

Reduced models for linearly elastic thin films allowing for fracture, debonding or delamination

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Presentation of the 3D model

The system

$$\Omega^\varepsilon = \Omega_f^\varepsilon \cup \Omega_b^\varepsilon \cup \Omega_s^\varepsilon$$

is made of

- a **film** $\Omega_f^\varepsilon := \omega \times (0, \varepsilon)$ ($\omega \subset \mathbb{R}^2$ open, bounded) deposited on
- an infinite **substrate** $\Omega_s^\varepsilon := \omega \times (-\infty, -\varepsilon)$ through
- a **bonding layer** $\Omega_b^\varepsilon := \omega \times [-\varepsilon, 0]$.

- ✓ **Kinematics** : isotropic linearized elasticity (displacement $v : \Omega^\varepsilon \rightarrow \mathbb{R}^3$);
- ✓ **Loadings** : inelastic strain e_0 , displacement in the substrate v_0 ;
- ✓ **Cracks** : sets $\Gamma \subset \overline{\Omega}^\varepsilon$ of finite area ;
- ✓ **Total energy** : for all $\Gamma \subset \overline{\Omega}^\varepsilon$ and $v : \Omega^\varepsilon \setminus \Gamma \rightarrow \mathbb{R}^3$ with $v = v_0$ in Ω_s^ε ,

$$J(\varepsilon)(v, \Gamma) = \int_{\Omega^\varepsilon \setminus \Gamma} \mathbb{C}^\varepsilon(e(v) - e_0) : (e(v) - e_0) \, dx + \int_\Gamma \kappa^\varepsilon \, d\mathcal{H}^2,$$

where $e(v) = (\nabla v + \nabla v^T)/2$ is the linearized strain.

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- ✓ **Total energy** : for all $v \in SBD(\Omega^\varepsilon)$ with $v = 0$ in Ω_s^ε ,

$$J(\varepsilon)(v) = \int_{\Omega^\varepsilon \setminus J_v} \mathbb{C}^\varepsilon e(v) : e(v) \, dx + \int_{J_v} \kappa^\varepsilon \, d\mathcal{H}^2,$$

where $E v = (Dv + Dv^T)/2 = e(v)dx + (v^+ - v^-) \odot \nu_v \delta_{J_v}$.

We assume that

$$\mathbb{C}^\varepsilon = \underbrace{\mathbb{C}_f}_{(\lambda_f, \mu_f)} \chi_{\Omega_f^\varepsilon} + \varepsilon^2 \underbrace{\mathbb{C}_b}_{(\lambda_b, \mu_b)} \chi_{\Omega_b^\varepsilon}, \quad \kappa^\varepsilon = \kappa_f \chi_{\Omega_f^\varepsilon} + \varepsilon \kappa_b \chi_{\Omega_b^\varepsilon}.$$

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$$\begin{aligned} J(\varepsilon)(v) &= \int_{\Omega_f^\varepsilon \setminus J_v} \mathbb{C}_f e(v) : e(v) \, dx + \kappa_f \mathcal{H}^2(J_v \cap \Omega_f^\varepsilon) \\ &\quad + \varepsilon^2 \int_{\Omega_b^\varepsilon \setminus J_v} \mathbb{C}_b e(v) : e(v) \, dx + \varepsilon \kappa_b \mathcal{H}^2(J_v \cap \Omega_b^\varepsilon), \end{aligned}$$

where $E v = (D v + D v^T)/2 = e(v) dx + (v^+ - v^-) \odot \nu_v \delta_{J_v}$.

Scaling

$$\Omega = \Omega^1, \Omega_f := \Omega_f^1, \Omega_b := \Omega_b^1, \Omega_s := \Omega_s^1$$

For all $x = (x', x_3) \in \Omega$,

$$u_\alpha(x', x_3) = v_\alpha(x', \varepsilon x_3), \quad u_3(x', x_3) = \varepsilon v_3(x', \varepsilon x_3).$$

