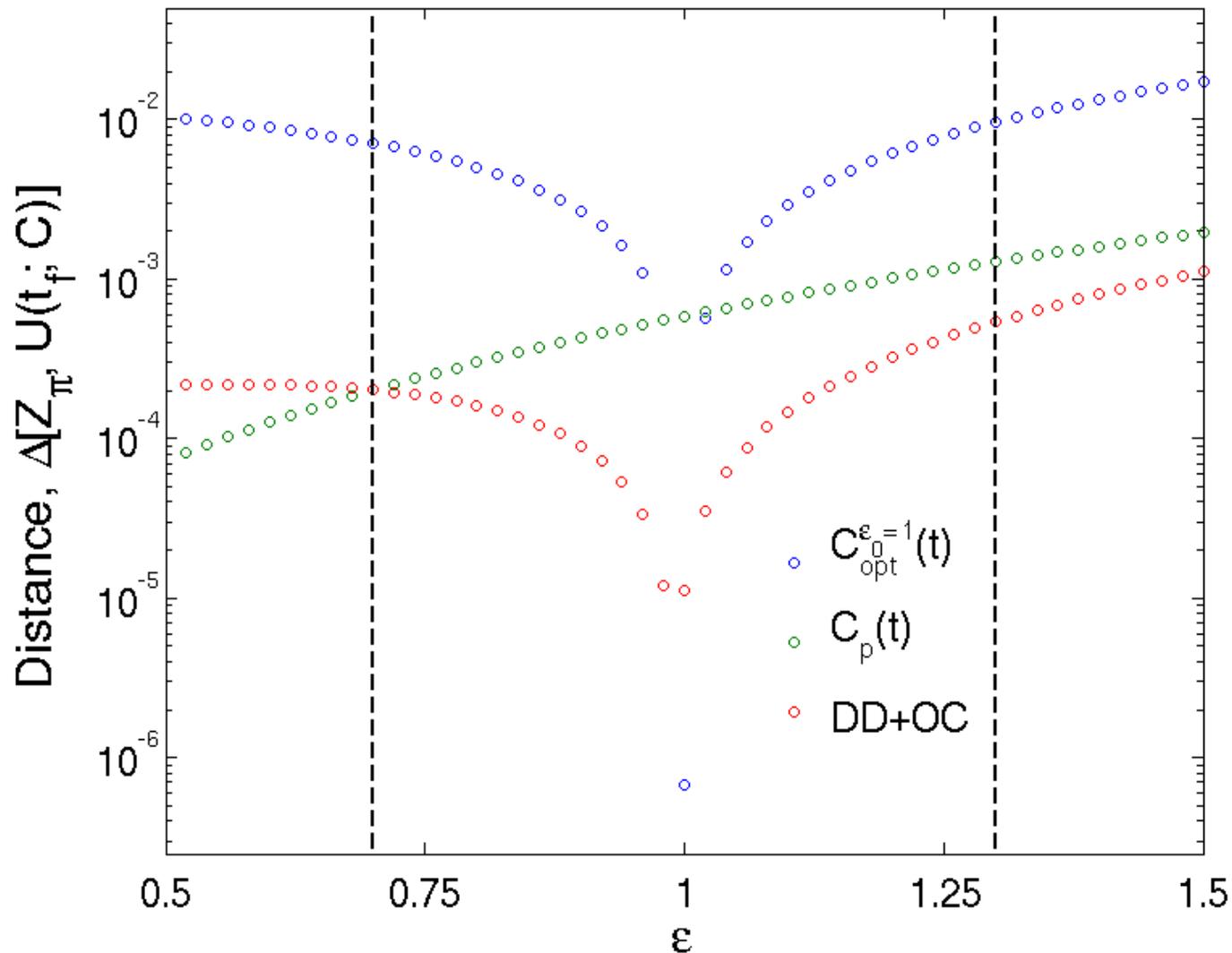
where \vec{p} is constructed from

$$G_{ij} = \langle \nabla_c \eta_i, \nabla_c \eta_j \rangle \text{ and } \langle \nabla_c \mathcal{J}, \nabla_c \eta_i \rangle.$$

