

# Some canonical metrics on four manifolds: rigidity and existence

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# Canonical Riemannian metrics

*"It is geometers dream to find a canonical metric  $g_{best}$  on a given smooth manifold  $M$  so that all topology of  $M$  will be captured by geometry." [M. Gromov]*

$(M^n, g)$  smooth Riemannian manifold,  $\dim(M^n) = n \geq 2$ ,  $\partial M = \emptyset$ .

$g$  metric  $\rightsquigarrow$  Riemann ( $Riem_g$ ), Ricci ( $Ric_g$ ) and scalar curvature ( $R_g$ )

In coordinates:

$$Riem_g = (Riem)_{ijkl} \xrightarrow{\text{trace}} Ric_g = (Ric)_{ik} = g^{jl} (Riem)_{ijkl} \xrightarrow{\text{trace}} R_g = g^{ik} (Ric)_{ik}$$

- Constant curvature:  $Riem_g = \lambda g \otimes g$       Space forms
- Constant Ricci curvature:  $Ric_g = \lambda g$       Einstein metrics
- Constant scalar curvature:  $R_g = \lambda$       Yamabe metrics

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# Einstein metrics I: variational point of view

Canonical metrics as **critical points** of curvature functionals. Let  $M$  closed,

$$\mathfrak{S}(g) = \text{Vol}_g(M)^{-\frac{n-2}{n}} \int_M R_g dV_g \quad \text{Einstein-Hilbert functional}$$

$g$  is **critical** for  $\mathfrak{S}(g) \iff g$  is **Einstein**, i.e.  $\text{Ric}_g = \lambda g$ ,  $\lambda \in \mathbb{R}$ .

- If  $n = 3$ , then Einstein metrics have constant curvature.
- If  $n = 4$ , it is well known that there are topological obstructions to the existence of an Einstein metric (e.g. Hitchin-Thorpe:  $\chi(M) \geq \frac{3}{2}|\tau(M)|$ ).
- If  $n > 4$ , still unknown.

On the other hand, the constrained problem in a conformal class (Yamabe problem) is unobstructed. More precisely the **Yamabe invariant**

$$\mathcal{Y}(M, [g]) = \inf_{\tilde{g} \in [g]} \mathfrak{S}(\tilde{g}) = \frac{4(n-1)}{n-2} \inf_{u \in W^{1,2}(M)} \frac{\int_M |\nabla u|^2 dV_g + \frac{n-2}{4(n-1)} \int_M R u^2 dV_g}{\left(\int_M |u|^{2n/(n-2)} dV_g\right)^{(n-2)/n}}$$

is always attained in every conformal class  $[g]$ .

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## Einstein metrics II: the Weyl tensor

Let  $(M^n, g)$  be an Einstein manifold, i.e.  $Ric_g = \lambda g$ , for some  $\lambda \in \mathbb{R}$  [Besse]. In particular, by tracing, the scalar curvature is constant  $R_g = n\lambda$ . By the decomposition of the curvature tensor

$$Riem_g = \frac{R_g}{2n(n-1)} (g \otimes g) + Weyl_g$$

So, all the *geometric* information of an Einstein manifold are contained in the **Weyl tensor**. In particular, if  $Weyl_g = 0$ , then  $(M^n, g)$  is a space form.

Some well known facts:

- In dimension  $n = 3$   $Weyl_g = 0$ , so every Einstein manifold is a space form.
- The Weyl tensor is **totally trace free**, i.e.  $g^{ik} W_{ijkl} = 0$ .
- Tracing the second Bianchi identity for  $Riem_g$ :  $\nabla_t R_{ijkl} + \nabla_l R_{ijtk} + \nabla_k R_{ijlt} = 0$  and using the decomposition, we get that the Weyl tensor has **zero divergence**, i.e.  $\nabla_t W_{ijkt} = 0$  (**harmonic Weyl curvature**).
- With some work, one can show that the Weyl tensor satisfies the **second Bianchi Identity**