Then for all $u \in SBD(\Omega)$ with $v = 0$ in Ω_s ,

$$J_\varepsilon(u) = \varepsilon^{-1} J(\varepsilon)(v) = J_\varepsilon^f(u) + J_\varepsilon^b(u),$$

where

$$J_\varepsilon^f(u) := \int_{\Omega_f \setminus J_u} \mathbb{C}_f e^\varepsilon(u) : e^\varepsilon(u) dx + \kappa_f \int_{J_u \cap \Omega_f} |((\nu_u)', \varepsilon^{-1}(\nu_u)_3)| d\mathcal{H}^2,$$

$$J_\varepsilon^b(u) := \varepsilon^2 \int_{\Omega_b \setminus J_u} \mathbb{C}_b e^\varepsilon(u) : e^\varepsilon(u) dx + \kappa_b \varepsilon \int_{J_u \cap \Omega_b} |((\nu_u)', \varepsilon^{-1}(\nu_u)_3)| d\mathcal{H}^2,$$

and

$$e^\varepsilon(u) := \begin{pmatrix} e_{11}(u) & e_{12}(u) & \varepsilon^{-1} e_{13}(u) \\ e_{12}(u) & e_{22}(u) & \varepsilon^{-1} e_{23}(u) \\ \varepsilon^{-1} e_{13}(u) & \varepsilon^{-1} e_{23}(u) & \varepsilon^{-2} e_{33}(u) \end{pmatrix}.$$

Derivation of elastic foundations

Derivation of transverse cracks

Fracture, debonding and delamination

For all $u \in H^1(\Omega; \mathbb{R}^3)$ with $u = 0$ in Ω_s ,

$$J_\varepsilon^f(u) := \int_{\Omega_f} \mathbb{C}_f e^\varepsilon(u) : e^\varepsilon(u) dx,$$

$$J_\varepsilon^b(u) := \varepsilon^2 \int_{\Omega_b} \mathbb{C}_b e^\varepsilon(u) : e^\varepsilon(u) dx.$$

For all $u \in H^1(\Omega; \mathbb{R}^3)$ with $u = 0$ in Ω_s ,

$$\begin{aligned}
 J_\varepsilon^f(u) &:= \int_{\Omega_f} \left[\frac{\lambda_f}{2} e_{\alpha\alpha}(u) e_{\beta\beta}(u) + \mu_f e_{\alpha\beta}(u) e_{\alpha\beta}(u) \right] dx \\
 &\quad + \varepsilon^{-2} \int_{\Omega_f} \left[\lambda_f e_{\alpha\alpha}(u) e_{33}(u) + 2\mu_f e_{\alpha 3}(u) e_{\alpha 3}(u) \right] dx \\
 &\quad + \varepsilon^{-4} \int_{\Omega_f} \frac{\lambda_f + 2\mu_f}{2} e_{33}(u) e_{33}(u) dx, \\
 J_\varepsilon^b(u) &:= \varepsilon^2 \int_{\Omega_b} \left[\frac{\lambda_b}{2} e_{\alpha\alpha}(u) e_{\beta\beta}(u) + \mu_b e_{\alpha\beta}(u) e_{\alpha\beta}(u) \right] dx \\
 &\quad + \int_{\Omega_b} \left[\lambda_b e_{\alpha\alpha}(u) e_{33}(u) + 2\mu_b e_{\alpha 3}(u) e_{\alpha 3}(u) \right] dx \\
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 &\quad + \int_{\Omega_b} \left[\lambda_b e_{\alpha\alpha}(u) e_{33}(u) + 2\mu_b e_{\alpha 3}(u) e_{\alpha 3}(u) \right] dx \\
 &\quad + \varepsilon^{-2} \int_{\Omega_b} \frac{\lambda_b + 2\mu_b}{2} e_{33}(u) e_{33}(u) dx \Rightarrow u_3 = 0
 \end{aligned}$$

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 J_\varepsilon^b(u) &:= \varepsilon^2 \int_{\Omega_b} \left[\frac{\lambda_b}{2} e_{\alpha\alpha}(u) e_{\beta\beta}(u) + \mu_b e_{\alpha\beta}(u) e_{\alpha\beta}(u) \right] dx \\
 &\quad + \int_{\Omega_b} \left[\lambda_b e_{\alpha\alpha}(u) e_{33}(u) + 2\mu_b e_{\alpha 3}(u) e_{\alpha 3}(u) \right] dx \xrightarrow{\substack{\int_{-1}^0 \partial_3 u_\alpha dx_3 \sim u_\alpha|_{\Omega_f} \\ \int_{-1}^0 \partial_\alpha u_3 dx_3 \sim 0}} \\
 &\quad + \varepsilon^{-2} \int_{\Omega_b} \frac{\lambda_b + 2\mu_b}{2} e_{33}(u) e_{33}(u) dx \Rightarrow u_3 = 0
 \end{aligned}$$

Theorem (B.-Henao)

