

Self-Adjoint Boundary Conditions for Singular Sturm–Liouville Operators and the computation of m -functions for Bessel, Legendre, and Laguerre operators

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**Herglotz–Nevanlinna Theory Applied to Passive, Causal
and Active Systems**

Banff International Research Station, Workshop, October 6–11, 2019

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Motivation

Try to extend the classical boundary values

$$g(a) = -W(u_a(\lambda_0, \cdot), g)(a) = \lim_{x \downarrow a} \frac{g(x)}{\widehat{u}_a(\lambda_0, x)}, \quad (*)$$

$$g^{[1]}(a) = (pg')(a) = W(\widehat{u}_a(\lambda_0, \cdot), g)(a) = \lim_{x \downarrow a} \frac{g(x) - g(a)\widehat{u}_a(\lambda_0, x)}{u_a(\lambda_0, x)}, \quad (**)$$

for **regular Sturm–Liouville operators** on $[a, b] \subset \mathbb{R}$ associated with differential expressions of the type

$$\tau = r(x)^{-1}[-(d/dx)p(x)(d/dx) + q(x)] \text{ for a.e. } x \in [a, b] \subset \mathbb{R},$$

to the case where τ is singular on $(a, b) \subseteq \mathbb{R}$ and the associated **minimal operator** T_{min} is **bounded from below**.

Here $u_a(\lambda_0, \cdot)$ and $\widehat{u}_a(\lambda_0, \cdot)$ denote appropriately normalized **principal** and **nonprincipal** solutions of $\tau u = \lambda_0 u$ for appropriate $\lambda_0 \in \mathbb{R}$, respectively.

While the l.h.s. in $(*)$, $(**)$ above will cease to be meaningful in the singular case, it will be shown that the r.h.s. remains valid!

Some Literature

Based on:

F.G., L. Littlejohn, and R. Nichols, *A note on self-adjoint boundary conditions for singular Sturm–Liouville operators*, in preparation.

Also relies on:

H.-D. Niessen and A. Zettl, *Singular Sturm–Liouville problems: the Friedrichs extension and comparison of eigenvalues*, Proc. London Math. Soc. (3) **64**, 545–578 (1992).

All Self-Adjoint B.C.'s in the Regular Case

All **separated** and **coupled** boundary conditions together describe **all self-adjoint** extensions of T_{min} :

Theorem.

Assume that τ is **regular** on $[a, b]$. Then the following items (i)–(iii) hold:

(i) **All self-adjoint** extensions $T_{\alpha, \beta}$ of T_{min} with **separated** boundary conditions are of the form

$$T_{\alpha, \beta} f = \tau f, \quad \alpha, \beta \in [0, \pi),$$

$$f \in \text{dom}(T_{\alpha, \beta}) = \{g \in \text{dom}(T_{max}) \mid g(a)\cos(\alpha) + g^{[1]}(a)\sin(\alpha) = 0; \\ g(b)\cos(\beta) + g^{[1]}(b)\sin(\beta) = 0\}.$$

Special cases: $\alpha = 0$, $g(a) = 0$ is called the **Dirichlet** boundary condition at a ; $\alpha = \frac{\pi}{2}$, $g^{[1]}(a) = 0$ is called the **Neumann** boundary condition at a (analogous facts hold at the endpoint b).

Note. Here $g^{[1]}(a) = \lim_{x \downarrow a} p(x)g'(x)$, $g^{[1]}(b) = \lim_{x \uparrow b} p(x)g'(x)$, denote the first quasi-derivatives of g at $x = a$, resp., at $x = b$.

All Self-Adjoint B.C.'s in the Regular Case (contd.)

Theorem (contd.).

(ii) **All self-adjoint** extensions $T_{\varphi,R}$ of T_{min} with **coupled** boundary conditions are of the type

$$T_{\varphi,R}f = \tau f,$$

$$f \in \text{dom}(T_{\varphi,R}) = \left\{ g \in \text{dom}(T_{max}) \mid \begin{pmatrix} g(b) \\ g^{[1]}(b) \end{pmatrix} = e^{i\varphi} R \begin{pmatrix} g(a) \\ g^{[1]}(a) \end{pmatrix} \right\},$$

where $\varphi \in [0, 2\pi)$, and R is a real 2×2 matrix with $\det(R) = 1$ (i.e., $R \in SL(2, \mathbb{R})$).

Special cases: $\varphi = 0$, $R = I_2$, $g(b) = g(a)$, $g^{[1]}(b) = g^{[1]}(a)$ are called **periodic** boundary conditions; similarly, $\varphi = \pi$, $R = I_2$, $g(b) = -g(a)$, $g^{[1]}(b) = -g^{[1]}(a)$ are called **antiperiodic** boundary conditions.

(iii) **Every self-adjoint** extension of T_{min} is either of type (i) (i.e., **separated**) or of type (ii) (i.e., **coupled**).

This completely characterizes the regular case (standard textbook literature).

The Singular Case. Basics

Hypothesis.

