Liouville Equations and Functional Determinants

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• While physicists may like these formulas, mathematicians usually have problems with infinite products of diverging numbers.

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$$\zeta(s) = \sum_{j=1}^{\infty} \lambda_j^{-s}.$$

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If ζ is regular near s = 0 one can define the *regularized determinant* $det'(-\Delta_g)$ via the following formula

$$\det'(-\Delta_g) = e^{-\zeta'(0)}.$$

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$$\zeta(s) \ = \ \sum_{j=1}^\infty \lambda_j^{-s} = \frac{1}{\Gamma(s)} \int_0^\infty \left(\sum_{j=1}^\infty e^{-\lambda_j t} \right) t^s \frac{dt}{t}$$

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$$\zeta(s) = \frac{1}{\Gamma(s)} \left\{ \frac{A(\Sigma)}{4\pi(s-1)} + \left(\frac{\chi(\Sigma)}{6} - 1\right) + \text{holom. in } s \right\},$$

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which is regular near zero. $\Rightarrow \det'(-\Delta_g)$ is well defined. Andrea Malchiodi (SNS, Pisa) Banff, 04-03-2018

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Existence of extremals is easy for positive genus. On spheres it can be achieved via a *balancing condition*, done in [Osgood-Phillips-Sarnak, '88] (see also [Aubin, '76], [Ghoussoub-Lin, '10], [Gui-Moradifam, '16]).

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Expanding the heat kernel (via parametrix) one can prove that

$$\sum_{j=1}^{\infty} e^{-\lambda_j t} =: Tr(e^{\Delta t}) = \frac{1}{t} \sum_{j=0}^{l} t^j \int_{\Sigma} U_j(x) dV + O(t^l),$$

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$$U_j \simeq \int_{\Sigma} K_g \Delta^{j-2} K_g dV \simeq \|u\|_{H^j(\Sigma)}$$

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therefore one gets bounds even in higher Soboley norms = > (문 > 문 > 오이어 (SNS, Pisa) Banff, 04-03-2018 6 / 31

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It was however shown in [Wolpert, '87] that

$$\det'(\hat{g}) \le \frac{1}{l} e^{-\frac{c_1}{l}}; \qquad c_1 = c_1(\chi(\Sigma)),$$

where l is the length of the shortest geodesic, so $l \not\rightarrow 0$.

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Finally, a theorem in [Mumford, '71] shows that if l is bounded below and if $K_{\hat{g}} = const.$, then there is smooth convergence of the metrics.

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On S^2 it is known that the supremum of the k-th eigenvalue is $8\pi k$ ([Karpukhin-Nadirashvili-Penskoi-Polterovich '17]), with previous results in [Hersch, '70], [Petrides, '14], [Nadirashvili-Sire, '17] (k = 1, 2, 3).

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With few exceptions, no explicit formulas are known in higher genus.

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A linear operator $A = A_g$ is conformally covariant of bi-degree (a, b) if $\tilde{g} = e^{2w}g$ implies

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$$L_g = -\frac{4(n-1)}{(n-2)}\Delta_g + R_g \qquad (a,b) = \left(\frac{n-2}{2}, \frac{n+2}{2}\right)$$

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2. The Paneitz operator P_g for n = 4

$$P_g \varphi = (-\Delta_g)^2 \varphi + \operatorname{div} \left[\left(\frac{2}{3} Rg - 2Ric \right) \circ \nabla \varphi \right], \qquad (a,b) = (0,4).$$

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$$L_g = -\frac{4(n-1)}{(n-2)}\Delta_g + R_g \qquad (a,b) = \left(\frac{n-2}{2}, \frac{n+2}{2}\right).$$

2. The Paneitz operator P_g for n = 4

$$P_g \varphi = (-\Delta_g)^2 \varphi + \operatorname{div} \left[\left(\frac{2}{3} Rg - 2Ric \right) \circ \nabla \varphi \right], \qquad (a,b) = (0,4).$$

3. The Dirac operator \mathcal{D} for $n \ge 2$: $(a,b) = \left(\frac{n-1}{\sqrt{2}}, \frac{n+1}{\sqrt{2}}\right)$.

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Here W_g is Weyl's curvature, while Q_g is the *Q*-curvature, a 4D conformal counterpart of the Gaussian curvature.