J_ε “ Γ -converges” to J_0 , where for all $\bar{u} \in H^1(\omega; \mathbb{R}^2)$,

$$J_0(\bar{u}) = \underbrace{\int_{\omega} \left[\frac{\lambda_f \mu_f}{\lambda_f + 2\mu_f} e_{\alpha\alpha}(\bar{u}) e_{\beta\beta}(\bar{u}) + \mu_f e_{\alpha\beta}(\bar{u}) e_{\alpha\beta}(\bar{u}) \right] dx'}_{\text{usual reduced energy (Ciarlet)}} + \underbrace{\frac{\mu_b}{2} \int_{\omega} |\bar{u}|^2 dx'}_{\text{bonding energy}}.$$

- ✓ Limit displacements are not only of Kirchhoff-Love type, but actually planar (condition on the substrate);
- ✓ Recovery sequence (if $\bar{u} \in H^2(\omega; \mathbb{R}^2)$) :

$$u_\varepsilon(x', x_3) = \begin{cases} (\bar{u}(x'), 0) - \varepsilon^2 x_3 \frac{\lambda_f}{\lambda_f + 2\mu_f} e_{\alpha\alpha}(\bar{u})(x') \vec{e}_3 & \text{if } x \in \Omega_f, \\ (x_3 + 1)(\bar{u}(x'), 0) & \text{if } x \in \Omega_b, \\ 0 & \text{if } x \in \Omega_s. \end{cases}$$

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For all $u \in SBD(\Omega_f)$,

$$J_\varepsilon^f(u) := \int_{\Omega_f} \mathbb{C}_f e^\varepsilon(u) : e^\varepsilon(u) dx + \kappa_f \int_{J_u \cap \Omega_f} |((\nu_u)', \varepsilon^{-1}(\nu_u)_3)| d\mathcal{H}^2$$

Compactness (Bellettini-Coscia-Dal Maso) : let $(u_\varepsilon) \subset SBD(\Omega_f)$ with $\|u_\varepsilon\|_\infty \leq C$ and $J_\varepsilon^f(u_\varepsilon) \leq C$, then " $u_\varepsilon \rightharpoonup u$ in $SBD(\Omega_f)$ " where

$$e_{i3}(u) = 0 \text{ in } \Omega_f, \quad (\nu_u)_3 = 0 \text{ on } J_u \cap \Omega_f \Rightarrow E_{33}u = D_3u_3 = 0.$$

✓ Contrary to pure elasticity or plasticity u is not of Kirchhoff-Love type :

$$E_{\alpha 3}u = \frac{[u]_3(\nu_u)_\alpha}{2} \mathcal{H}^2 \llcorner J_u \neq 0.$$

For all $u \in SBD(\Omega_f)$,

$$\begin{aligned} J_\varepsilon^f(u) := & \int_{\Omega_f} \left[\frac{\lambda_f}{2} e_{\alpha\alpha}(u) e_{\beta\beta}(u) + \mu_f e_{\alpha\beta}(u) e_{\alpha\beta}(u) \right] dx \\ & + \varepsilon^{-2} \int_{\Omega_f} \left[\lambda_f e_{\alpha\alpha}(u) e_{33}(u) + 2\mu_f e_{\alpha 3}(u) e_{\alpha 3}(u) \right] dx \\ & + \varepsilon^{-4} \int_{\Omega_f} \frac{\lambda_f + 2\mu_f}{2} e_{33}(u) e_{33}(u) dx + \kappa_f \int_{J_u \cap \Omega_f} |((\nu_u)', \varepsilon^{-1}(\nu_u)_3)| d\mathcal{H}^2. \end{aligned}$$

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$$u_3 \in SBV(\omega), \quad \nabla u_3 \in SBD(\omega), \quad \bar{u} := \int_0^1 (u_1(\cdot, x_3), u_2(\cdot, x_3)) dx_3 \in SBD(\omega),$$

$$u_\alpha(x) = \bar{u}_\alpha(x') + \left(\frac{1}{2} - x_3 \right) \partial_\alpha u_3(x'), \quad J_u = (J_{\bar{u}} \cup J_{u_3} \cup J_{\nabla u_3}) \times (0, 1),$$

and J_ε^f Γ -converges to

$$\begin{aligned} J_0(u) &= \int_\omega \left[\frac{\lambda_f \mu_f}{\lambda_f + 2\mu_f} e_{\alpha\alpha}(\bar{u}) e_{\beta\beta}(\bar{u}) + \mu_f e_{\alpha\beta}(\bar{u}) e_{\alpha\beta}(\bar{u}) \right] dx' \\ &\quad + \frac{1}{12} \int_\omega \left[\frac{\lambda_f \mu_f}{\lambda_f + 2\mu_f} e_{\alpha\alpha}(\nabla u_3) e_{\beta\beta}(\nabla u_3) + \mu_f e_{\alpha\beta}(\nabla u_3) e_{\alpha\beta}(\nabla u_3) \right] dx' \\ &\quad + \underbrace{\kappa_f \mathcal{H}^1(J_{\bar{u}} \cup J_{u_3} \cup J_{\nabla u_3})}_{\text{Griffith energy}}. \end{aligned}$$