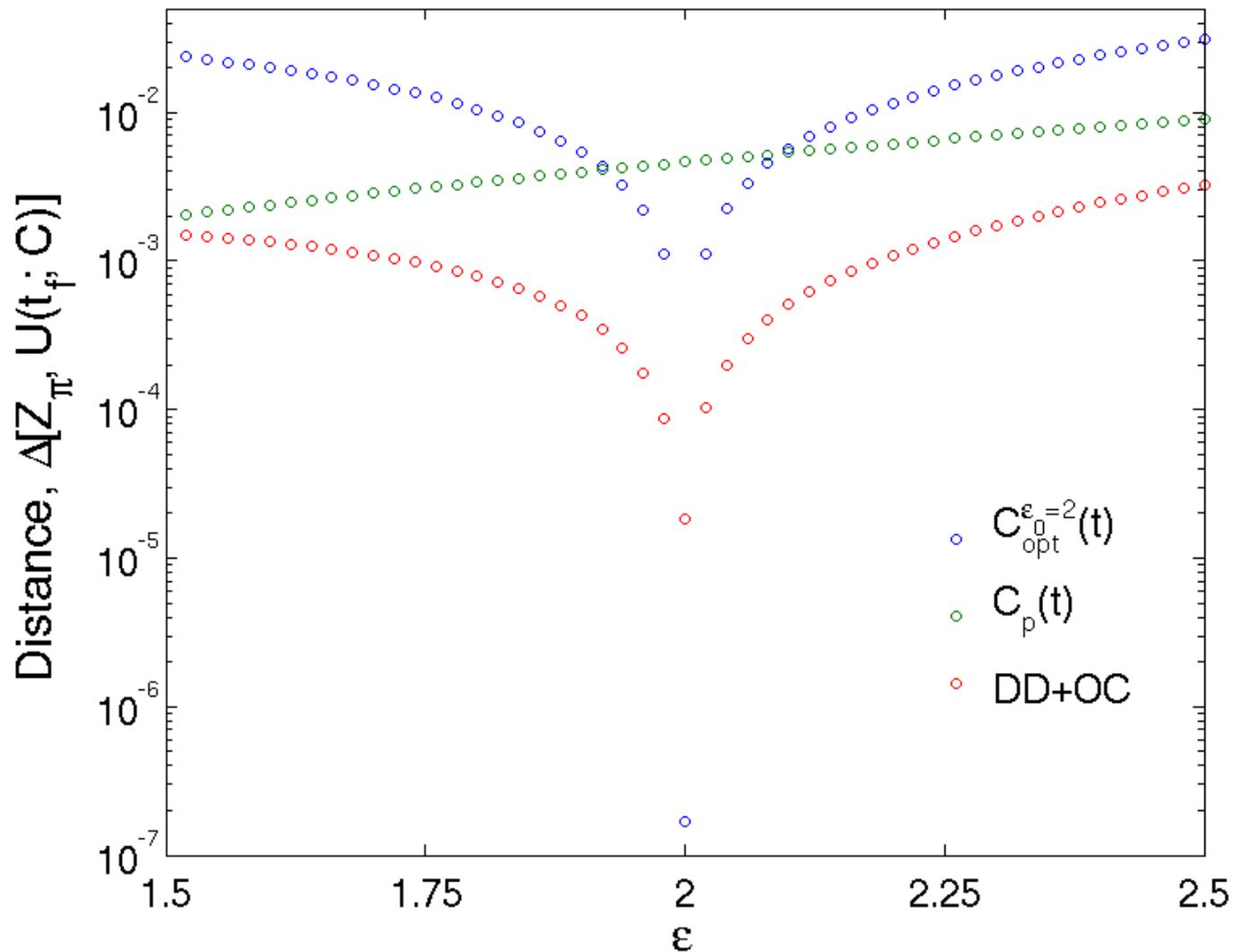
Gate Distances from OC and DD+OC



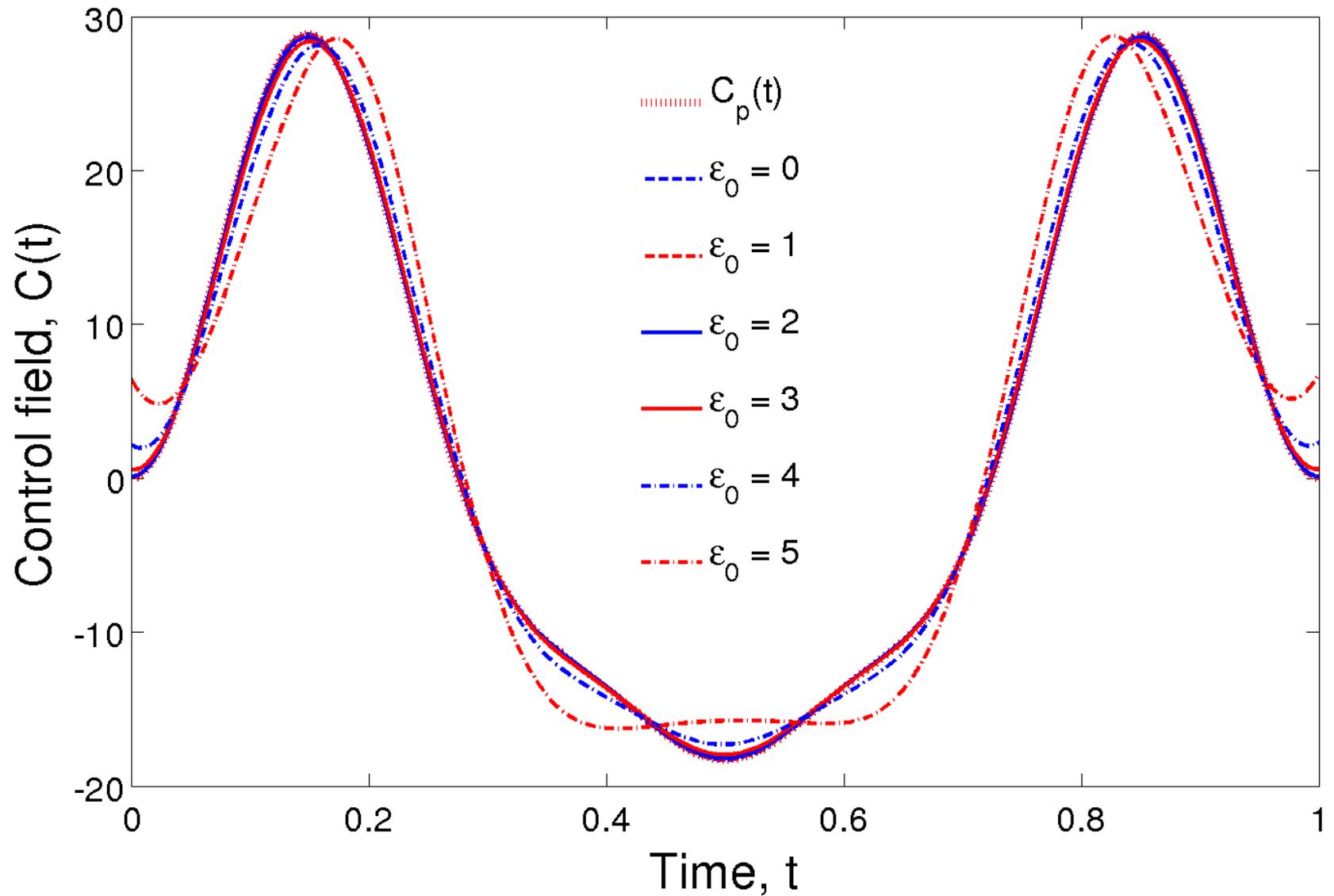
DD+OC Pulses: Improved Robustness 1



DD+OC Pulses: Improved Robustness 2



DD+OC Pulses



Conclusions and Current Work

- Demonstrated dynamical decoupling + optimal control for improved gate fidelity and robustness
- Extend formalism to arbitrary rotation axes and perturbative expansions about arbitrary ϵ .
- Explore robustness to control field variations

Calculating the Distance Measure with the Hilbert-Schmidt Norm

$$\Delta(U, V) = \lambda_{\min}_{\Phi_e} \left\{ \|U - (\mathbb{1}_s \otimes \Phi_e)V\|_{\text{HS}} \right\}$$