Let $(a, b) \subseteq \mathbb{R}$ and suppose that p, q, r are (Lebesgue) measurable functions on (a, b) such that the following items (i)–(iii) hold:

- (i) $r > 0$ a.e. on (a, b) , $r \in L^1_{loc}((a, b); dx)$.
- (ii) $p > 0$ a.e. on (a, b) , $1/p \in L^1_{loc}((a, b); dx)$.
- (iii) q is real-valued a.e. on (a, b) , $q \in L^1_{loc}((a, b); dx)$.

Definition.

The **maximal** operator T_{max} in $L^2((a, b); rdx)$ associated with τ is defined by

$$T_{max}f = \tau f,$$

$$f \in \text{dom}(T_{max}) = \left\{ g \in L^2((a, b); rdx) \mid g, g^{[1]} \in AC_{loc}((a, b)); \right. \\ \left. \tau g \in L^2((a, b); rdx) \right\}.$$

The Singular Case. Basics (cont.)

Definition (contd.).

The **minimal** operator $T_{min,0}$ in $L^2((a, b); rdx)$ associated with τ is defined by

$$T_{min,0}f = \tau f,$$

$$f \in \text{dom}(T_{min,0}) = \left\{ g \in L^2((a, b); rdx) \mid g, g^{[1]} \in AC_{loc}((a, b)); \right. \\ \left. \text{supp}(g) \subset (a, b) \text{ is compact; } \tau g \in L^2((a, b); rdx) \right\}.$$

One can prove that $T_{min,0}$ is closable and then defines T_{min} as the closure of $T_{min,0}$, $T_{min,0} = \overline{T_{min,0}}$.

The Singular Case. Basics (cont.)

Theorem (Weyl's Alternative).

The following **alternative** holds:

- (i) For every $z \in \mathbb{C}$, all solutions u of $(\tau - z)u = 0$ are in $L^2((a, b); rdx)$ near b (resp., near a).
- (ii) For every $z \in \mathbb{C}$, there exists at least one solution u of $(\tau - z)u = 0$ which is not in $L^2((a, b); rdx)$ near b (resp., near a). In this case, for each $z \in \mathbb{C} \setminus \mathbb{R}$, there exists precisely one solution u_b (resp., u_a) of $(\tau - z)u = 0$ (up to constant multiples) which lies in $L^2((a, b); rdx)$ near b (resp., near a).

This yields the **limit circle/limit point** classification of τ at an interval endpoint:

Definition (Limit Circle/Limit Point).

In case (i) in the Theorem, τ is said to be in the **limit circle case** at b (resp., at a). (Frequently, τ is then called **quasi-regular** at b (resp., a).)

In case (ii) in the Theorem, τ is said to be in the **limit point case** at b (resp., at a).

If τ is in the **limit circle case** at a and b then τ is called **quasi-regular** on (a, b) .

All Self-Adjoint B.C.'s in the Singular Case

Theorem.

Assume that τ is in the **limit circle case** at a and b (i.e., τ is quasi-regular on (a, b)). In addition, assume that $v_j \in \text{dom}(T_{max})$, $j = 1, 2$, satisfy

$$W(\overline{v_1}, v_2)(a) = W(\overline{v_1}, v_2)(b) = 1, \quad W(\overline{v_j}, v_j)(a) = W(\overline{v_j}, v_j)(b) = 0, \quad j = 1, 2.$$

(E.g., real-valued sols. v_j , $j = 1, 2$, of $(\tau - \lambda)u = 0$ with $\lambda \in \mathbb{R}$, s.t. $W(v_1, v_2) = 1$.) For $g \in \text{dom}(T_{max})$ we introduce the **generalized boundary values**

$$\begin{aligned} \tilde{g}_1(a) &= -W(v_2, g)(a), & \tilde{g}_1(b) &= -W(v_2, g)(b), \\ \tilde{g}_2(a) &= W(v_1, g)(a), & \tilde{g}_2(b) &= W(v_1, g)(b). \end{aligned}$$

Then the following items (i)–(iii) hold:

(i) **All self-adjoint** extensions $T_{\alpha, \beta}$ of T_{min} with **separated** b.c.'s are of the form

$$T_{\alpha, \beta} f = \tau f, \quad \alpha, \beta \in [0, \pi),$$

$$f \in \text{dom}(T_{\alpha, \beta}) = \left\{ g \in \text{dom}(T_{max}) \mid \begin{aligned} &\tilde{g}_1(a)\cos(\alpha) + \tilde{g}_2(a)\sin(\alpha) = 0; \\ &\tilde{g}_1(b)\cos(\beta) + \tilde{g}_2(b)\sin(\beta) = 0 \end{aligned} \right\}.$$

All Self-Adjoint B.C.'s in the Singular Case (contd.)

Theorem (contd.).