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# Einstein metrics III: Bochner-Weitzenböck formula

Proposition (Derdzinski, '83)

Let  $(M^4, g)$  be a four dimensional Einstein manifold. Then

$$\Delta W_{ijkl} = \frac{1}{2} R_g W_{ijkl} - 4W_{ipkq} W_{pjql} - W_{klpq} W_{pqij}.$$

In particular,

$$\frac{1}{2} \Delta |Weyl_g|^2 = |\nabla Weyl_g|^2 + \frac{1}{2} R_g |Weyl_g|^2 - 3W_{ijkl} W_{ijpq} W_{klpq}.$$

More in general, the last formula holds on four-manifolds with harmonic Weyl curvature.

**Proof:** Taking the divergence of the second Bianchi identity for  $Weyl_g$  and commuting, we get

$$\begin{aligned} 0 &= \nabla_t \nabla_t W_{ijkl} + \nabla_t \nabla_l W_{ijtk} + \nabla_t \nabla_k W_{ijlt} \\ &= \Delta W_{ijkl} + \nabla_l \nabla_t W_{ijtk} + \nabla_k \nabla_t W_{ijlt} + Riem_g * Weyl_g \\ &= \Delta W_{ijkl} + Riem_g * Weyl_g \end{aligned}$$

Since  $Riem_g = (R_g/24)(g \otimes g) + Weyl_g$  and  $Weyl_g$  is trace free  $\implies \square$

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Let  $(M^4, g)$  be a four dimensional Einstein manifold. Then

$$\Delta W_{ijkl} = \frac{1}{2} R_g W_{ijkl} - 4W_{ipkq} W_{pjql} - W_{klpq} W_{pqij}.$$

In particular,

$$\frac{1}{2} \Delta |Weyl_g|^2 = |\nabla Weyl_g|^2 + \frac{1}{2} R_g |Weyl_g|^2 - 3W_{ijkl} W_{ijpq} W_{klpq}.$$

More in general, the last formula holds on four-manifolds with harmonic Weyl curvature.

**Proof:** Taking the divergence of the second Bianchi identity for  $Weyl_g$  and commuting, we get

$$\begin{aligned} 0 &= \nabla_t \nabla_t W_{ijkl} + \nabla_t \nabla_l W_{ijtk} + \nabla_t \nabla_k W_{ijlt} \\ &= \Delta W_{ijkl} + \nabla_l \nabla_t W_{ijtk} + \nabla_k \nabla_t W_{ijlt} + Riem_g * Weyl_g \\ &= \Delta W_{ijkl} + Riem_g * Weyl_g \end{aligned}$$

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# Bochner type formula of high order

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Thus, under suitable assumptions, one can deduce Liouville type results for this class of PDEs. To this aim we computed a **Bochner-Weitzenböck formula** for the covariant derivatives of the Weyl tensor, obtaining:

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Let  $(M^4, g)$  be a four dimensional Einstein manifold. Then,

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By standard commutation rules, it is quite easy to derive “rough” Bochner type identity for the covariant derivative of Weyl, and, with some work, even a formula for the  $k$ -th covariant derivative  $\nabla^k W$ .

The proof of the theorem, instead, relies heavily on the *algebraic structure* of curvature operators in dimension four. In fact, on an oriented Riemannian manifold of dimension four  $(M^4, g)$ ,  $\Lambda^2$  decomposes as the sum of two subbundles  $\Lambda^2 = \Lambda^+ \oplus \Lambda^-$ , which are the eigenspaces of the Hodge operator  $\star : \Lambda^2 \rightarrow \Lambda^2$  corresponding respectively to the eigenvalue  $\pm 1$ . Since the Weyl tensor acts on  $\Lambda^2$ , we have the decomposition

