# Liouville Equations and Functional Determinants 

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- While physicists may like these formulas, mathematicians usually have problems with infinite products of diverging numbers.


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If $\zeta$ is regular near $s=0$ one can define the regularized determinant $\operatorname{det}^{\prime}\left(-\Delta_{g}\right)$ via the following formula

$$
\operatorname{det}^{\prime}\left(-\Delta_{g}\right)=e^{-\zeta^{\prime}(0)}
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Existence of extremals is easy for positive genus. On spheres it can be achieved via a balancing condition, done in [Osgood-Phillips-Sarnak, '88] (see also [Aubin, '76], [Ghoussoub-Lin, '10], [Gui-Moradifam, '16]).

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Finally, a theorem in [Mumford, '71] shows that if $l$ is bounded below and if $K_{\hat{g}}=$ const., then there is smooth convergence of the metrics.

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On $S^{2}$ it is known that the supremum of the $k$-th eigenvalue is $8 \pi k$ ([Karpukhin-Nadirashvili-Penskoi-Polterovich '17]), with previous results in [Hersch , '70], [Petrides, '14], [Nadirashvili-Sire, '17] $(k=1,2,3)$.

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With few exceptions, no explicit formulas are known in higher genus.

## Conformally covariant operators

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2. The Paneitz operator $P_{g}$ for $n=4$

$$
P_{g} \varphi=\left(-\Delta_{g}\right)^{2} \varphi+\operatorname{div}\left[\left(\frac{2}{3} R g-2 R i c\right) \circ \nabla \varphi\right], \quad(a, b)=(0,4)
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2. The Paneitz operator $P_{g}$ for $n=4$

$$
P_{g} \varphi=\left(-\Delta_{g}\right)^{2} \varphi+\operatorname{div}\left[\left(\frac{2}{3} R g-2 R i c\right) \circ \nabla \varphi\right], \quad(a, b)=(0,4)
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3. The Dirac operator $\mathcal{D}$ for $n \geq 2:(a, b)=\left(\frac{n-1}{2}, \frac{n+1}{2}\right)$.

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- Both $P_{g}$ and $Q_{g}$ have a crucial role in the study of the topology of 4-manifolds (works by Chang, Gursky, Yang, Qing, ョ: )


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- The theorem applies to $L_{g}$ and $\mathcal{D}$, but not to the Paneitz operator $P_{g}$ (discussed later).


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A consequence of these improved inequalities is that, for example, if $k_{Q} \in\left(8 \pi^{2}, 16 \pi^{2}\right)$ and if $F_{A}$ is large, then the conformal volume must concentrate near a single point of $M$. One can then exploit the topology of $M$ to find critical point of $F_{A}$ of saddle type.

## Min-max methods and compactness

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The topological structure of the energy (with Struwe's monotonicity argument) allows to produce solutions of perturbed equations

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Finally, using an integration by parts (Pohozaev), one shows that $\beta_{i}=$ $8 \pi^{2}$ for all $i$, a contradiction to $k_{Q} \notin 8 \pi^{2} \mathbb{N}$.

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Some results were available for the $p$-Laplacian ([Serrin, '64], [VeronKichenassamy, '86]), but for that one has homogeneity of the operator, plus the maximum principle.

## Nonlinear Green's function

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For the regularity, one can use an approximate solution $u_{\text {app }}$ of the form

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## The determinant of the Paneitz operator

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On $S^{4}$ instead one has

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\begin{aligned}
F_{P}[w] & =\int_{S^{4}}\left[18(\Delta w)^{2}+64|\nabla w|^{2} \Delta w+32|\nabla w|^{4}-60|\nabla w|^{2}\right] d v \\
& +112 \pi^{2} \log \left(f_{S^{4}} e^{4(w-\bar{w})} d v\right)
\end{aligned}
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- It goes similarly with compact hyperbolic manifolds.


## A second solution on $S^{4}$

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(c) The mountain pass structure suggests to use a variational approach. However this strategy is now out of reach: we used ODEs instead.
(d) A similar result holds in $\mathbb{R}^{4}$, much easier to prove.

## A convenient change of variables

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Once the north and south poles are removed, $S^{4}$ is conformally equivalent to the cylinder $S^{3} \times \mathbb{R}$


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Asking for axial symmetry is equivalent to having solutions independent of the $S^{3}$ component. Therefore we just solve for $u=u(t), t \in \mathbb{R}$.

- We can also assume that $u(t)$ is even in $t$.


## The ODE on the cylinder

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With the above change of variables one finds the equation
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This is integrable, with a one-parameter family of periodic solutions.

## Conservation laws

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- The second formula reduces $(E)$ to a third order, autonomous equation in $u^{\prime}$ (the exponential term disappears).


## The autonomous system

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x(t)=-u^{\prime}(t) ; \quad y(t)=-u^{\prime \prime}(t) ; \quad z(t)=-u^{\prime \prime \prime}(t)
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the system becomes
(A) $\left\{\begin{array}{l}x^{\prime}=y, \\ y^{\prime}=z, \\ z^{\prime}=\frac{8}{3}\left(x^{2}-1\right)\left(4 x^{2}-1\right)-4 x z+\frac{32}{3} x^{2} y+2 y^{2}-\frac{20}{3} y .\end{array}\right.$

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Thanks to this (miracolous) result the asymptotics of the solutions of $(E)$ can be made rigorous.
Goal Look for solutions of $(A)$ starting from the $y$-axis and converging asymptotically to the point $(1,0,0)$.

## An invariant set for $(A)$

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- The transversal dynamics is attractive


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Let us try now to vary the initial data, hoping to find another admissible solution.
For $\varepsilon>0$, let $\vec{X}_{\varepsilon}(t)$ be the solution of $(A)$ with initial data

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The proof uses refined asymptotic analysis, a Gronwall inequality and the construction of two (sort of) Lyapunov functions.

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Technically, one needs to rule out infinitely-many oscillations.

## Comments and open problems

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It is an interesting problem to find extremals of this quotient in $\mathbb{R}^{4}$.

## The Euler equation

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On $\mathbb{R}^{4}$ critical points satisfy

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A natural question is whether a critical point always exists for $F_{P}$. This is be a natural counterpart of the Uniformization problems or the Yamabe problem. Apart from the compactness issues, new sharp Moser-Trudinger inequalities would be expected.

# Thanks for your attention 