(ii) **All self-adjoint** extensions $T_{\varphi,R}$ of T_{min} with **coupled** boundary conditions are of the type

$$T_{\varphi,R}f = \tau f,$$

$$f \in \text{dom}(T_{\varphi,R}) = \left\{ g \in \text{dom}(T_{max}) \mid \begin{pmatrix} \tilde{g}_1(b) \\ \tilde{g}_2(b) \end{pmatrix} = e^{i\varphi} R \begin{pmatrix} \tilde{g}_1(a) \\ \tilde{g}_2(a) \end{pmatrix} \right\},$$

where $\varphi \in [0, 2\pi)$, and R is a real 2×2 matrix with $\det(R) = 1$ (i.e., $R \in SL(2, \mathbb{R})$).

(iii) **Every self-adjoint** extension of T_{min} is either of type (i) (i.e., **separated**) or of type (ii) (i.e., **coupled**).

T_{min} Bounded from Below. Basics

Definition.

- (i) Fix $c \in (a, b)$ and $\lambda \in \mathbb{R}$. Then $\tau - \lambda$ is called **nonoscillatory** at a (resp., b), if every real-valued solution $u(\lambda, \cdot)$ of $\tau u = \lambda u$ has finitely many zeros in (a, c) (resp., (c, b)). Otherwise, $\tau - \lambda$ is called **oscillatory** at a (resp., b).
- (ii) Let $\lambda_0 \in \mathbb{R}$. Then T_{min} is called **bounded from below** by λ_0 , and one writes $T_{min} \geq \lambda_0 I$, if

$$(u, [T_{min} - \lambda_0 I]u)_{L^2((a,b);rdx)} \geq 0, \quad u \in \text{dom}(T_{min}).$$

The following is a key result.

Theorem.

The following items (i)–(ii) are equivalent:

- (i) T_{min} (and hence any symmetric extension of T_{min}) is **bounded from below**.
- (ii) There exists a $\nu_0 \in \mathbb{R}$ such that for all $\lambda < \nu_0$, $\tau - \lambda$ is **nonoscillatory** at a and b .

T_{min} Bounded from Below. Basics (contd.)**Definition.**

Suppose that T_{min} is **bounded from below**, and let $\lambda \in \mathbb{R}$.

(i) Then $u_a(\lambda, \cdot)$ (resp., $u_b(\lambda, \cdot)$) is called a **principal** (or **minimal**) solution of $\tau u = \lambda u$ at a (resp., b) if $u_a(\lambda, \cdot)$ and $u_b(\lambda, \cdot)$ are minimal solutions of $\tau u = \lambda u$ in the sense that

$$u(\lambda, x)^{-1} u_a(\lambda, x) = o(1) \text{ as } x \downarrow a,$$

$$u(\lambda, x)^{-1} u_b(\lambda, x) = o(1) \text{ as } x \uparrow b,$$

for any other solution $u(\lambda, \cdot)$ of $\tau u = \lambda u$ (which is nonvanishing near a , resp., b) with $W(u_a(\lambda, \cdot), u(\lambda, \cdot)) \neq 0$, respectively, $W(u_b(\lambda, \cdot), u(\lambda, \cdot)) \neq 0$.

(ii) A real-valued solution $\hat{u}_a(\lambda, \cdot)$ (resp., $\hat{u}_b(\lambda, \cdot)$) of $\tau u = \lambda u$ linearly independent of $u_a(\lambda, \cdot)$ (resp., $u_b(\lambda, \cdot)$) is called **nonprincipal** at a (resp., b).

Boundary Values if T_{min} is Bounded from Below**Theorem.**

Assume that τ is in the **limit circle case** at a and b (i.e., τ is quasi-regular on (a, b)). In addition, assume that $T_{min} \geq \lambda_0 I$ for some $\lambda_0 \in \mathbb{R}$, and denote by $u_a(\lambda_0, \cdot)$ and $\widehat{u}_a(\lambda_0, \cdot)$ (resp., $u_b(\lambda_0, \cdot)$ and $\widehat{u}_b(\lambda_0, \cdot)$) principal and nonprincipal solutions of $\tau u = \lambda_0 u$ at a (resp., b), satisfying (a normalization)

$$W(\widehat{u}_a(\lambda_0, \cdot), u_a(\lambda_0, \cdot)) = W(\widehat{u}_b(\lambda_0, \cdot), u_b(\lambda_0, \cdot)) = 1.$$

Introduce $v_j \in \text{dom}(T_{max})$, $j = 1, 2$, via

$$v_1(x) = \begin{cases} \widehat{u}_a(\lambda_0, x), & \text{for } x \text{ near } a, \\ \widehat{u}_b(\lambda_0, x), & \text{for } x \text{ near } b, \end{cases} \quad v_2(x) = \begin{cases} u_a(\lambda_0, x), & \text{for } x \text{ near } a, \\ u_b(\lambda_0, x), & \text{for } x \text{ near } b, \end{cases}$$

Boundary Values if T_{min} is Bounded from Below

Theorem (contd.).