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The proof of the theorem, instead, relies heavily on the *algebraic structure* of curvature operators in dimension four. In fact, on an oriented Riemannian manifold of dimension four  $(M^4, g)$ ,  $\Lambda^2$  decomposes as the sum of two subbundles  $\Lambda^2 = \Lambda^+ \otimes \Lambda^-$ , which are the eigenspaces of the Hodge operator  $\star : \Lambda^2 \rightarrow \Lambda^2$  corresponding respectively to the eigenvalue  $\pm 1$ . Since the Weyl tensor acts on  $\Lambda^2$ , we have the decomposition

$$Weyl_g = W_g^+ + W_g^-$$

where the **self-dual** and **anti-self-dual**  $W^\pm$  are trace-free endomorphisms of  $\Lambda^\pm$ .

# Einstein metrics VI: applications

As a consequence, we can show the following:

## Corollary 1 (C.-Mastrolia)

Let  $(M^4, g)$  be a closed four dimensional Einstein manifold. Then

$$\int |\nabla^2 W_g^\pm|^2 - \frac{5}{3} \int |\Delta W_g^\pm|^2 + \frac{1}{4} R_g \int |\nabla W_g^\pm|^2 = 0,$$

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Let  $(M^4, g)$  be a four dimensional Einstein manifold with positive scalar curvature. If

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# Harmonic Weyl curvature

$$\mathcal{SF} \subset \mathcal{E} \subset \mathcal{Y}$$
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- Einstein metrics have harmonic Weyl curvature.
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$$\frac{1}{2} \Delta |W_g|^2 = |\nabla W_g|^2 + \frac{1}{2} R_g |W_g|^2 - 3 W_{ijkl} W_{ijpq} W_{klpq}.$$

Actually this formula **characterizes** harmonic Weyl metrics on closed four manifolds. This follows from the integral identity [Chang-Gursky-Yang]

$$\int_M (|\nabla W_g|^2 - 4|\delta W_g|^2 + \frac{1}{2} R_g |W_g|^2 - 3 W_{ijkl} W_{ijpq} W_{klpq}) dV_g = 0$$

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# A new variational problem

From a variational point of view it seems natural to consider the quadratic scaling-invariant Riemannian functional

$$\mathfrak{D}(g) := \text{Vol}_g(M)^{\frac{1}{2}} \int_M |\delta_g W_g|_g^2 dV_g$$

Obviously harmonic Weyl metrics are critical points (absolute minima) of  $\mathfrak{D}(g)$ . In the same spirit of the Yamabe problem, we define the conformal invariant

$$\mathcal{D}(M, [g]) := \inf_{\tilde{g} \in [g]} \mathfrak{D}(\tilde{g})$$

## Questions:

1. What are the geometric properties of critical metrics in the conformal class for the functional  $g \mapsto \mathfrak{D}(g)$ ?
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# Weak harmonic Weyl metrics

We have the following characterization of critical metrics in the conformal class for the functional

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A metric is critical in the conformal class for the functional  $g \mapsto \mathfrak{D}(g)$  if and only if it satisfies the Weitzenböck formula

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In this case we say that  $g$  is a **weak harmonic Weyl metric**.

By Derdzinski formula, harmonic Weyl implies weak harmonic Weyl.

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# Existence of WHW metrics

## Theorem 2 (C.-Mastrolia-Monticelli-Punzo)

On every closed four-dimensional manifold there exists a weak harmonic Weyl metric.

- Aubin proved that every closed Riemannian manifold admits a constant negative scalar curvature metric. Besides this one, Theorem 2 is the only existence result of a *canonical* metric, which generalizes the Einstein condition, on every four-dimensional Riemannian manifold, without any topological obstructions.
- The metric in Theorem 2 is constructed as follows: first, thanks to a result of Aubin, on every four-dimensional manifold  $M^4$  we can choose a reference metric  $g_0$  with  $|W_{g_0}|_{g_0} > 0$ . Then, we prove that on  $(M^4, g_0)$  the infimum  $\mathcal{D}(M, [g_0])$  is attained by a conformal metric  $g \in [g_0]$ , which is a weak harmonic Weyl metric. Moreover, we show that every critical point in the conformal class  $[g_0]$  is necessarily a minimum point.