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Also, in 4D there is a Gauss-Bonnet formula

$$\int_{M} \left(Q_g + \frac{1}{8} |W_g|^2 \right) dv = 4\pi^2 \chi(M).$$

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• Both P_g and Q_g have a crucial role in the study of the topology of 4-manifolds (works by Chang, Gursky, Yang, Qing, reg.), reg.

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- If $A_g = \mathcal{D}$, the square of the Dirac operator, then

$$(\gamma_1, \gamma_2, \gamma_3) = \left(-7, -88, -\frac{14}{3}\right).$$

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- Uniqueness holds for $-\gamma_1 \int_M |W_g|^2 dv \gamma_2 \int_M Q_g dv < 0.$
- The theorem applies to L_g and \mathcal{D} , but not to the Paneitz operator P_g (discussed later).

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The topological structure of the energy (with Struwe's monotonicity argument) allows to produce solutions of *perturbed equations*

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Finally, using an integration by parts (Pohozaev), one shows that $\beta_i = 8\pi^2$ for all i, a contradiction to $k_Q \notin 8\pi^2 \mathbb{N}$.

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Some results were available for the *p*-Laplacian ([Serrin, '64], [Veron-Kichenassamy, '86]), but for that one has homogeneity of the operator, plus the maximum principle.

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Nonlinear Green's function

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Nonlinear Green's function

The natural space to work with variationally is $W^{2,2}$. However, this is not possible for singular solutions.

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For the regularity, one can use an approximate solution u_{app} of the form

$$u_{\text{app}}(x) \simeq \sum_{i=1}^{l} \alpha_i \log \frac{1}{d(x, p_i)}; \qquad \alpha_i = \alpha_i(\beta_i).$$

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This implies the desired regularity on the original closed manifold.

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$$F_P[w] = \int_{\mathbb{T}^4} \left[18(\Delta w)^2 + 64 |\nabla w|^2 \Delta w + 32 |\nabla w|^4 \right]$$

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This functional has a <u>triple homogeneity</u>. Moreover, by the borderline embedding $W^{2,2}(\mathbb{R}^4) \hookrightarrow \overline{W^{1,4}(\mathbb{R}^4)}$, it is also <u>doubly critical</u>.

On S^4 instead one has

$$F_{P}[w] = \int_{S^{4}} \left[18(\Delta w)^{2} + 64|\nabla w|^{2}\Delta w + 32|\nabla w|^{4} - 60|\nabla w|^{2} \right] dv + 112\pi^{2} \log \left(\int_{S^{4}} e^{4(w-\overline{w})} dv \right).$$

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Proposition 1

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- Geometrically, this conformal factor generates cylinder, not a bubble.

- Loss of coercivity may happen in *different ways* (e.g., at many points), differently e.g. from the Q-curvature equation.
- It goes similarly with compact hyperbolic manifolds.

A second solution on S^4

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A second solution on S^4

Theorem ([Gursky-M., '12])

Let (S^4, g_0) be the standard 4-sphere. Then F_P admits a non-trivial axially symmetric solution.

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Let (S^4, g_0) be the standard 4-sphere. Then F_P admits a non-trivial axially symmetric solution.

Remarks (a) For most geometric problems the round metric is *the only critical point*. One has indeed uniqueness of the round metric for constant mean curvature, Gaussian curvature, scalar curvature and Q-curvature.

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(d) A similar result holds in \mathbb{R}^4 , much easier to prove.
A convenient change of variables

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• We can also assume that u(t) is even in t.

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With the above change of variables one finds the equation

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• $u_0(t) = -\log \cosh(t)$ represents the standard spherical metric

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• $u_0(t) = -\log \cosh(t)$ represents the standard spherical metric

• if $u(t) \to -\infty$ as $t \to +\infty$, then u(t) shadows a solution of

$$(E_{\infty}) \qquad \qquad 9v'''' - 96v''(v')^2 + 60v'' = 0.$$

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With the above change of variables one finds the equation

(E)
$$9u'''' - 96u''(u')^2 + 60u'' + 42e^{4u} = 0.$$

Evenness in t implies u'(0) = u'''(0) = 0. We also require

$$u'(t) \to -1, u''(t) \to 0$$
 as $t \to +\infty$ $(u(t)$ extends to $S^4)$;

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This is integrable, with a one-parameter family of periodic solutions.Andrea Malchiodi (SNS, Pisa)Banff, 04-03-201822 / 31

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Integrating by parts one finds the following result. **Proposition 2** Admissible solutions of (E) satisfy

$$-\frac{9}{2}[u''(0)]^2 + \frac{21}{2}e^{4u(0)} = 6$$

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and also the equation

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• By the first formula, the initial conditions are completely determined by u''(0) (recall that u'(0) = u'''(0) = 0).