Then one obtains for all $g \in \text{dom}(T_{max})$,

$$\begin{aligned}\tilde{g}(a) &= -W(v_2, g)(a) = \tilde{g}_1(a) = -W(u_a(\lambda_0, \cdot), g)(a) \\ &= \lim_{x \downarrow a} \frac{g(x)}{\widehat{u}_a(\lambda_0, x)},\end{aligned}\tag{*}$$

$$\begin{aligned}\tilde{g}(b) &= -W(v_2, g)(b) = \tilde{g}_1(b) = -W(u_b(\lambda_0, \cdot), g)(b) \\ &= \lim_{x \uparrow b} \frac{g(x)}{\widehat{u}_b(\lambda_0, x)},\end{aligned}\tag{**}$$

$$\begin{aligned}\tilde{g}'(a) &= W(v_1, g)(a) = \tilde{g}_2(a) = W(\widehat{u}_a(\lambda_0, \cdot), g)(a) \\ &= \lim_{x \downarrow a} \frac{g(x) - \tilde{g}(a)\widehat{u}_a(\lambda_0, x)}{u_a(\lambda_0, x)},\end{aligned}\tag{***}$$

$$\begin{aligned}\tilde{g}'(b) &= W(v_1, g)(b) = \tilde{g}_2(b) = W(\widehat{u}_b(\lambda_0, \cdot), g)(b) \\ &= \lim_{x \uparrow b} \frac{g(x) - \tilde{g}(b)\widehat{u}_b(\lambda_0, x)}{u_b(\lambda_0, x)}.\end{aligned}\tag{****}$$

In particular, the limits on the right-hand sides in (*) – (****) exist.

The Friedrichs Extension

The Friedrichs extension is characterized in the expected manner:

Theorem (Niessen and Zettl 1992).

Assume that τ is in the **limit circle case** at a and b (i.e., τ is quasi-regular on (a, b)). In addition, assume that $T_{min} \geq \lambda_0 I$ for some $\lambda_0 \in \mathbb{R}$. Then the **Friedrichs extension** T_F of T_{min} is characterized by

$$T_F f = \tau f, \quad f \in \text{dom}(T_F) = \{g \in \text{dom}(T_{max}) \mid \tilde{g}(a) = \tilde{g}(b) = 0\}.$$

We recall,

$$\tilde{g}(a) = \lim_{x \downarrow a} \frac{g(x)}{\hat{u}_a(\lambda_0, x)}, \quad \tilde{g}(b) = \lim_{x \uparrow b} \frac{g(x)}{\hat{u}_b(\lambda_0, x)}.$$

One can now express Weyl–Titchmarsh m -functions directly in terms of the boundary values \tilde{g} , \tilde{g}' , but this needs a few preparations:

Basics of m -function Theory

In the singular Sturm–Liouville operator case this is a bit more involved!
First, one needs a (rather benign) spectral hypothesis:

Spectral Hypothesis

In addition to the standard assumptions on p, q, r , suppose that for some (and hence for all) $c \in (a, b)$, the self-adjoint operator $T_{\alpha_0, 0, a, c}$ in $L^2((a, c); rdx)$, associated with $\tau|_{(a, c)}$ and a **Dirichlet boundary condition** at c (i.e., $g(c) = 0$, $g \in \text{dom}(T_{\max, a, c})$, the maximal operator associated with $\tau|_{(a, c)}$ in $L^2((a, c); rdx)$), has **purely discrete spectrum**.

This Hypothesis is equivalent to the existence of an **entire** solution $\phi_{\alpha_0}(z, \cdot)$ of $\tau u = zu$, $z \in \mathbb{C}$, that is **real-valued** for $z \in \mathbb{R}$, and lies in $\text{dom}(T_{\alpha_0, \beta_0})$ near the point a . In particular, $\phi_{\alpha_0}(z, \cdot)$ satisfies the **boundary condition** indexed by α_0 at the left endpoint a if τ is in the **limit circle case** at a , and $\phi_{\alpha_0}(z, \cdot) \in L^2((a, c); rdx)$ if τ is in the **limit point case** at a . In addition, a second, linearly independent **entire** solution $\theta_{\alpha_0}(z, \cdot)$ of $\tau u = zu$ exists, with $\theta_{\alpha_0}(z, \cdot)$ **real-valued** for $z \in \mathbb{R}$, satisfying (the normalization)

$$W(\theta_{\alpha_0}(z, \cdot), \phi_{\alpha_0}(z, \cdot)) = 1, \quad z \in \mathbb{C}.$$

Basics of m -function Theory (contd.)