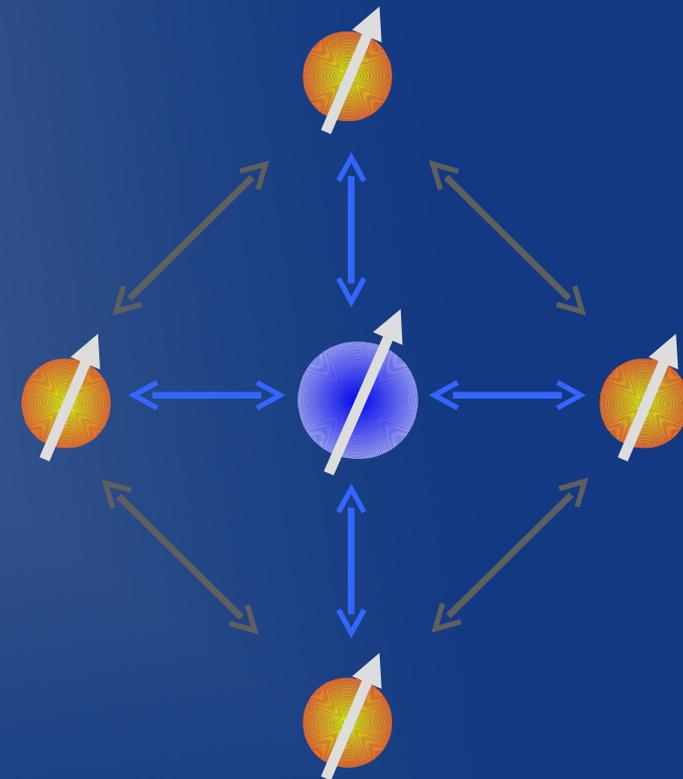
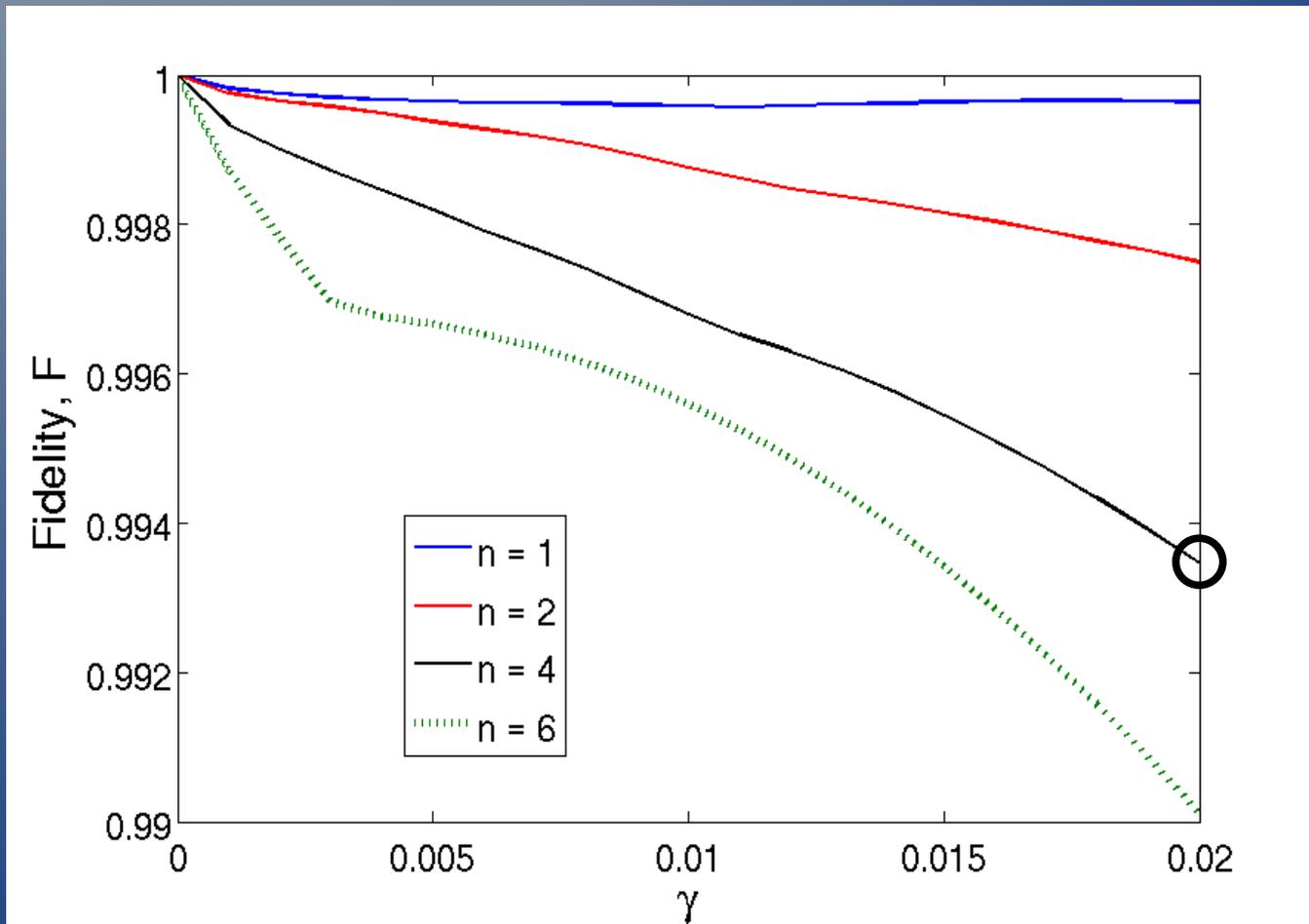
After some relatively straightforward linear algebra . . .

$$= \left[1 - \frac{1}{n} \|\text{Tr}_s(UV^\dagger)\|_{\text{Tr}} \right]^{1/2}$$

Nice analytical result; solved numerically in practice.

Quantum Control Results: One-Qubit Operations

Multiparticle environment: Hadamard gate



Gate Robustness to System Variations

Optimal Hadamard gate with a four-particle environment:

$$\gamma = 0.02, \quad \gamma' = 0.0175, \quad \text{and} \quad F \approx 0.9934.$$

This control is applied to an ensemble of systems with random variations in γ and γ' given by $\Delta\gamma/\gamma = 1/8$.