✓ Recovery sequence (if \bar{u} is “smooth” enough) :

$$u_\varepsilon(x', x_3) = u(x) - \varepsilon^2 x_3 \frac{\lambda_f}{\lambda_f + 2\mu_f} e_{\alpha\alpha}(\bar{u})(x') \vec{e}_3.$$

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We assume that the original (3D) displacements are of the form

$$u(x) = u(x_1, x_3) \vec{e}_2, \quad \text{for all } (x_1, x_3) \in I \times (-\infty, 1),$$

then for all $u \in SBV(I \times (-\infty, 1))$ with $u = 0$ in $I \times (-\infty, -1)$,

$$J_\varepsilon^f(u) = \frac{\mu_f}{2} \int_{I \times (0,1)} (|\partial_1 u|^2 + \varepsilon^{-2} |\partial_3 u|^2) dx_1 dx_3 + \kappa_f \int_{J_u \cap [I \times (0,1)]} |((\nu_u)_1, \varepsilon^{-1}(\nu_u)_3)| d\mathcal{H}^1,$$

$$J_\varepsilon^b(u) = \frac{\mu_b}{2} \int_{I \times [-1,0]} (\varepsilon^2 |\partial_1 u|^2 + |\partial_3 u|^2) dx_1 dx_3 + \kappa_b \int_{J_u \cap [I \times [-1,0]]} |(\varepsilon(\nu_u)_1, (\nu_u)_3)| d\mathcal{H}^1,$$

and

Theorem (Leon Baldelli-B.-Bourdin-Henao-Maurini)

$J_\varepsilon = J_\varepsilon^f + J_\varepsilon^b$ ‘Γ-converges’ to J_0 , where for all $u \in SBV(I)$,

$$J_0(u) := \frac{\mu_f}{2} \int_I |u'|^2 dx_1 + \kappa_f \#(J_u) + \frac{\mu_b}{2} \int_{I \setminus \Delta} |u|^2 dx_1 + \underbrace{\kappa_b \mathcal{L}^1(\Delta)}_{\text{Delamination energy}},$$

with $\Delta := \{|u| > \sqrt{2\kappa_b/\mu_b}\}$.

In the general case, we conjecture that J_ε Γ -converges to

$$\begin{aligned} J_0(u) = & \int_{\omega} \left[\frac{\lambda_f \mu_f}{\lambda_f + 2\mu_f} e_{\alpha\alpha}(\bar{u}) e_{\beta\beta}(\bar{u}) + \mu_f e_{\alpha\beta}(\bar{u}) e_{\alpha\beta}(\bar{u}) \right] dx' \\ & + \frac{1}{12} \int_{\omega} \left[\frac{\lambda_f \mu_f}{\lambda_f + 2\mu_f} e_{\alpha\alpha}(\nabla u_3) e_{\beta\beta}(\nabla u_3) + \mu_f e_{\alpha\beta}(\nabla u_3) e_{\alpha\beta}(\nabla u_3) \right] dx' \\ & + \kappa_f \mathcal{H}^1(J_{\bar{u}} \cup J_{u_3} \cup J_{\nabla u_3}) + \frac{\mu_b}{2} \int_{\omega \setminus \Delta} |\bar{u}|^2 dx' + \kappa_b \mathcal{L}^2(\Delta), \end{aligned}$$

where $\Delta := \{|\bar{u}| > \sqrt{2\kappa_b/\mu_b}\} \cup \{u_3 \neq 0\}$ is the delamination set.

✓ The upper bound is true. Recovery sequence :

$$u_\varepsilon(x', x_3) = \begin{cases} u(x) - \varepsilon^2 x_3 \frac{\lambda_f}{\lambda_f + 2\mu_f} e_{\alpha\alpha}(\bar{u})(x') \vec{e}_3 & \text{if } x \in \Omega_f, \\ (x_3 + 1)(\bar{u}(x'), 0) & \text{if } x \in (\omega \setminus \Delta) \times [-1, 0], \\ 0 & \text{if } x \in (\Delta \times [-1, 0]) \cup \Omega_s, \end{cases}$$

where Δ and \bar{u} are suitably regularized (Chambolle).

✓ We get a too low lower bound (with a bad multiplicative constant), under additional compactness assumptions on minimizing sequences.

Thanks for your attention !