We note that $\phi_{\alpha_0}(z, \cdot)$ is unique up to a nonvanishing entire factor (real on the real line) with respect to $z \in \mathbb{C}$. Hence, we may normalize $\phi_{\alpha_0}(z, \cdot)$ such that

$$\tilde{\phi}_{\alpha_0}(z, a) = -\sin(\alpha_0), \quad \tilde{\phi}'_{\alpha_0}(z, a) = \cos(\alpha_0), \quad z \in \mathbb{C},$$

and thus,

$$\tilde{\theta}_{\alpha_0}(z, a) = \cos(\alpha_0), \quad \tilde{\theta}'_{\alpha_0}(z, a) = \sin(\alpha_0), \quad z \in \mathbb{C},$$

Given (for $z \in \mathbb{C} \setminus \mathbb{R}$),

$$\psi_{\beta_0,+}(z, \cdot) = \theta_0(z, \cdot) + m_{0,\beta_0}(z)\phi_0(z, \cdot) \begin{cases} \text{satisfies the } \mathbf{b.c.} \text{ at } x = b \\ \text{if } \tau \text{ is } \mathbf{l.c.c.} \text{ at } b, \\ \in L^2((a, b); r(x)dx) \text{ if } \tau \text{ is } \mathbf{l.p.c.} \text{ at } b, \end{cases}$$

one verifies that the **Dirichlet m -function**, where $\alpha = 0$, can be computed via

$$m_{0,\beta_0}(z) = \tilde{\psi}'_{0,\beta_0}(z, a) / \tilde{\psi}_{0,\beta_0}(z, a), \quad z \in \mathbb{C} \setminus \mathbb{R}.$$

For other (i.e., **non-Dirichlet**) **b.c.'s**, where $\alpha_0 \neq 0$, use the usual linear fractional transformations (keep β_0 , the **b.c.** at $x = b$, fixed).

Here $\tilde{\psi}, \tilde{\psi}'$ denote precisely the generalized boundary values we introduced before.

The Bessel Operator on $(0, \infty)$

Example (Bessel Operator).

Let $a = 0$, $b = \infty$,

$$p(x) = r(x) = 1, \quad q(x) := q_\gamma(x) = \frac{\gamma^2 - (1/4)}{x^2}, \quad \gamma \in [0, 1), \quad x \in (0, \infty).$$

Then $\tau_\gamma = -d^2/dx^2 + [\gamma^2 - (1/4)]x^{-2}$, $\gamma \in [0, 1)$, $x \in (0, \infty)$, is in the limit circle case at the endpoint 0 and in the limit point case at ∞ . It suffices to focus on the generalized boundary values at the singular endpoint $x = 0$. To this end we introduce principal and nonprincipal solutions $u_{0,\gamma}(0, \cdot)$ and $\widehat{u}_{0,\gamma}(0, \cdot)$ of $\tau_\gamma u = 0$ by

$$u_{0,\gamma}(0, x) = x^{(1/2)+\gamma}, \quad \gamma \in [0, 1), \quad x \in (0, \infty),$$

$$\widehat{u}_{0,\gamma}(0, x) = \begin{cases} (2\gamma)^{-1}x^{(1/2)-\gamma}, & \gamma \in (0, 1), \\ x^{1/2}\ln(1/x), & \gamma = 0; \end{cases} \quad x \in (0, \infty).$$

The Bessel Operator on $(0, \infty)$ (contd.)

Example (Bessel Operator (contd.)).

The generalized boundary values for $g \in \text{dom}(T_{\max, \gamma})$ (the maximal operator associated with τ_γ) are then of the form

$$\begin{aligned} \tilde{g}(0) &= -W(u_{0, \gamma}(0, \cdot), g)(0) \\ &= \begin{cases} \lim_{x \downarrow 0} g(x) / [(2\gamma)^{-1} x^{(1/2) - \gamma}], & \gamma \in (0, 1), \\ \lim_{x \downarrow 0} g(x) / [x^{1/2} \ln(1/x)], & \gamma = 0, \end{cases} \end{aligned}$$

$$\begin{aligned} \tilde{g}'(0) &= W(\hat{u}_{0, \gamma}(0, \cdot), g)(0) \\ &= \begin{cases} \lim_{x \downarrow 0} [g(x) - \tilde{g}(0)(2\gamma)^{-1} x^{(1/2) - \gamma}] / x^{(1/2) + \gamma}, & \gamma \in (0, 1), \\ \lim_{x \downarrow 0} [g(x) - \tilde{g}(0)x^{1/2} \ln(1/x)] / x^{1/2}, & \gamma = 0. \end{cases} \end{aligned}$$

The Bessel Operator on $(0, \infty)$ (contd.)

Theorem (Bessel operator m -function)

For the (Dirichlet-type) m -function one obtains the **Nevanlinna–Herglotz** fct.

$$m_0(z; \gamma) = \begin{cases} -e^{-i\pi\gamma} 2^{-2\gamma-1} \gamma^{-1} [\Gamma(1-\gamma)/\Gamma(1+\gamma)] z^\gamma, & \gamma \in (0, 1), \\ i(\pi/2) + \ln(2) - \gamma_E - 2^{-1} \ln(z), & \gamma = 0, \end{cases}$$

$$z \in \mathbb{C} \setminus [0, \infty).$$

Here $\gamma_E = 0.57721\dots$ represents Euler's constant, and $\Gamma(\cdot)$ is the Gamma fct.

Theorem (Bessel operator m -function, contd.)