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On every closed four-dimensional manifold there exists a weak harmonic Weyl metric.

- Aubin proved that every closed Riemannian manifold admits a constant negative scalar curvature metric. Besides this one, Theorem 2 is the only existence result of a *canonical* metric, which generalizes the Einstein condition, on every four-dimensional Riemannian manifold, without any topological obstructions.
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In order to prove this theorem, we endow a closed four-manifolds  $M^4$  with the metric  $g_0$  constructed by Aubin and we consider the functional

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where all the geometric quantities are referred to  $g_0$  and the function  $v$  belongs to the convex cone

$$H(M) := \left\{ u \in H^1(M) : u > 0 \text{ a.e. and } \int_M u^{-4} dV < \infty \right\}.$$

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One has

$$\mathfrak{D}(v) = \left( \int_M v^{-4} dV \right)^{\frac{1}{2}} \int_M (a|\nabla v|^2 + c v^2) dV = \left( \int_M v^{-4} dV \right)^{\frac{1}{2}} \int_M v Lv dV$$

with  $a \in C^\infty(M)$ ,  $a > 0$ ,  $c \in C^\infty(M)$  and the uniformly elliptic self-adjoint operator  $L$  is given by

$$Lv := -\operatorname{div}(a \nabla v) + c v.$$

Since, by definition,  $\mathfrak{D}(v) \geq 0$ , we get

$$\lambda_1 := \inf_{u \in H^1(M), u \neq 0} \frac{\int_M (a|\nabla u|^2 + c u^2) dV}{\int_M u^2 dV} \geq 0.$$

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## Sketch of the proof III

We have the following (strong) maximum principle.

### Lemma 1

Let  $\lambda_1 > 0$ . If  $u \in H^1(M)$  satisfies  $Lu \geq 0$  in the weak sense, then either  $u = 0$  a.e. on  $M$  or  $\operatorname{ess\,inf}_M u > 0$ .

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$$\operatorname{Vol}(M)^{\frac{3}{2}} \lambda_1 \leq \mathcal{D} \leq \frac{\int_M \varphi_1^2 dV}{\left(\int_M \varphi_1^{-4} dV\right)^{\frac{1}{2}}} \lambda_1.$$

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**Conclusion:** if we choose a reference metric  $g_0$  with  $|W_{g_0}|_{g_0} > 0$ , then we can find a conformal metric  $g = v^{-2} g_0$ ,  $v \in C^\infty(M)$ , minimizing the functional  $\mathfrak{D}(g)$ .



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## Degenerate case

What happens in the degenerate case, i.e. if  $|W_{g_0}|_{g_0} = 0$  somewhere in  $M$ ?

We can show that uniqueness (up to scaling) of smooth ( $C^2$ ) solutions to the equation still holds, unless  $g_0$  is locally conformally flat, i.e.  $W_{g_0} \equiv 0$ . Moreover we have the following non-existence results:

- If  $|W_{g_0}|_{g_0} \equiv 0$  on some open set  $\Omega \subset M$ , then we can show that a smooth metric  $g = v^{-2}g_0$  is critical if and only if it is locally conformally flat, i.e.  $W_g \equiv 0$  on  $M$  (and thus  $W_{g_0} \equiv 0$ ).
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## Degenerate case

What happens in the degenerate case, i.e. if  $|W_{g_0}|_{g_0} = 0$  somewhere in  $M$ ?

We can show that uniqueness (up to scaling) of smooth ( $C^2$ ) solutions to the equation still holds, unless  $g_0$  is locally conformally flat, i.e.  $W_{g_0} \equiv 0$ . Moreover we have the following non-existence results:

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Thank you.

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