• The second formula reduces (E) to a third order, <u>autonomous</u> equation in u' (the exponential term disappears).

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Setting

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the system becomes

(A)
$$\begin{cases} x' = y, \\ y' = z, \\ z' = \frac{8}{3}(x^2 - 1)(4x^2 - 1) - 4xz + \frac{32}{3}x^2y + 2y^2 - \frac{20}{3}y. \end{cases}$$

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Proposition 3

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Proposition 3 System (A) contains both solutions of (E) and (E_{∞}) .

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Proposition 3 System (A) contains both solutions of (E) and (E_{∞}) . Thanks to this (miracolous) result the asymptotics of the solutions of (E) can be made rigorous.

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Proposition 3 System (A) contains both solutions of (E) and (E_{∞}) .

Thanks to this (miracolous) result the asymptotics of the solutions of (E) can be made rigorous.

<u>Goal</u> Look for solutions of (A) starting from the *y*-axis and converging asymptotically to the point (1, 0, 0).

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Fact System (A) contains a one-parameter family of periodic orbits forming a topological disk \mathfrak{D} .

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The *center* is the point $p_0 = (\frac{1}{2}, 0, 0)$

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The *center* is the point $p_0 = (\frac{1}{2}, 0, 0)$, while the most external orbit is a homoclinic, with limit point $p_1 = (1, 0, 0)$.

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The *center* is the point $p_0 = (\frac{1}{2}, 0, 0)$, while the most external orbit is a homoclinic, with limit point $p_1 = (1, 0, 0)$.

• The transversal dynamics is attractive

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Shooting method

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Shooting method

Recall that the spherical metric corresponds to $u(t) = -\log \cosh t$.

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$$\overrightarrow{X}_0(0) = (0, 1, 0).$$

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Let us try now to vary the initial data, hoping to find another admissible solution.

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$$\vec{X}_0(0) = (0, 1, 0).$$

Let us try now to vary the initial data, hoping to find another admissible solution.

For $\varepsilon > 0$, let $\overrightarrow{X}_{\varepsilon}(t)$ be the solution of (A) with initial data

$$\overrightarrow{X}_{\varepsilon}(0) = (0, 1 - \varepsilon, 0).$$

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Proposition 4 For $\varepsilon > 0$ small enough, $\overrightarrow{X}_{\varepsilon}(t)$ is globally defined in time

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Proposition 4 For $\varepsilon > 0$ small enough, $\overrightarrow{X}_{\varepsilon}(t)$ is globally defined in time, and shadows one of the periodic orbits in \mathfrak{D} .

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The proof uses refined asymptotic analysis, a Gronwall inequality and the construction of two (sort of) Lyapunov functions.

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Technically, one needs to rule out infinitely-many oscillations.

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Comments and open problems

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Comments and open problems

Our proof is very specific and does not exploit the structure of the problem.

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$$F_P[w] = \int_{\mathbb{T}^4} \left[18(\Delta w)^2 + 64 |\nabla w|^2 \Delta w + 32 |\nabla w|^4 \right].$$

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$$\inf_{u \neq 0} \frac{\int_{\mathbb{T}^4} (\Delta u)^2 dx}{\left(\int_{\mathbb{T}^4} |\nabla u|^4 dx\right)^{\frac{1}{2}}}.$$

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It is an interesting problem to find extremals of this quotient in \mathbb{R}^4 .

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On \mathbb{R}^4 critical points satisfy

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A natural question is whether a critical point always exists for F_P . This is be a natural counterpart of the Uniformization problems or the Yamabe problem. Apart from the compactness issues, new sharp Moser-Trudinger inequalities would be expected.

Thanks for your attention

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