In the limit point case where $\gamma \geq 1$, one obtains

$$m_0(z; \gamma) = \begin{cases} -C_\gamma e^{-i\pi\gamma} (2/\pi) \sin(\pi\gamma) z^\gamma, & \gamma \in [1, \infty) \setminus \mathbb{N}, \\ C_0 (2/\pi) z^n [i - (1/\pi) \ln(z)], & \gamma \in \mathbb{N}, \end{cases} \quad z \in \mathbb{C} \setminus [0, \infty).$$

Thus, the limit point case, $\gamma \geq 1$ naturally leads to a **generalized Nevanlinna–Herglotz** function $m_0(\cdot; \gamma)$.

The Legendre Operator on $(-1, 1)$

Example (Legendre Operator).

Let $a = -1$, $b = 1$,

$$p(x) = 1 - x^2, \quad r(x) = 1, \quad q(x) = 0, \quad x \in (-1, 1).$$

Then $\tau_L = -(d/dx)(1 - x^2)(d/dx)$, $x \in (-1, 1)$, is in the limit circle case and singular at both endpoints ± 1 . Principal and nonprincipal solutions $u_{\pm 1, L}(0, \cdot)$ and $\hat{u}_{\pm 1, L}(0, \cdot)$ of $\tau_L u = 0$ at ± 1 are then given by

$$u_{\pm 1, L}(0, x) = 1, \quad \hat{u}_{\pm 1, L}(0, x) = 2^{-1} \ln((1 - x)/(1 + x)), \quad x \in (-1, 1).$$

The **generalized boundary values** for $g \in \text{dom}(T_{max, L})$ (the maximal operator associated with τ_L) are then of the form

$$\begin{aligned} \tilde{g}(\pm 1) &= -W(u_{\pm 1, L}(0, \cdot), g)(\pm 1) \\ &= -(pg')(\pm 1) = \lim_{x \rightarrow \pm 1} g(x) / [2^{-1} \ln((1 - x)/(1 + x))], \\ \tilde{g}'(\pm 1) &= W(\hat{u}_{\pm 1, L}(0, \cdot), g)(\pm 1) \\ &= \lim_{x \rightarrow \pm 1} [g(x) - \tilde{g}(\pm 1) 2^{-1} \ln((1 - x)/(1 + x))]. \end{aligned}$$

The Legendre Operator on $(-1, 1)$ (contd.)

One observes the curious fact that the **Friedrichs** extension $T_{F,L}$ of $T_{min,L}$ (the minimal operator associated with τ_L) then satisfies the boundary conditions

$$(pg')(-1) = (pg')(1) = 0,$$

which resembles the **Neumann** (and **not** the **Dirichlet**) boundary conditions in the context of a regular Sturm–Liouville differential expression on the interval $[-1, 1]$. However, since τ_L is singular at both endpoints ± 1 , this represents no conundrum.

While this is well-known to experts, I will not lie, this fact served as one of the prime motivations to write our paper on this topic!

In addition, we note that the spectrum of $T_{F,L}$ may be computed explicitly,

$$\sigma(T_{F,L}) = \{n^2 - n\}_{n \in \mathbb{N}}.$$

The Legendre Operator on $(-1, 1)$ (contd.)

Theorem (Legendre operator m -function)

For the (Dirichlet-type) m -function one obtains the **Nevanlinna–Herglotz** fct.

$$m_{0,L}(z) = -\frac{1}{2} \left[\pi \cot(\nu(z)\pi) + \gamma_E + 2\psi(1 + \nu(z)) \right], \quad z \in \rho(T_{F,L}),$$

where we abbreviated

$$\nu(z) := 2^{-1} \left[-1 + (1 + 4z)^{1/2} \right],$$

and where

$$\psi(z) = \Gamma'(z)/\Gamma(z), \quad z \in \mathbb{C} \setminus \mathbb{N}_0,$$

denotes the Digamma function.

Now prove from scratch that this is indeed a Nevanlinna–Herglotz fct.!!!!

Note. $\nu(z) := 2^{-1} \left[-1 + (1 + 4z)^{1/2} \right]$ is indeed a **Nevanlinna–Herglotz** function.

The Legendre Operator on $(-1, 1)$ (contd.)

Proving this Nevanlinna–Herglotz property of $m_{0,L}(z)$ is tricky: Consider

$$-\pi \cot(z\pi) = \sum_{n \in \mathbb{Z}} \left[\frac{1}{n-z} - \frac{n\pi^2}{n^2\pi^2+1} \right], \quad z \in \mathbb{C} \setminus \mathbb{Z},$$

and

$$\psi(1+z) = -\gamma_E + \sum_{n \in \mathbb{N}} \left[\frac{1}{n} - \frac{1}{n+z} \right], \quad z \in \mathbb{C} \setminus (-\mathbb{N})$$

Then

$$\begin{aligned} m_{0,L}(z) &= -(\pi/2) \cot(\nu(z)\pi) - \gamma_E - \psi(1 + \nu(z)) \\ &= \frac{1}{2} \sum_{n \in \mathbb{Z}} \left[\frac{1}{n - \nu(z)} - \frac{n\pi^2}{n^2\pi^2 + 1} \right] + \sum_{n \in \mathbb{N}} \left[\frac{1}{n + \nu(z)} - \frac{1}{n} \right] \end{aligned}$$

The Legendre Operator on $(-1, 1)$ (contd.)

$$\begin{aligned}
 &= -\frac{1}{2\nu(z)} + \frac{1}{2} \sum_{n \in \mathbb{N}} \left[\frac{1}{n - \nu(z)} - \frac{n\pi^2}{n^2\pi^2 + 1} \right] \\
 &\quad + \frac{1}{2} \sum_{n \in \mathbb{N}} \left[\frac{-1}{n + \nu(z)} + \frac{n\pi^2}{n^2\pi^2 + 1} \right] \\
 &\quad + \frac{1}{2} \sum_{n \in \mathbb{N}} \left[\frac{1}{n + \nu(z)} - \frac{1}{n} \right] + \frac{1}{2} \sum_{n \in \mathbb{N}} \left[\frac{1}{n + \nu(z)} - \frac{1}{n} \right] \\
 &= -\frac{1}{2\nu(z)} + \frac{1}{2} \sum_{n \in \mathbb{N}} \left[\frac{1}{n - \nu(z)} - \frac{1}{n} \right] + \frac{1}{2} \sum_{n \in \mathbb{N}} \left[\frac{1}{n + \nu(z)} - \frac{1}{n} \right] \\
 &= -\frac{1}{2\nu(z)} + \frac{1}{2} \sum_{n \in \mathbb{N}} \left[\frac{1}{n - \nu(z)} - \frac{1}{n} \right] + \frac{1}{2} \left[\frac{1}{1 + \nu(z)} - 1 \right] \\
 &\quad + \frac{1}{2} \sum_{n=2}^{\infty} \left[\frac{1}{n + \nu(z)} - \frac{1}{n} \right]
 \end{aligned}$$

The Legendre Operator on $(-1, 1)$ (contd.)

$$\begin{aligned}
 &= -\frac{1}{2} \frac{1}{\nu(z)[\nu(z)+1]} - \frac{1}{2} + \frac{1}{2} \sum_{n \in \mathbb{N}} \left[\frac{1}{n - \nu(z)} - \frac{1}{n} \right] \\
 &\quad + \frac{1}{2} \sum_{n \in \mathbb{N}} \left[\frac{1}{n+1 + \nu(z)} - \frac{1}{n} \right] + \frac{1}{2} \sum_{n \in \mathbb{N}} \left[\frac{1}{n} - \frac{1}{n+1} \right] \\
 &= -\frac{1}{2z} - \frac{1}{2} + \frac{1}{2} \sum_{n \in \mathbb{N}} \frac{1}{n(n+1)} \\
 &\quad + \frac{1}{2} \sum_{n \in \mathbb{N}} \left[\frac{1}{n+2^{-1} - 2^{-1}(1+4z)^{1/2}} - \frac{1}{n} \right] \\
 &\quad + \frac{1}{2} \sum_{n \in \mathbb{N}} \left[\frac{1}{n+2^{-1} + 2^{-1}(1+4z)^{1/2}} - \frac{1}{n} \right] \\
 &= -\frac{1}{2z} + \sum_{n \in \mathbb{N}} \left[\frac{n+2^{-1}}{(n+2^{-1})^2 - 4^{-1} - z} - \frac{1}{n} \right] \\
 &= -\frac{1}{2z} + \sum_{n \in \mathbb{Z}} \left[\frac{n+2^{-1}}{n(n+1) - z} - \frac{1}{n} \right], \quad z \in \mathbb{Z} \setminus \{n(n+1)\}_{n \in \mathbb{N}_0}.
 \end{aligned}$$

The Legendre Operator on $(-1, 1)$ (contd.)

Here we used

$$\sum_{n \in \mathbb{N}} \frac{1}{n(n+1)} = 1,$$

and

$$\nu(z)[\nu(z) + 1] = z, \quad z \in \mathbb{C}.$$

Once again, one confirms explicitly that the set of poles of $m_{0,L}(\cdot)$ coincides with the spectrum of $T_{F,L}$,

$$\sigma(T_{F,L}) = \{n^2 - n\}_{n \in \mathbb{N}}.$$

The Laguerre (resp., Kummer, or Confluent Hypergeometric) Operator on $(0, \infty)$

Example (Laguerre Operator).

Let $a = 0$, $b = \infty$,

$$p(x) = p_\beta(x) = x^\beta e^{-x}, \quad q(x) = 0, \quad r(x) := r_\beta(x) = x^{\beta-1} e^{-x}, \\ \beta \in (0, 2), \quad x \in (0, \infty).$$

Then $\tau_\beta = -x^{1-\beta} e^x \frac{d}{dx} x^\beta e^{-x} \frac{d}{dx}$, $x \in (0, \infty)$, and the underlying Hilbert space is $L^2((0, \infty); x^{\beta-1} e^{-x} dx)$. At $x = 0$, τ_β is regular for $\beta \in (0, 1)$ and singular for $\beta \in [1, 2)$.

For $z \in \mathbb{C}$, solutions to the Kummer equation $\tau_\beta y = zy$ are given by

$$y_{1,\beta}(z, x) = F(-z, \beta; x), \quad \beta \in (0, 2), \quad z \in \mathbb{C}, \quad x \in (0, \infty), \\ y_{2,\beta}(z, x) = \begin{cases} x^{1-\beta} F(1 - \beta - z, 2 - \beta; x), & \beta \in (0, 2) \setminus \{1\}, \quad z \in \mathbb{C}, \\ \Gamma(-z) U(-z, 1; x), & \beta = 1, \quad z \in \mathbb{C} \setminus \{0\}, \\ -\int_1^x dt t^{-1} e^t, & \beta = 1, \quad z = 0; \quad x \in (0, \infty), \end{cases}$$

The Laguerre Operator on $(0, \infty)$ (contd.)

Example (Laguerre Operator (contd.)).

where $F(\cdot, \cdot; \cdot)$ (also frequently denoted by ${}_1F_1(\cdot, \cdot; \cdot)$ or $M(\cdot, \cdot; \cdot)$) denotes the confluent hypergeometric function and $U(\cdot, 1; \cdot)$ represents an associated logarithmic case.

A principal solution of $\tau_\beta u = \lambda u$, $\lambda \leq 0$, at $x = 0$ is given by

$$u_{0,\beta}(\lambda, \cdot) = \begin{cases} (1 - \beta)^{-1} y_{2,\beta}(\lambda, \cdot), & \beta \in (0, 1), \\ -(1 - \beta)^{-1} y_{1,\beta}(\lambda, \cdot), & \beta \in (1, 2), \\ y_{1,1}(\lambda, \cdot), & \beta = 1, \end{cases} \quad \lambda \leq 0, \quad (7.1)$$

and a nonprincipal solution of $\tau_\beta u = \lambda u$ at $x = 0$ is given by

$$\widehat{u}_{0,\beta}(\lambda, \cdot) = \begin{cases} y_{1,\beta}(\lambda, \cdot), & \beta \in (0, 1), \\ y_{2,\beta}(\lambda, \cdot), & \beta \in [1, 2). \end{cases} \quad \lambda \leq 0. \quad (7.2)$$

The Laguerre Operator on $(0, \infty)$ (contd.)

Example (Laguerre Operator (contd.)).

The generalized boundary values for $g \in \text{dom}(T_{\max, \beta})$ (the maximal operator associated with τ_β) are then of the form

$$\tilde{g}(0) = -W(u_{0, \beta}(0, \cdot), g)(0) = \lim_{x \downarrow 0} \frac{g(x)}{\hat{u}_{0, \beta}(0, x)} = \begin{cases} g(0), & \beta \in (0, 1), \\ \lim_{x \downarrow 0} \frac{g(x)}{x^{1-\beta}}, & \beta \in (1, 2), \\ \lim_{x \downarrow 0} \frac{g(x)}{[-\ln(x)]}, & \beta = 1, \end{cases}$$

$$\tilde{g}'(0) = W(\hat{u}_{0, \beta}(0, \cdot), g)(0) = \lim_{x \downarrow 0} \frac{g(x) - \tilde{g}(0)\hat{u}_{0, \beta}(0, x)}{u_{0, \beta}(0, x)}$$

$$= \begin{cases} \lim_{x \downarrow 0} \frac{g(x) - g(0)}{(1-\beta)^{-1}x^{1-\beta}} = \frac{0}{0} = \lim_{x \downarrow 0} \frac{g'(x)}{x^{-\beta}} = g^{[1]}(0), & \beta \in (0, 1), \\ (\beta - 1) \lim_{x \downarrow 0} [g(x) - \tilde{g}(0)x^{1-\beta}], & \beta \in (1, 2), \\ \lim_{x \downarrow 0} \{g(x) - \tilde{g}(0)[- \ln(x)]\}, & \beta = 1. \end{cases}$$

The Laguerre Operator on $(0, \infty)$ (contd.)

Theorem (Laguerre operator m -function)

For the (Dirichlet-type) m -function one obtains the **Nevanlinna–Herglotz** fct.

$$m_{0,\beta}(z) = \begin{cases} \frac{(1-\beta)\Gamma(2-\beta)\Gamma(-z)}{\Gamma(\beta)\Gamma(1-\beta-z)}, & \beta \in (1, 2), z \in \rho(T_{F,\beta}), \\ -\psi(-z), & \beta = 1, z \in \rho(T_{F,1}). \end{cases}$$

Once again, here

$$\psi(z) = \Gamma'(z)/\Gamma(z), \quad z \in \mathbb{C} \setminus \mathbb{N}_0,$$

denotes the Digamma function.

We recall

$$\sigma(T_{F,\beta}) = \begin{cases} \{n+1-\beta\}_{n \in \mathbb{N}_0}, & \beta \in (0, 1), \\ \mathbb{N}_0, & \beta \in [1, 2]. \end{cases}$$